

Diffraction from non-crystalline materials

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Stefan U. Egelhaaf
Condensed Matter Physics Laboratory
Heinrich Heine University
Düsseldorf, Germany

17th Oxford School of Neutron Scattering
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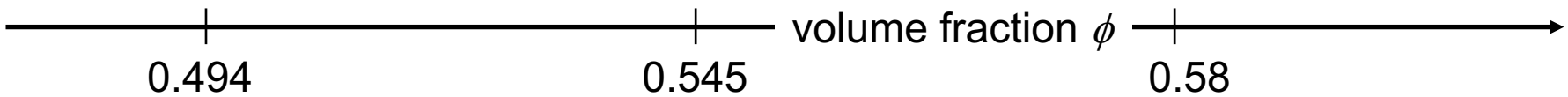
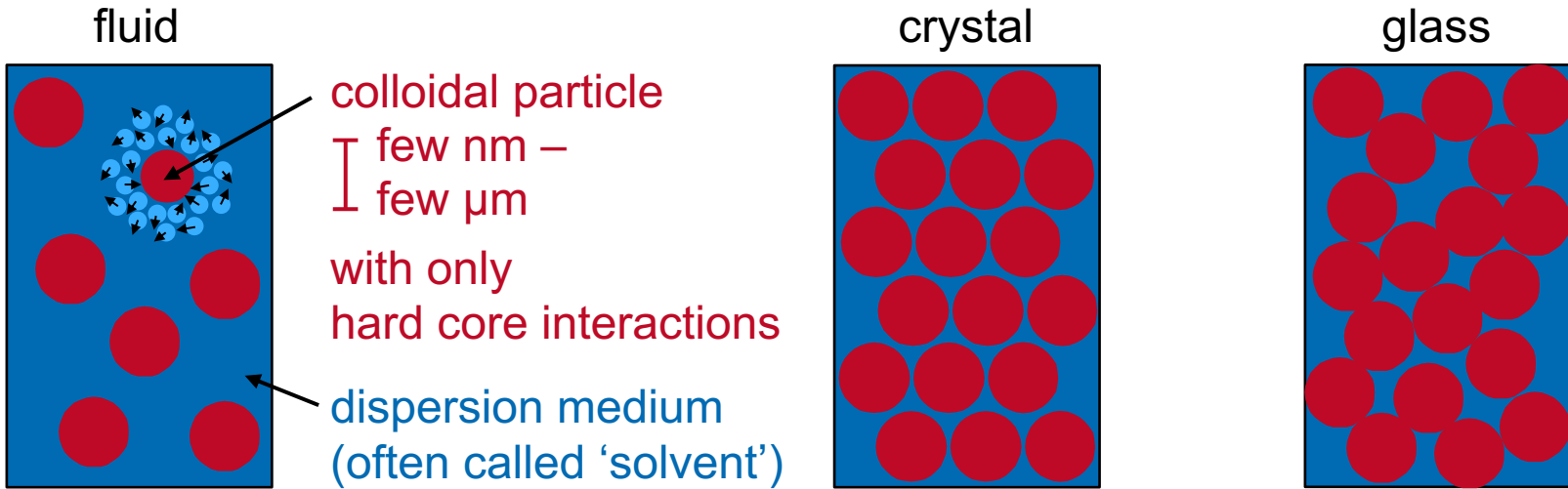
examples of non-crystalline materials

characterization of non-crystalline materials

diffraction from non-crystalline materials

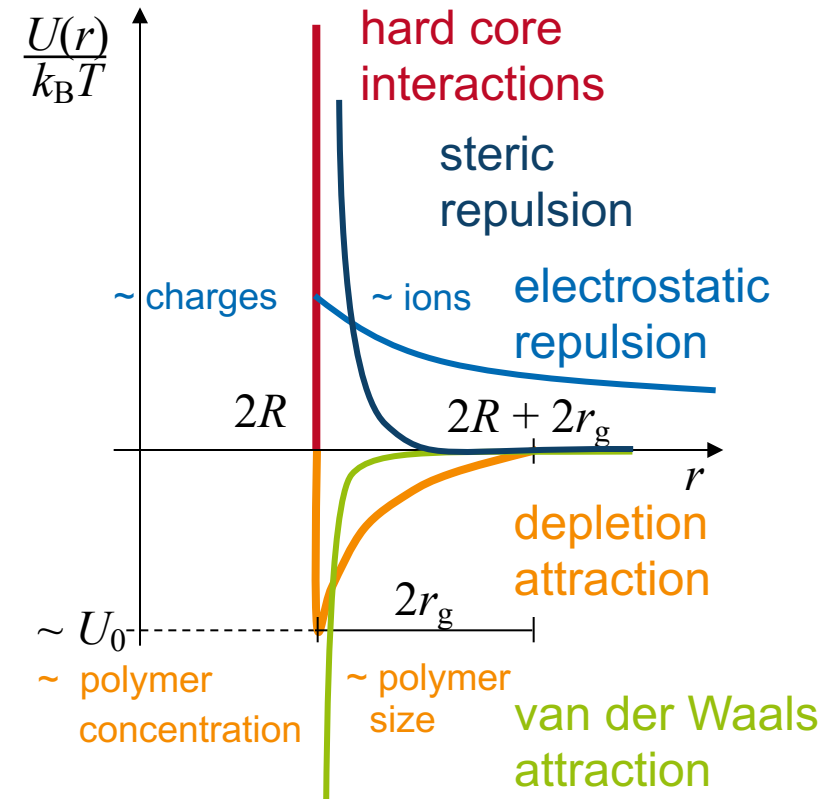
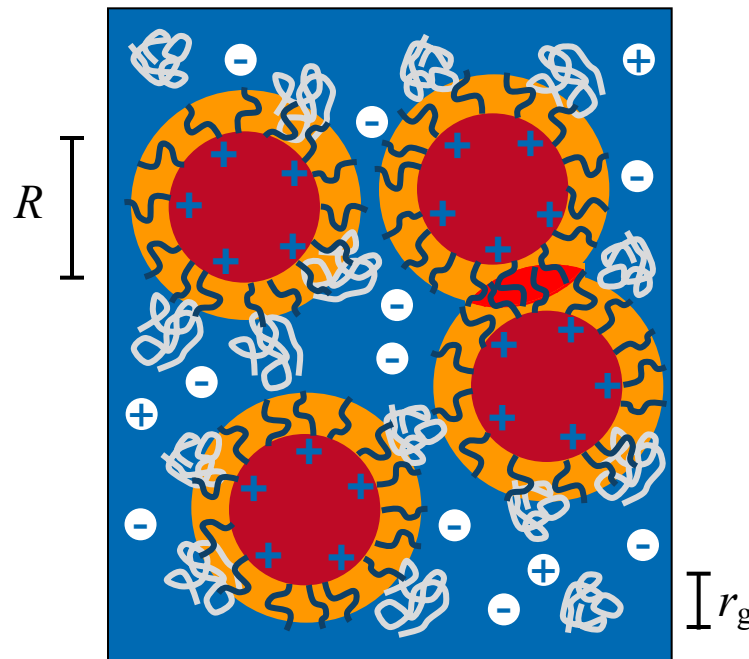
- general concept
- form factor – particle properties
- structure factor – particle arrangements
- ‘prefactor’ – contrast variation

Colloidal Suspensions



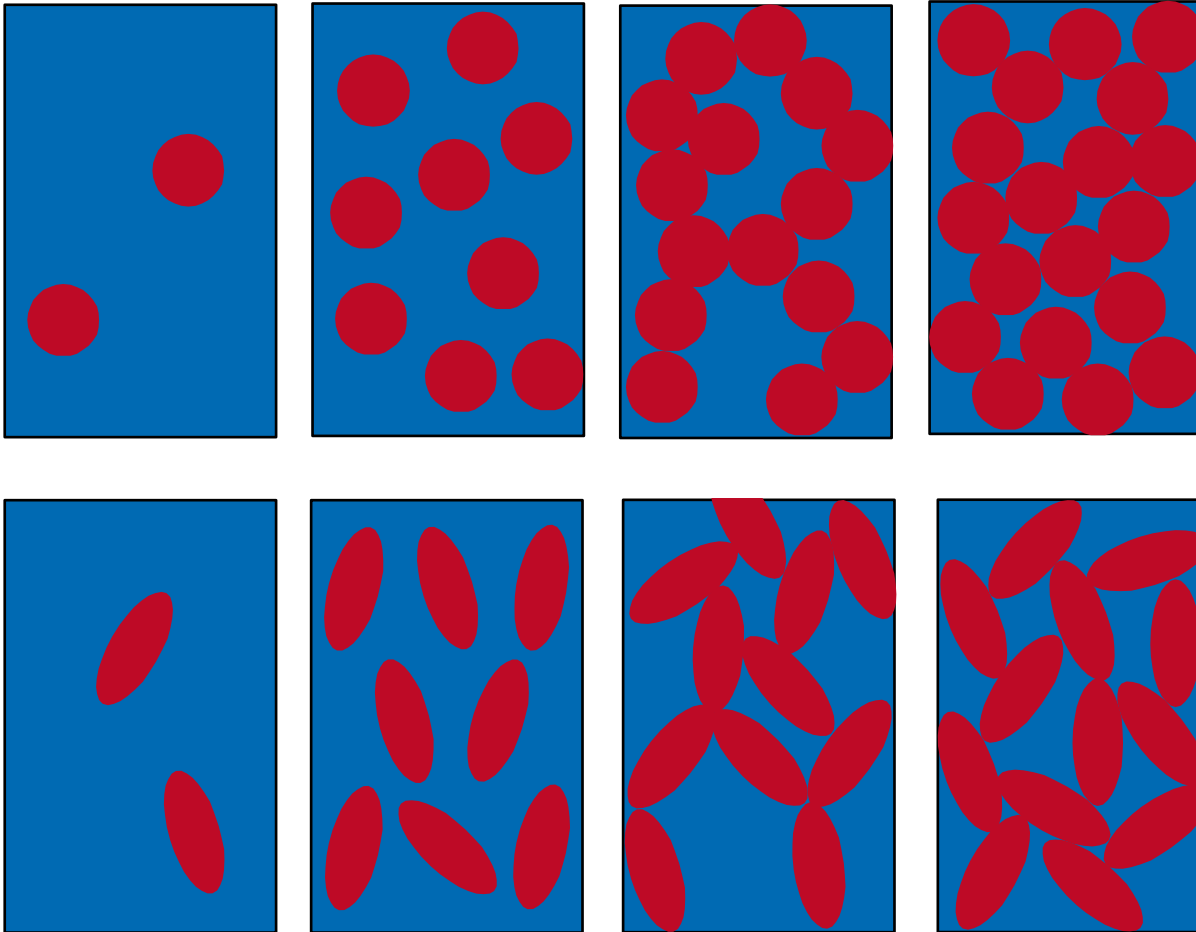
equilibrium

non-equilibrium

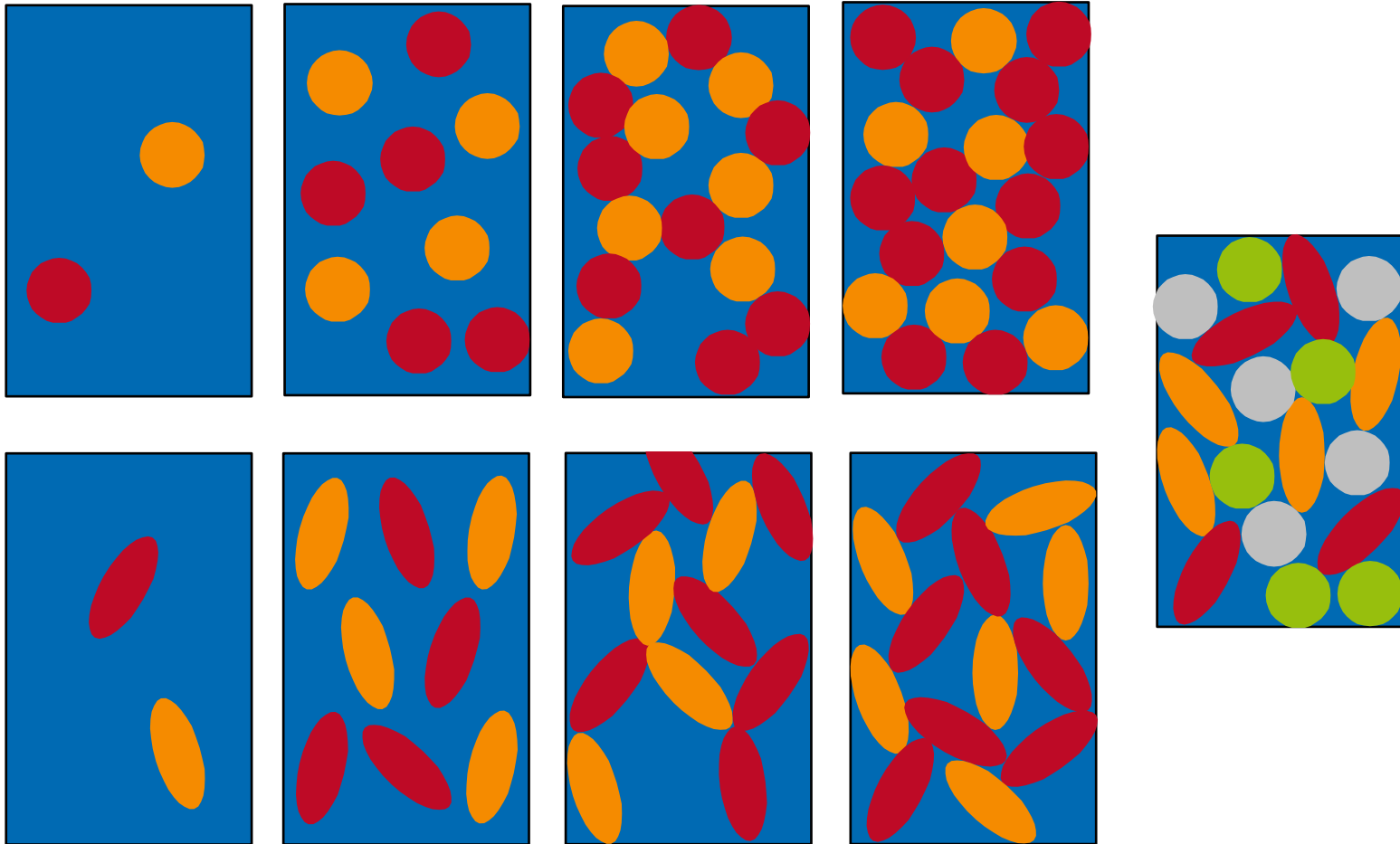


as well as anisotropic interactions,
(e.g. dipol or 'patchy' interactions)

Non-crystalline materials



Non-crystalline materials



examples of non-crystalline materials

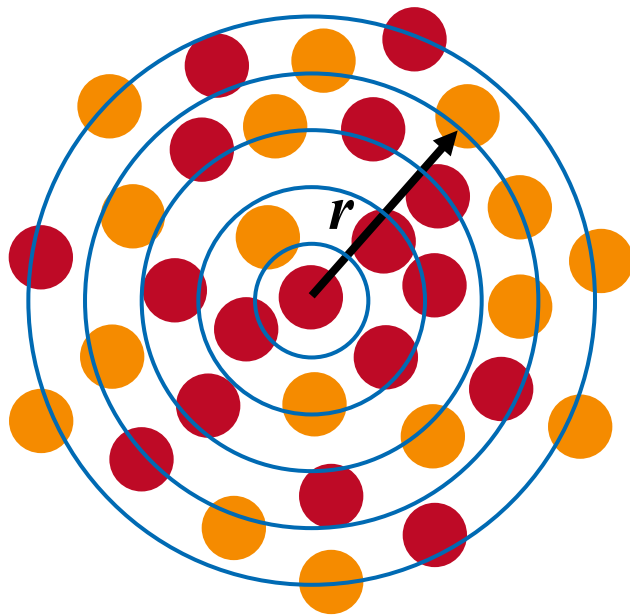
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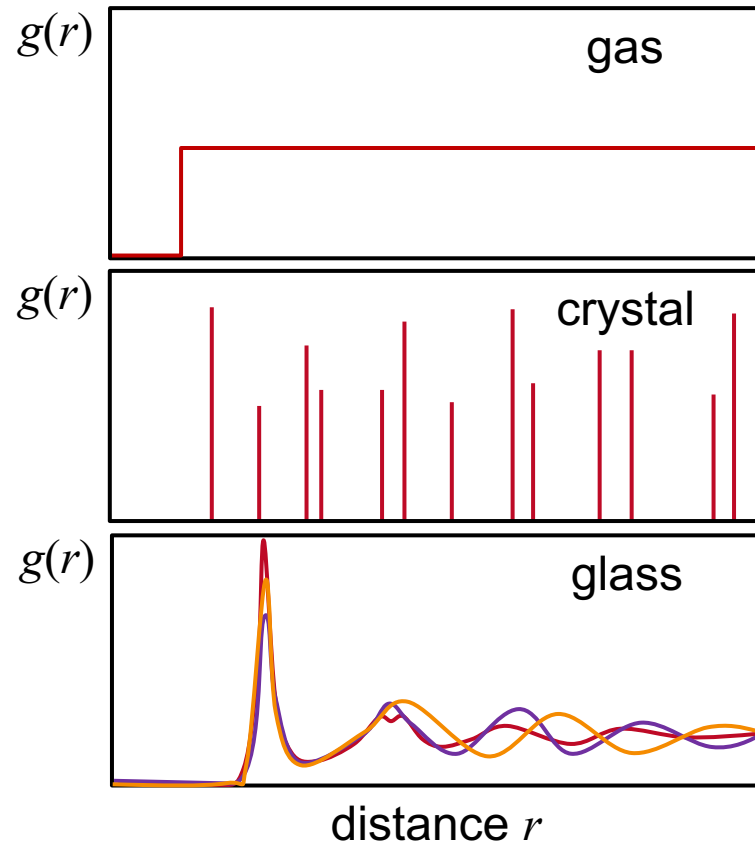
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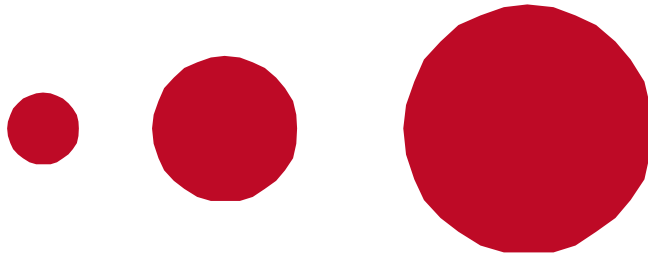
radial distribution function $g(r)$

$(N/V) g(r) d^3r$ is the number of particles in d^3r at r



contains information on interactions





particle size



particle shape



particle structure

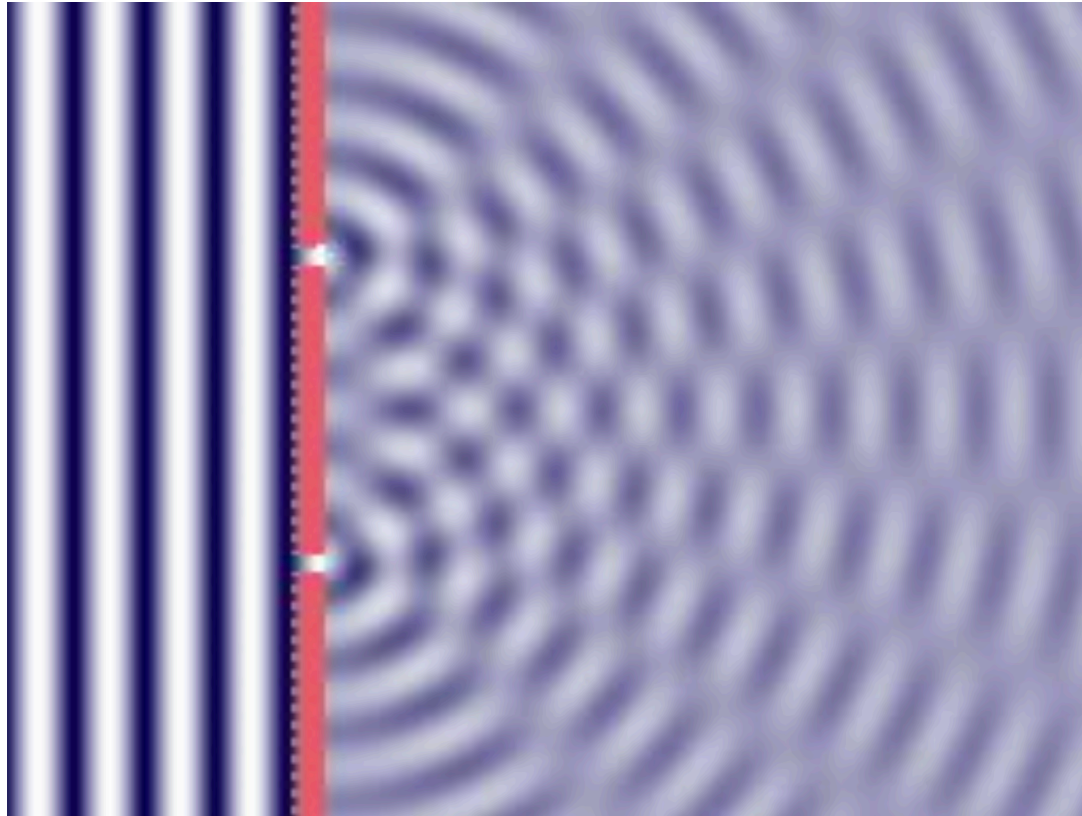
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Two-slit experiment



particle = ensemble of small volume elements dV (using Born approx.)

$$\psi_p^{sc} = \int_V e^{i\delta\phi} d\psi^{sc} \sim \int_V e^{i\mathbf{Q}\cdot\mathbf{r}} \Delta b(\mathbf{r}) dV$$

Fourier transform

$$\delta\phi = \frac{2\pi}{\lambda} \delta L = \mathbf{k}_i \cdot \mathbf{r} - \mathbf{k}_f \cdot \mathbf{r} = \mathbf{Q} \cdot \mathbf{r}$$

$$d\psi^{sc} \sim \Delta b(\mathbf{r}) dV \quad \text{scatt. length density } b(\mathbf{r})$$

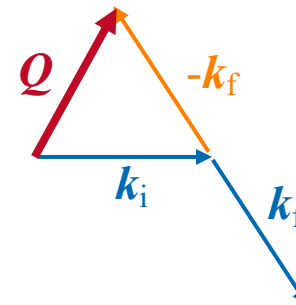
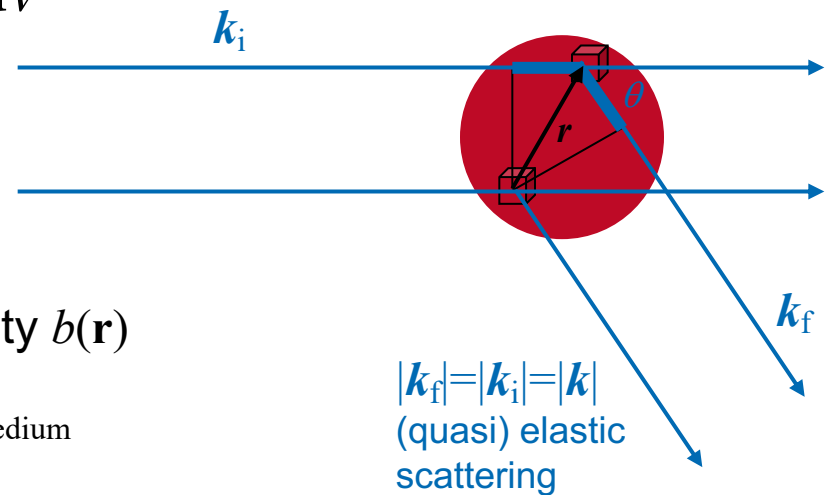
$$\Delta b(\mathbf{r}) = b_{\text{part}}(\mathbf{r}) - b_{\text{medium}}$$

with scattering vector $\mathbf{Q} = \mathbf{k}_i - \mathbf{k}_f$

$$|\mathbf{Q}| = 2k \sin \frac{\theta}{2} = \frac{4\pi}{\lambda} \sin \frac{\theta}{2} = \frac{4\pi n}{\lambda_0} \sin \frac{\theta}{2}$$

$$|\mathbf{Q}| \sim (\text{length scale})^{-1}$$

$$\delta L = \delta L_1 + \delta L_2$$



many particles = ensemble of particles $j=1..N$ (using Born approx.)

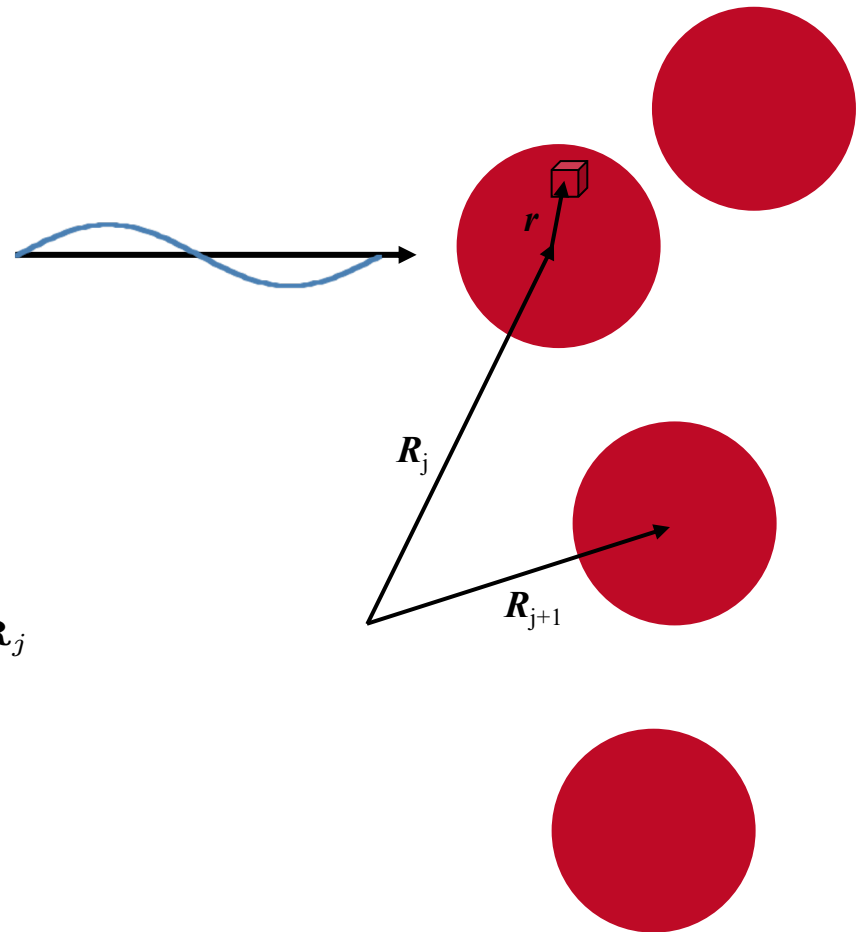
$$\psi^{\text{sc}} = \sum_j \psi_{\text{p},j}^{\text{sc}} e^{i\mathbf{Q}\cdot\mathbf{R}_j}$$

'Fourier transform'

$$\psi_{\text{p},j}^{\text{sc}} \sim \int_{V_j} \Delta b(\mathbf{r}) e^{i\mathbf{Q}\cdot\mathbf{r}} dV$$

Fourier transform

$$\begin{aligned} \psi^{\text{sc}} &\sim \sum_j \int_{V_j} \Delta b(\mathbf{r}) e^{i\mathbf{Q}\cdot\mathbf{r}} dV e^{i\mathbf{Q}\cdot\mathbf{R}_j} \\ &\sim \sum_j \int_{V_j} \Delta b(\mathbf{r}) e^{i\mathbf{Q}\cdot(\mathbf{r}+\mathbf{R}_j)} dV \end{aligned}$$

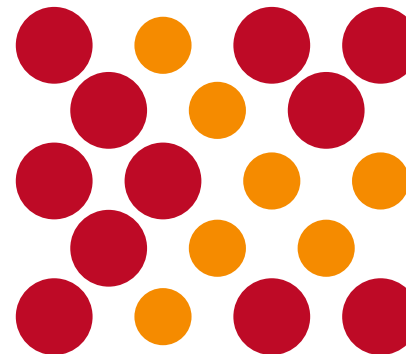
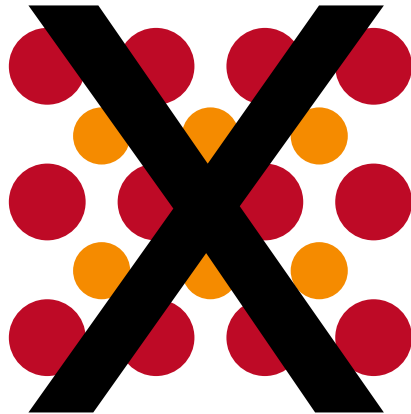


time-averaged (or ensemble-averaged) differential cross section is detected

$$\frac{d\sigma}{d\Omega}(\mathbf{Q}) \sim |\psi^{\text{sc}}|^2 \sim \sum_k \sum_j \left\langle \psi_{\mathbf{p},j}^{\text{sc}}(\mathbf{Q}) \psi_{\mathbf{p},k}^{\text{sc}*}(\mathbf{Q}) e^{i\mathbf{Q}\cdot(\mathbf{R}_j - \mathbf{R}_k)} \right\rangle$$

assumption: all particles are identical, i.e. $\psi_{\mathbf{p},j}^{\text{sc}}(\mathbf{Q}) = \psi_{\mathbf{p},k}^{\text{sc}}(\mathbf{Q}) = \psi_{\mathbf{p}}^{\text{sc}}(\mathbf{Q})$

(important: particle properties must not be linked to their positions)



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$$\frac{d\sigma}{d\Omega}(\mathbf{Q}) \sim |\psi_{\text{p}}^{\text{sc}}(\mathbf{Q})|^2 \sum_k \sum_j \left\langle e^{i\mathbf{Q}\cdot(\mathbf{R}_j(t) - \mathbf{R}_k(t))} \right\rangle$$

$$\frac{d\sigma}{d\Omega}(\mathbf{Q}) \sim N |\psi_{\text{p}}^{\text{sc}}(0)|^2 \frac{|\psi_{\text{p}}^{\text{sc}}(\mathbf{Q})|^2}{|\psi_{\text{p}}^{\text{sc}}(0)|^2} \frac{1}{N} \sum_k \sum_j \left\langle e^{i\mathbf{Q}\cdot(\mathbf{R}_j(t) - \mathbf{R}_k(t))} \right\rangle$$

'prefactor'
scattering
by N particles
(random walk)

$= P(\mathbf{q})$
form factor
(intraparticle)

$= S(\mathbf{q})$
structure factor
(interparticle)

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intraparticle interference
depends on particle size and shape

$$P(\mathbf{Q}) = \frac{|\psi_{\text{p}}^{\text{sc}}(\mathbf{Q})|^2}{|\psi_{\text{p}}^{\text{sc}}(0)|^2}$$

with $\psi_{\text{p}}^{\text{sc}}(\mathbf{Q}) \sim \int_V \Delta b(\mathbf{r}) e^{i\mathbf{Q}\cdot\mathbf{r}} dV$

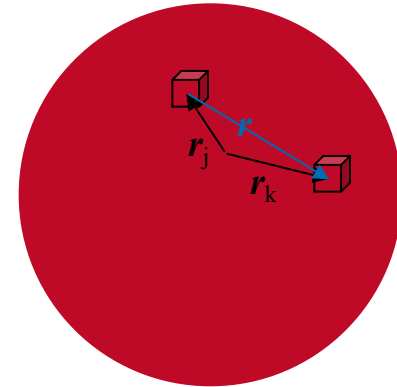
⋮

$$|\psi_{\text{p}}^{\text{sc}}(\mathbf{Q})|^2 = 4\pi \int_0^\infty p(r) \frac{\sin(qr)}{qr} dr$$



pair distance distribution function

$$p(r) = r^2 \int_V \Delta b(\mathbf{r}') \Delta b(\mathbf{r}-\mathbf{r}') d^3\mathbf{r}$$



any shape – Guinier approximation

for small Q ($Q < R_g^{-1}$):

$$P(Q) = 1 - \frac{1}{3}(QR_g)^2$$

with radius of gyration R_g

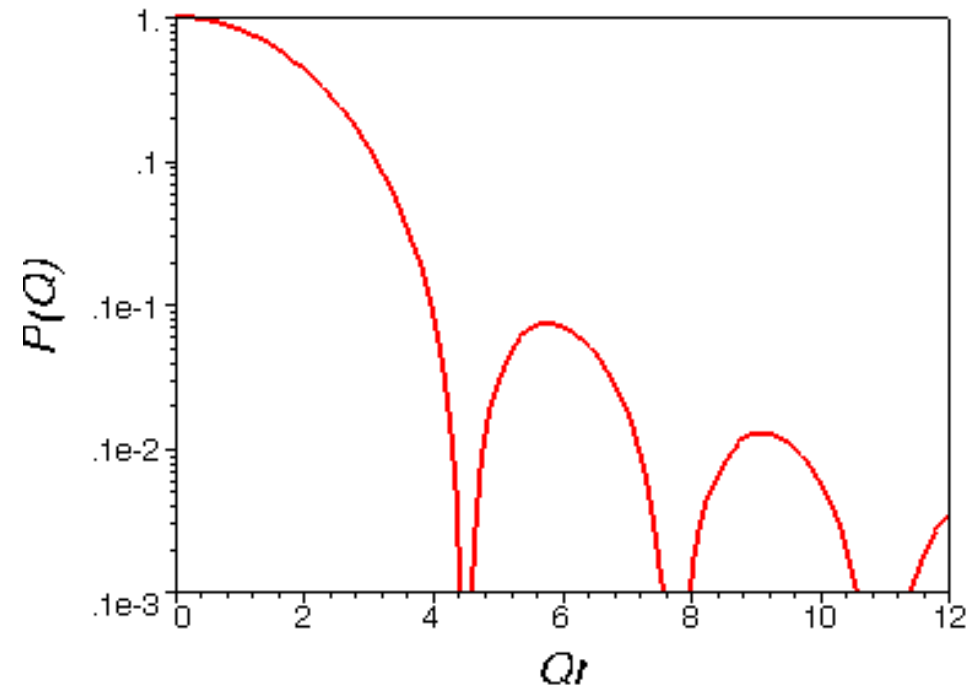
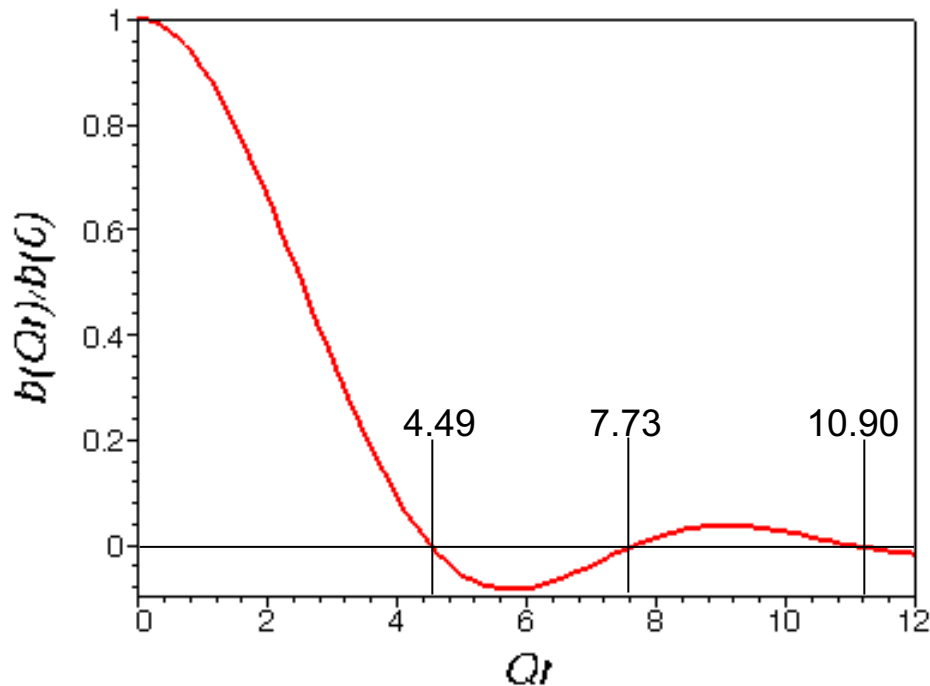
$$R_g^2 = \int_V \rho(\mathbf{r}) (\mathbf{r} - \mathbf{r}_{\text{cm}})^2 d\mathbf{r}$$

homogenous sphere

$$\psi_p^{sc}(\mathbf{Q}) = \int \Delta b(\mathbf{r}) e^{i\mathbf{Q}\cdot\mathbf{r}} dV = \Delta b \int_{V_{\text{sphere}}} e^{i\mathbf{Q}\cdot\mathbf{r}} dV = \Delta b \int d\phi \int r^2 dr \int e^{iQr \cos \theta} \sin \theta d\theta$$

$$\vdots$$

$$\sim \frac{3}{(QR)^3} (\sin(QR) - QR \cos(QR))$$



polydisperse homogenous spheres

radius of gyration

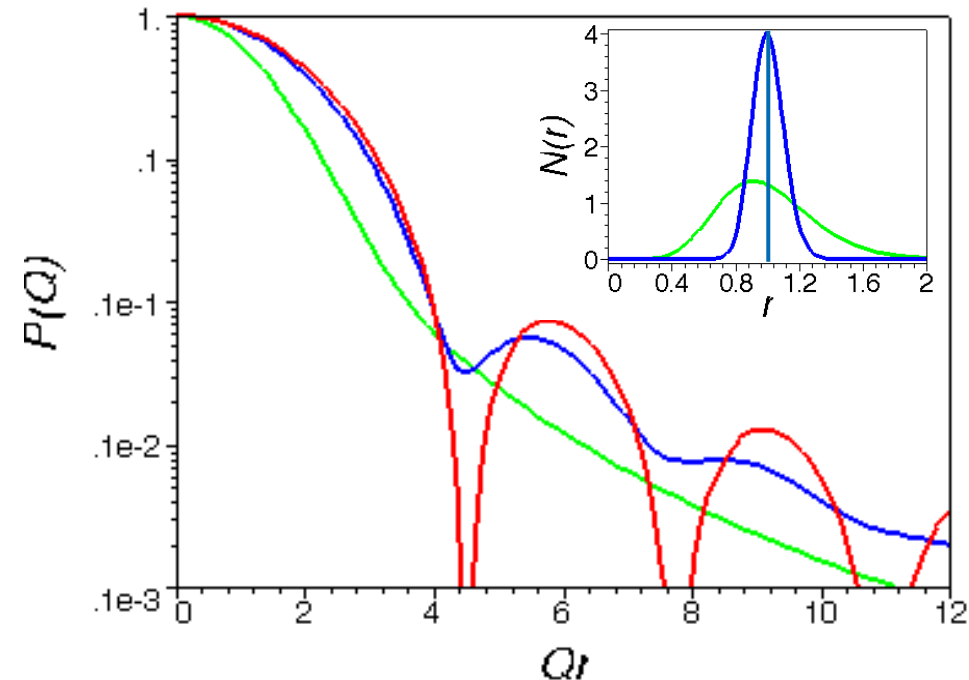
$$\langle R_g^2 \rangle = \frac{\int_0^\infty N(r) M^2(r) R_g^2 dr}{\int_0^\infty N(r) M^2(r) dr}$$

molar mass

$$\langle M \rangle = \frac{\int_0^\infty N(r) M^2(r) dr}{\int_0^\infty N(r) M(r) dr}$$

form factor

$$\langle P(\mathbf{q}) \rangle = \frac{\int_0^\infty N(r) M^2(r) P(\mathbf{q}, r) dr}{\int_0^\infty N(r) M^2(r) dr}$$



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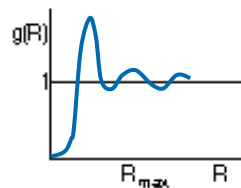
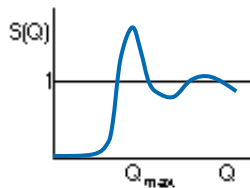
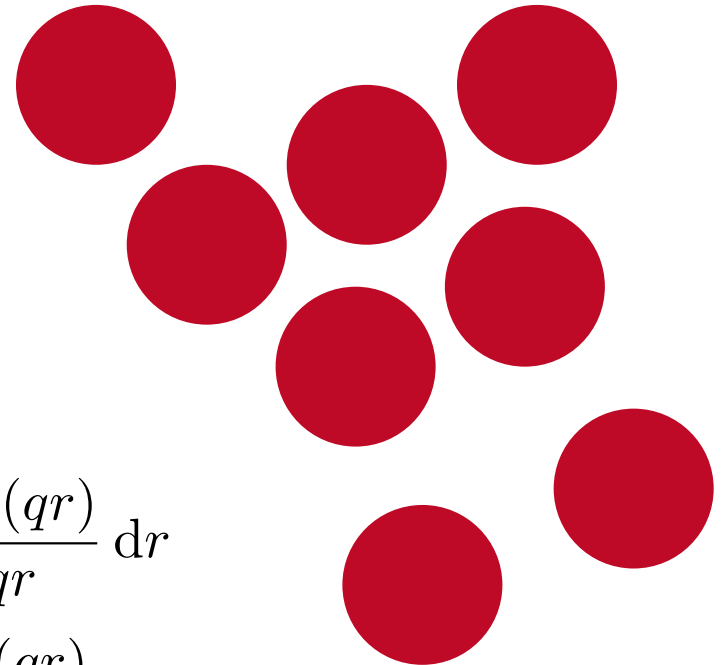
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interparticle interference
depends on particle arrangement

$$S(\mathbf{Q}) = \frac{1}{N} \sum_j \sum_k \langle e^{i\mathbf{Q} \cdot (\mathbf{R}_j - \mathbf{R}_k)} \rangle$$

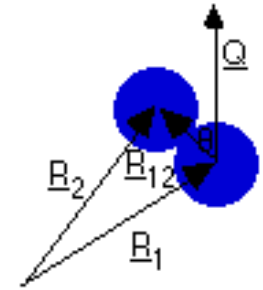
$$S(\mathbf{Q}) = 1 + 4\pi \frac{N}{V} \int (g(\mathbf{r}) - 1) r^2 \frac{\sin(qr)}{qr} dr$$

$$S(\mathbf{Q}) - 1 = 4\pi \frac{N}{V} \int (g(\mathbf{r}) - 1) r^2 \frac{\sin(qr)}{qr} dr$$



dumbbell

$$\begin{aligned} S(\mathbf{Q}) &= \frac{1}{N} \sum_j \sum_k \langle e^{i\mathbf{Q} \cdot (\mathbf{R}_j - \mathbf{R}_k)} \rangle = \frac{1}{2} \sum_{j=1}^2 \sum_{k=1}^2 \langle e^{i\mathbf{Q} \cdot (\mathbf{R}_j - \mathbf{R}_k)} \rangle \\ &= \frac{1}{2} \langle 1 + e^{i\mathbf{Q} \cdot (\mathbf{R}_1 - \mathbf{R}_2)} + e^{i\mathbf{Q} \cdot (\mathbf{R}_2 - \mathbf{R}_1)} + 1 \rangle \\ &= \langle 1 + \cos(\mathbf{Q} \cdot \mathbf{R}_{12}) \rangle \end{aligned}$$



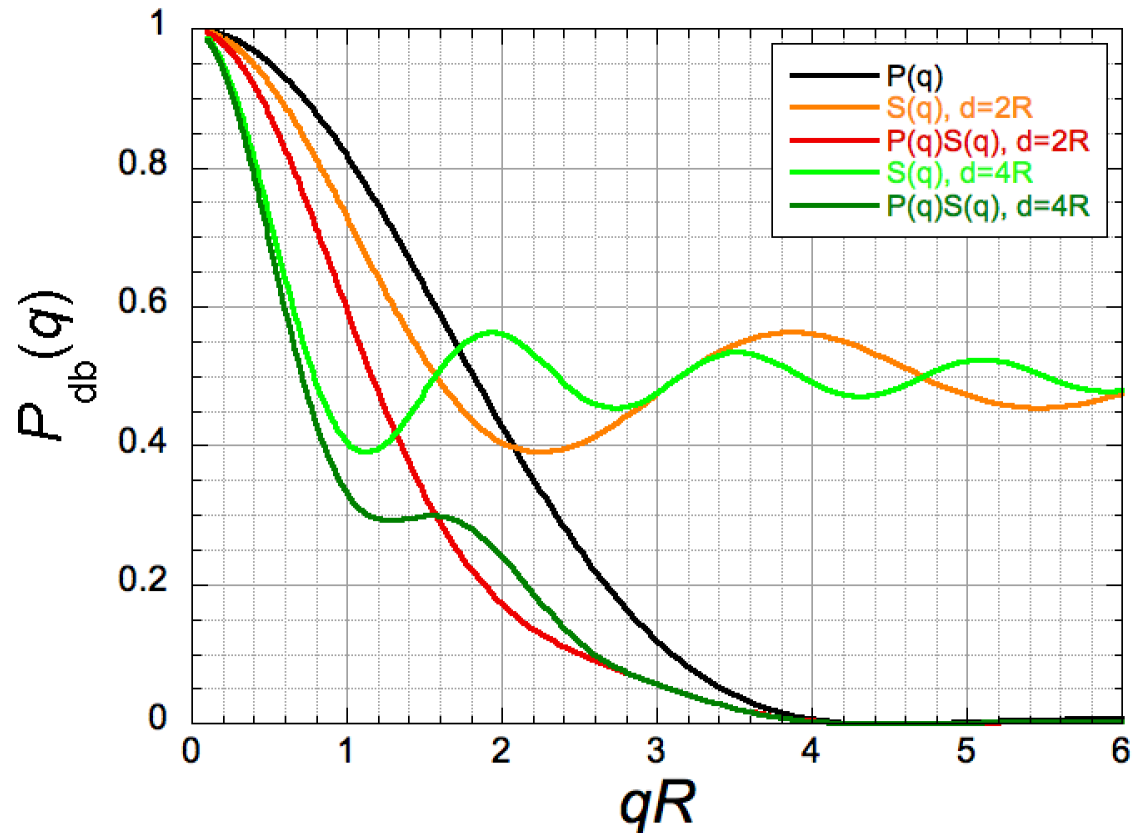
↓ spherical averaging

$$\begin{aligned} S(\mathbf{Q}) &= 1 + \frac{1}{4\pi} \int_0^{2\pi} \int_0^\pi \cos(\mathbf{Q} \cdot \mathbf{R}_{12}) \sin \theta d\theta d\phi \\ &= 1 + \frac{\sin(2QR_{12})}{2QR_{12}} \end{aligned}$$

↓ $P(q)$ of a sphere
dumbbell = 2 spheres

.....

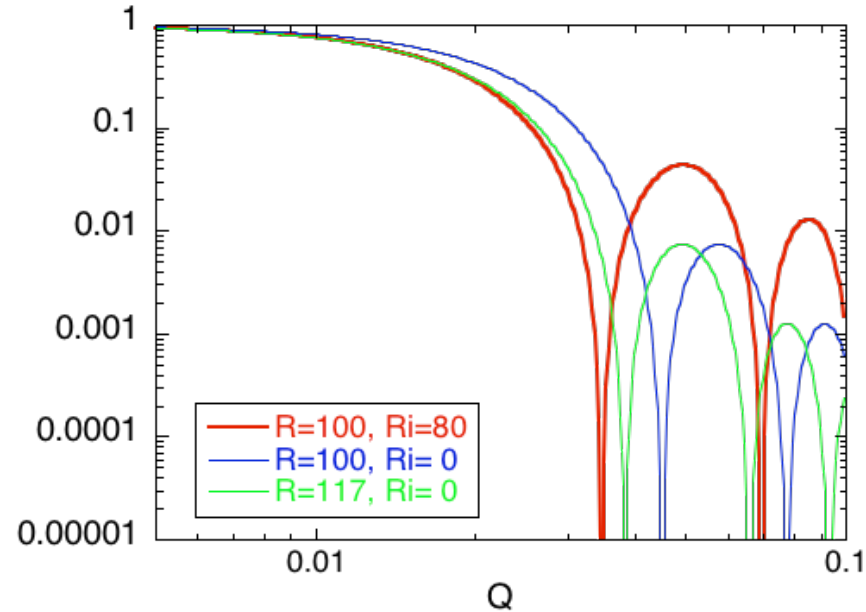
dumbbell



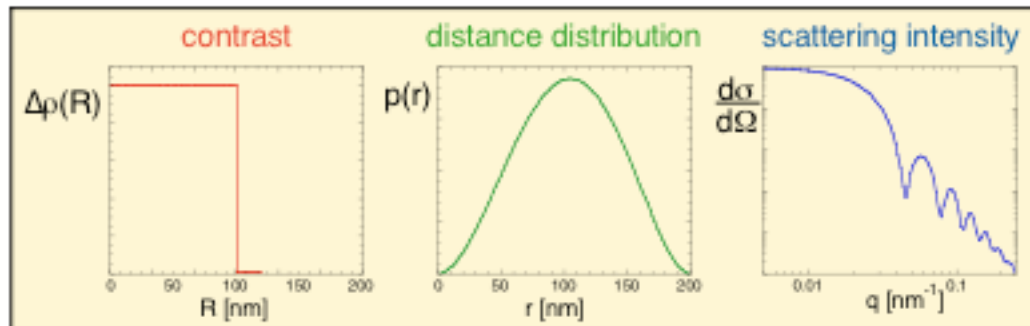
inhomogenous sphere (core-shell)



$$P(q) = \left(\frac{3}{(QR_s)^3 - (QR_c)^3} \{ (\sin(QR_s) - QR_s \cos(QR_s)) - (\sin(QR_c) - QR_c \cos(QR_c)) \} \right)^2$$



indirect Fourier transformation



model fitting (analytical, numerical, simulations)

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$$\frac{d\sigma}{d\Omega}(\mathbf{Q}) \sim N|\psi_p^{\text{sc}}(0)|^2$$

$$\psi_p^{\text{sc}}(\mathbf{Q}) = \int \Delta b(\mathbf{r}) e^{i\mathbf{Q}\cdot\mathbf{r}} dV \quad \rightarrow \quad \psi_p^{\text{sc}}(0) \sim \Delta b V \quad \text{with } \Delta b = \langle \Delta b(\mathbf{r}) \rangle \\ \sim (b_{\text{part}} - b_{\text{medium}}) V$$

$$\frac{d\sigma}{d\Omega}(0) \sim (b_{\text{part}} - b_{\text{medium}})^2$$

consider that solvent consists of a mixture of solvent A (b_A) and solvent B (b_B) with volume fraction ϕ_B of solvent B:

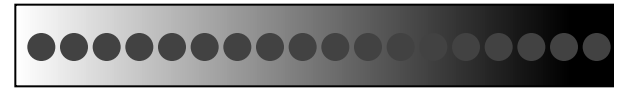
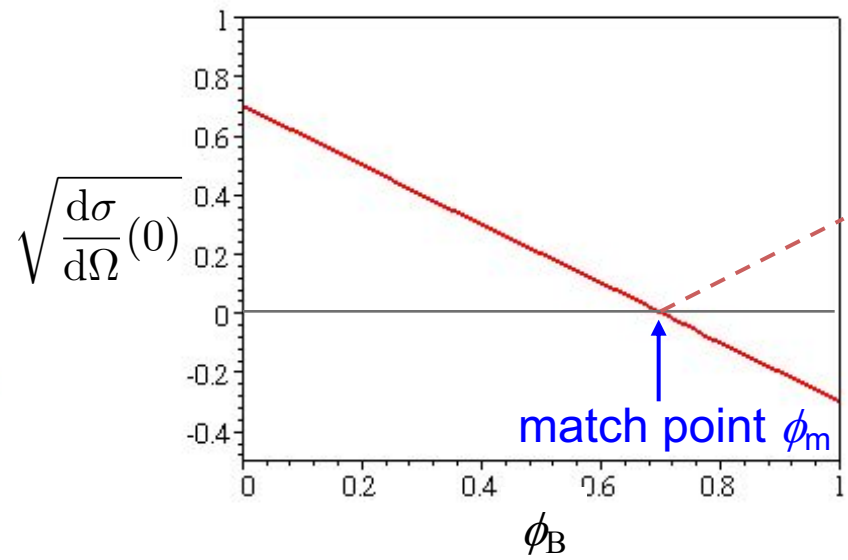
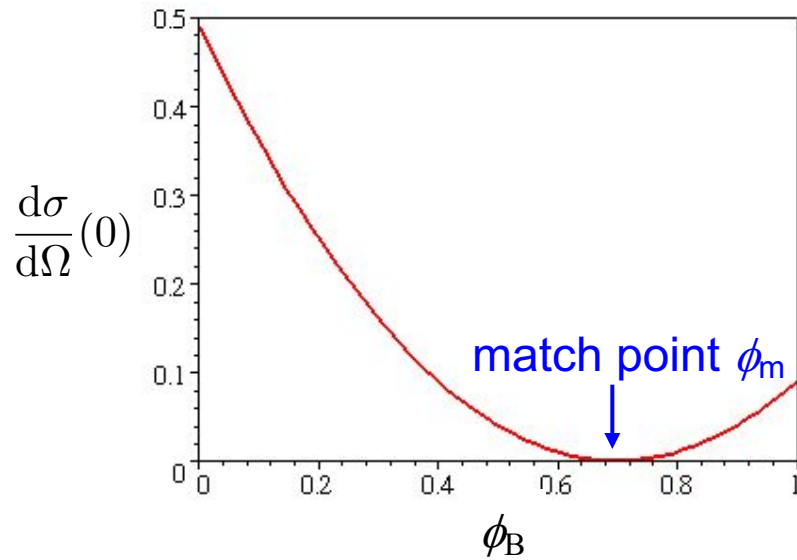
$$b_{\text{medium}} = (1 - \phi_B)b_A + \phi_B b_B = b_A + \phi_B(b_B - b_A)$$



$$\frac{d\sigma}{d\Omega} \sim (b_{\text{part}} - b_{\text{medium}})^2 = (b_{\text{part}} - \{b_A + \phi_B(b_B - b_A)\})^2$$

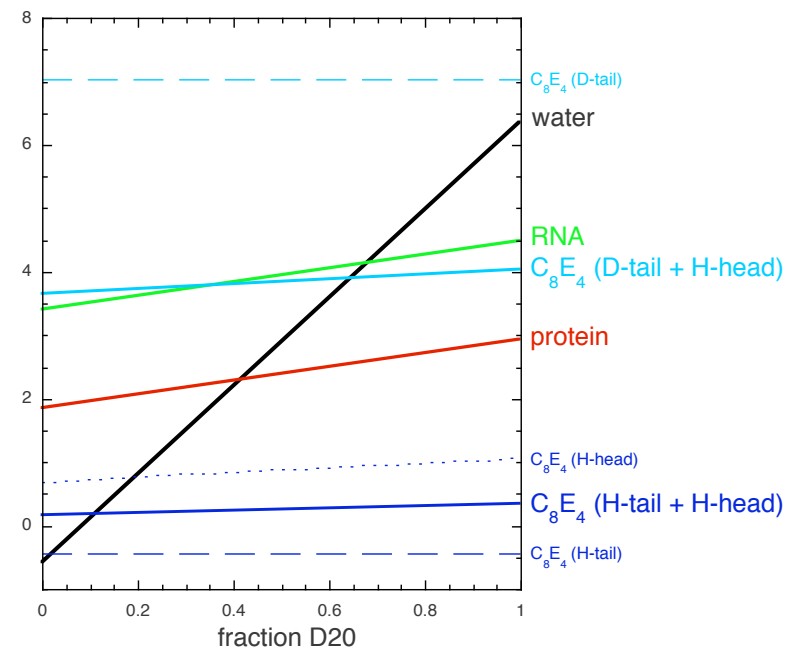
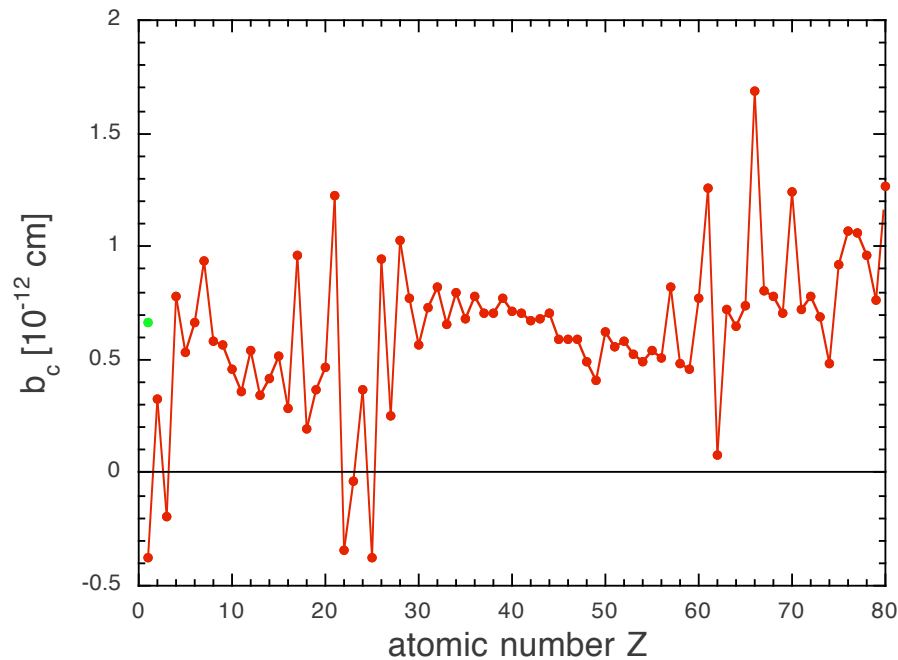
$$\sim (b_B - b_A)^2 (\phi_m - \phi_B)^2, \quad \phi_m = \frac{b_{\text{part}} - b_A}{b_B - b_A}$$

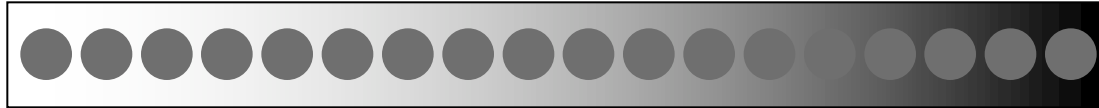
(Solvent) contrast variation



scattering length depends on:

- nucleus, i.e. isotope (non-systematic)
- nuclear spin (magnetic scattering)





- increase signal
improve statistics
- reduce signal
reduce multiple scattering
- reduce 'noise'
reduce background due to solvent

example: dilute polystyrene spheres in cyclohexane
with $R = 10$ nm, $c = 10$ mg/ml ($\approx 1\%$)

substance	formula	molar mass M (g/mol)	density \bar{v}^{-1} (g/cm ³)	scatt. length dens. b (10 ¹⁰ cm ⁻²)
polystyrene	(C ₈ H ₈) _n	(104.2) _n	1.04	1.44
	(C ₈ D ₈) _n	(112.2) _n	~1.12	6.44
cyclohexane	C ₆ H ₁₂	84.2	0.779	- 0.24
	C ₆ D ₁₂	96.2	0.893	6.74

example: dilute polystyrene spheres in cyclohexane
with $R = 10$ nm, $c = 10$ mg/ml ($\approx 1\%$)

excess spheres ('signal')

H in **H**: 10.8 cm^{-1}

H in **D**: 113.1 cm^{-1}

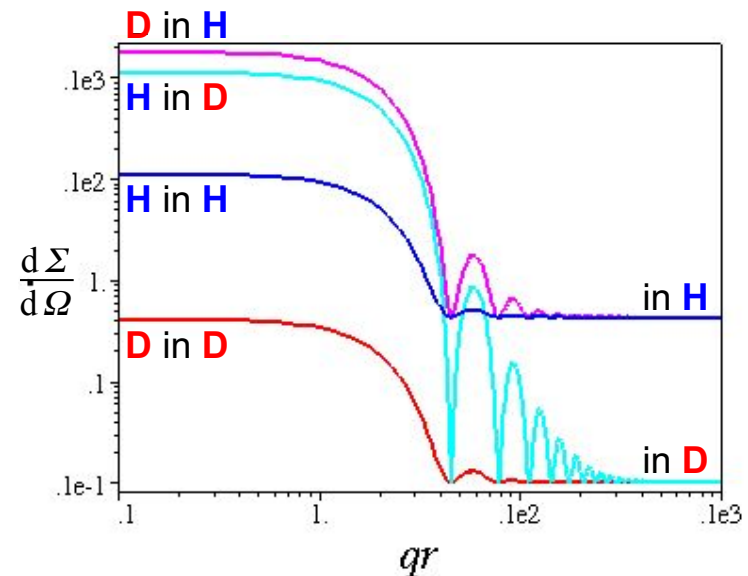
D in **H**: 179.7 cm^{-1}

D in **D**: 0.4 cm^{-1}

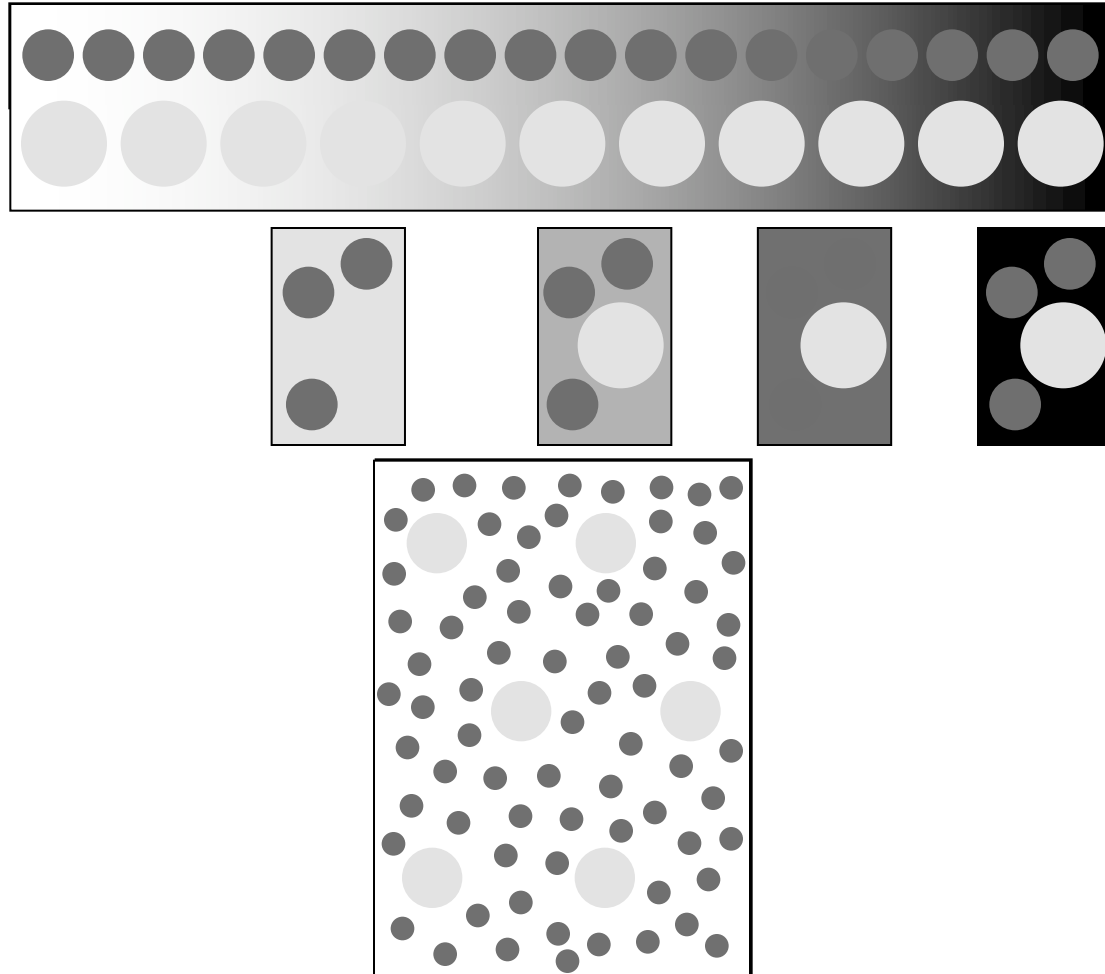
solvent ('noise')

in **H**: 0.43 cm^{-1}

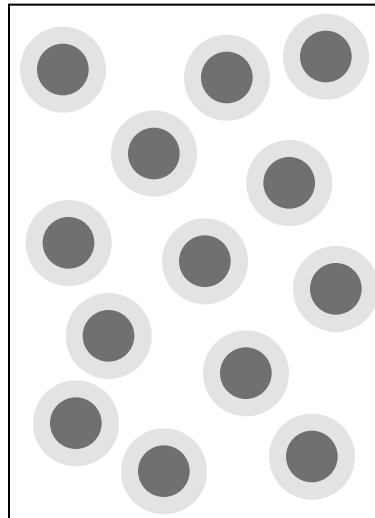
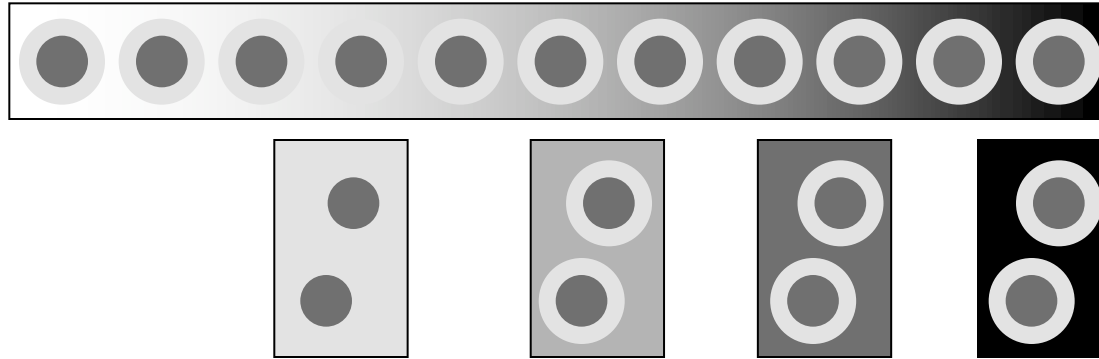
in **D**: 0.01 cm^{-1}



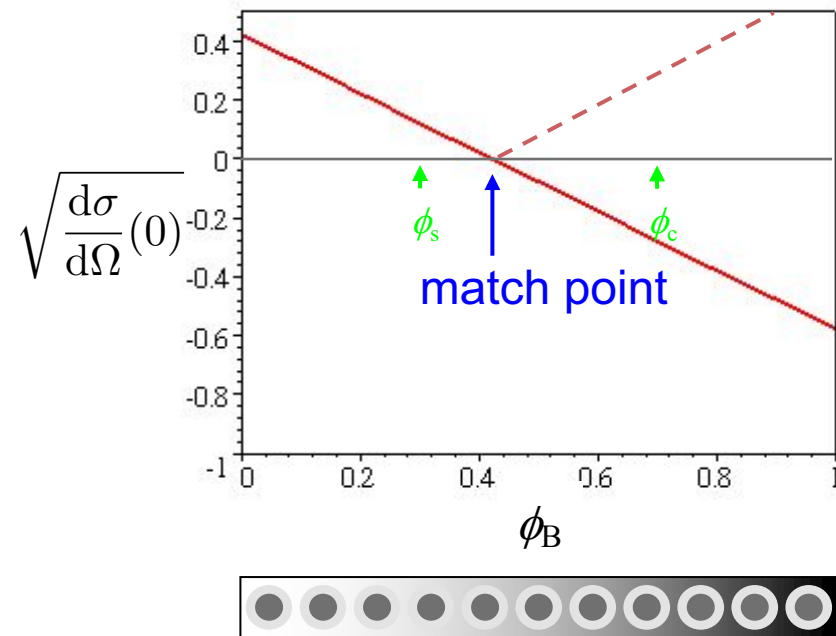
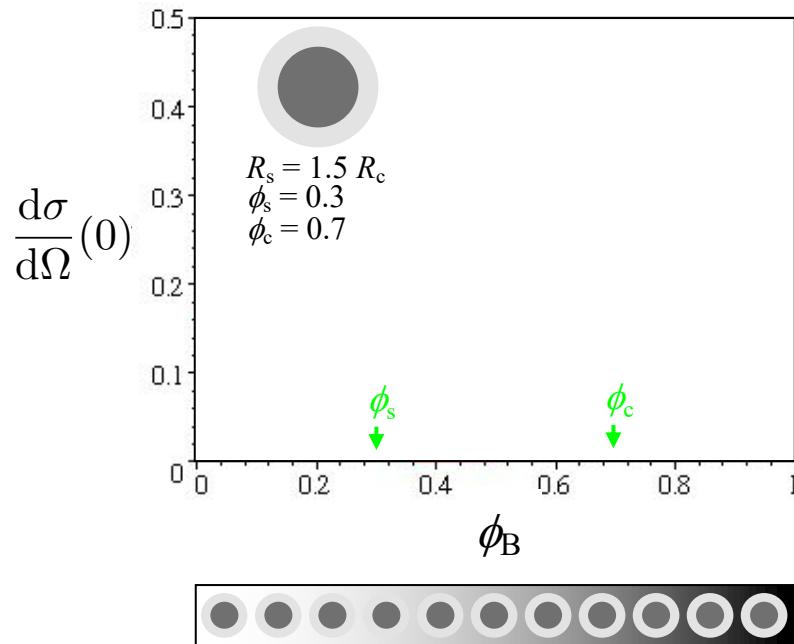
Heterogeneous samples



Heterogeneous particles



(Solvent) contrast variation

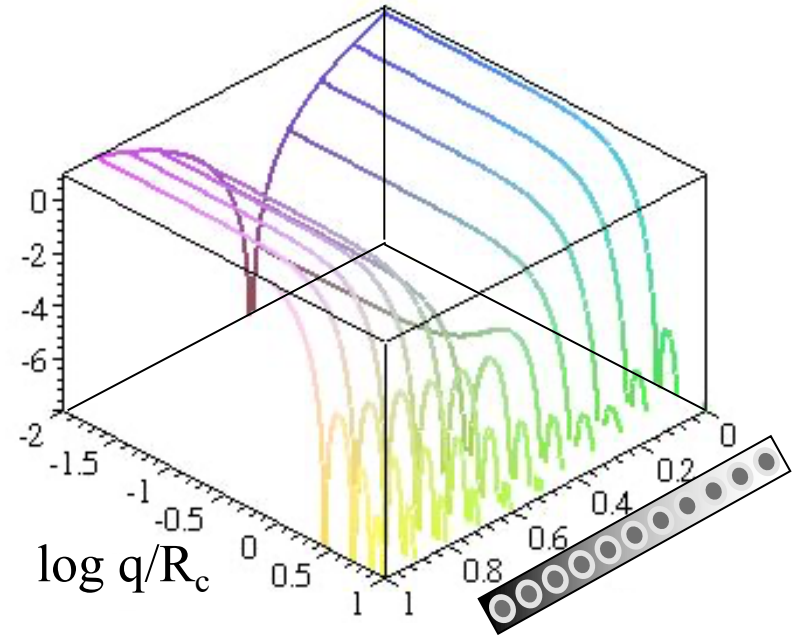
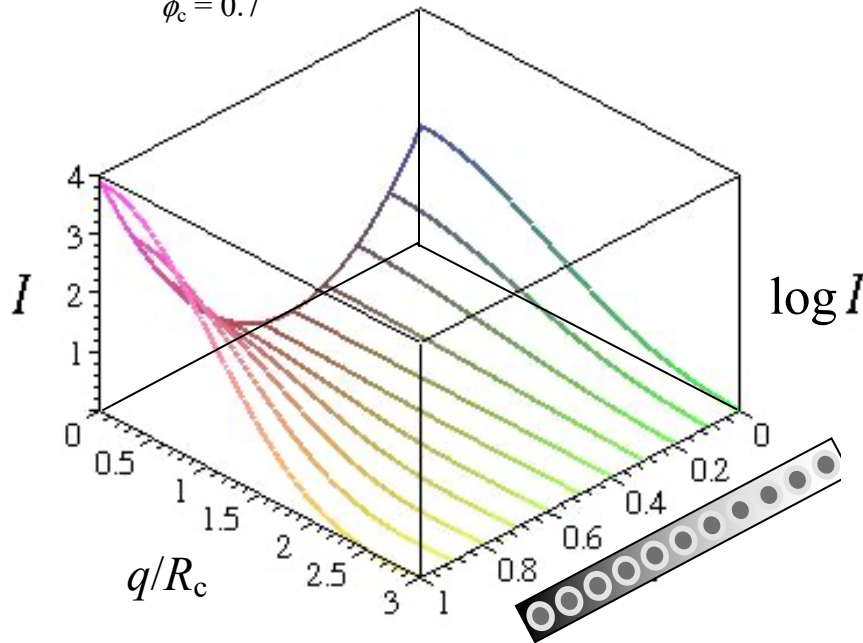


$$\phi_{part} = \frac{\phi_s(V_s - V_c) + \phi_c V_c}{V_s} = \frac{0.3 \times (1.5^3 - 1^3) + 0.7 \times 1^3}{1.5^3} = 0.42$$

(Solvent) contrast variation



$$R_s = 1.5 R_c$$
$$\phi_s = 0.3$$
$$\phi_c = 0.7$$



$$\frac{d\sigma}{d\Omega}(\mathbf{Q}) \sim |\psi^{\text{sc}}|^2 \sim \sum_k \sum_j \left\langle \psi_{p,j}^{\text{sc}}(\mathbf{Q}) \psi_{p,k}^{\text{sc}*}(\mathbf{Q}) e^{i\mathbf{Q}\cdot(\mathbf{R}_j - \mathbf{R}_k)} \right\rangle$$

↓ identical scatterers

$$\frac{d\sigma}{d\Omega}(\mathbf{Q}) \sim N |\psi_p^{\text{sc}}(0)|^2 \frac{|\psi_p^{\text{sc}}(\mathbf{Q})|^2}{|\psi_p^{\text{sc}}(0)|^2} \frac{1}{N} \sum_k \sum_j \left\langle e^{i\mathbf{Q}\cdot(\mathbf{R}_j(t) - \mathbf{R}_k(t))} \right\rangle$$

but b_j, b_k depend on nucleus (isotope) and spin

thus for same isotope

$$j \neq k : \langle b_j b_k e^{i\mathbf{Q}\cdot\mathbf{r}_{jk}} \rangle = \langle b_j \rangle \langle b_k \rangle \langle e^{i\mathbf{Q}\cdot\mathbf{r}_{jk}} \rangle = \langle b \rangle^2 \langle e^{i\mathbf{Q}\cdot\mathbf{r}_{jk}} \rangle$$

$$j = k : \langle b_j b_k e^{i\mathbf{Q}\cdot\mathbf{r}_{jk}} \rangle = \langle b_j^2 \rangle = \langle b^2 \rangle$$

↓

$$\frac{d\sigma}{d\Omega} \sim \sum_j \sum_k \langle b_j b_k e^{i\mathbf{Q}\cdot\mathbf{r}_{jk}} \rangle = \sum_j \sum_{k \neq j} \langle b_j b_k e^{i\mathbf{Q}\cdot\mathbf{r}_{jk}} \rangle + \sum_j \langle b_j b_j e^{i\mathbf{Q}\cdot\mathbf{r}_{jk}} \rangle$$

$$\frac{d\sigma}{d\Omega} = \langle b \rangle^2 \sum_j \sum_{k \neq j} \langle e^{i\mathbf{Q}\cdot\mathbf{r}_{jk}} \rangle + N \langle b^2 \rangle$$

$$\frac{d\sigma}{d\Omega} = \langle b \rangle^2 \sum_j \sum_k \langle e^{i\mathbf{Q}\cdot\mathbf{r}_{jk}} \rangle + N \{ \langle b^2 \rangle - \langle b \rangle^2 \}$$

coherent cross section

$$\sigma_{coh} = 4\pi \langle b \rangle^2$$

- contains structural information on particles & arrangement ($e^{i\mathbf{Q}r}$ term)
- collective properties of particles ($S(\mathbf{Q}), D_c$)

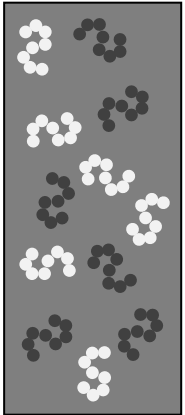
incoherent cross section

$$\sigma_{incoh} = 4\pi \left(\langle b^2 \rangle - \langle b \rangle^2 \right)$$

- contains no structural information on arrangement (no interference term, missing phase relation)
- properties of individual particles (D_s)

$$\frac{d\sigma}{d\Omega} = \langle b \rangle^2 \sum_j \sum_k \langle e^{i\mathbf{Q} \cdot \mathbf{r}_{jk}} \rangle$$

$$+ N \{ \langle b^2 \rangle - \langle b \rangle^2 \}$$



identical particles, but hydrogenated (b_H) and deuterated (b_D) species in H/D-solvent such that $\Delta b_H = -\Delta b_D$

$$\Rightarrow \langle b \rangle = \frac{1}{2}(\Delta b_H + \Delta b_D) = 0$$

\Rightarrow coherent contribution = 0
(i.e. structural contribution)

$$\langle b^2 \rangle = \frac{1}{2}(\Delta b_H^2 + \Delta b_D^2) = \Delta b_H^2 = \Delta b_D^2$$

incoherent contribution $\neq 0$
(i.e. individual particle contribution)

coherent cross section

$$\sigma_{coh} = 4\pi \langle b \rangle^2$$

- contains structural information on particles & arrangement ($e^{i\mathbf{Q}\cdot\mathbf{r}}$ term)
- collective properties of particles ($S(\mathbf{Q}), D_c$)

incoherent cross section

$$\sigma_{incoh} = 4\pi (\langle b^2 \rangle - \langle b \rangle^2)$$

- contains no structural information on arrangement (no interference term, missing phase relation)
- properties of individual particles (D_s)

$$\frac{d\sigma}{d\Omega} = \langle b \rangle^2 \sum_j \sum_k \langle e^{i\mathbf{Q}\cdot\mathbf{r}_{jk}} \rangle$$

$$+ N \{ \langle b^2 \rangle - \langle b \rangle^2 \}$$

solvent isotope substitution can cause

- shift in cmc (generally higher in H₂O)
- shift in Θ temperature ($\pm 4^\circ\text{C}$ for PS in cyclohexane) or critical point
- shift in melting point (6°C for polyethylene)
- selective adsorption
- exchange of particle H/D with solvent (e.g. OH, NH, COOH groups, scattering length density becomes a function of the solvent H/D ratio)
- ...

'particle' isotope substitution can cause

- phase separation (e.g. H- and D-polymer)
- ...

non-crystalline materials are wonderful

K. Paul, V. Thomas, *Oxford University Press* (1987)

stefan.egelhaaf@uni-duesseldorf.de