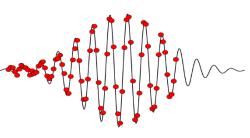
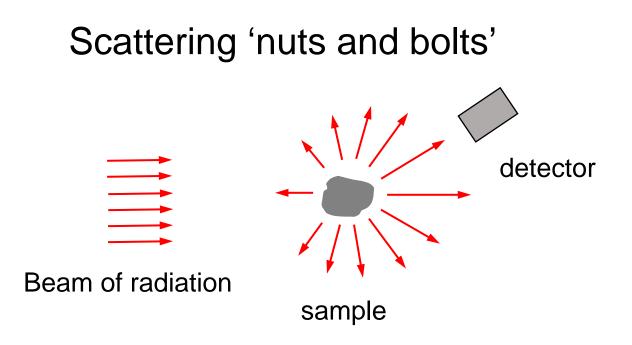


## The neutron: properties and interactions

Andrew Boothroyd Oxford University

- What is a neutron?
- How does it interact with matter?
- Prepare the ground for understanding the principles of neutron scattering





- Source of radiation: photons, neutrons, electrons, atoms, ...
- Measure distribution of radiation scattered from a sample
- Form of interaction determines what you can measure
- Radiation must be coherent to measure correlations

### Discovery of the neutron

1932 (Chadwick)

NATURE

[FEBRUARY 27, 1932

#### Letters to the Editor

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#### Possible Existence of a Neutron

It has been shown by Bothe and others that beryllium when bombarded by  $\alpha$ -particles of polonium emits a radiation of great penetrating power, which has an absorption coefficient in lead of about 0.3 (cm.)<sup>-1</sup>. Recently Mme. Curie-Joliot and M. Joliot found, when measuring the ionisation produced by this beryllium radiation in a vessel with a thin window, that the ionisation increased when matter containing hydrogen was placed in front of the window. The This again receives a simple explanation on the neutron hypothesis.

If it be supposed that the radiation consists of quanta, then the capture of the *a*-particle by the Be<sup>9</sup> nucleus will form a C<sup>13</sup> nucleus. The mass defect of C<sup>13</sup> is known with sufficient accuracy to show that the energy of the quantum emitted in this process cannot be greater than about  $14 \times 10^6$  volts. It is difficult to make such a quantum responsible for the effects observed.

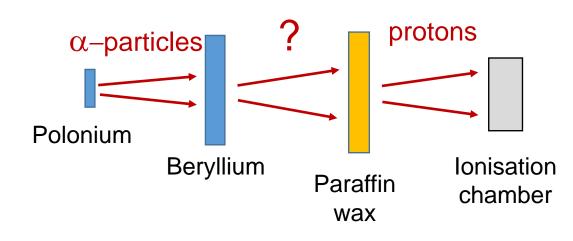
It is to be expected that many of the effects of a neutron in passing through matter should resemble those of a quantum of high energy, and it is not easy to reach the final decision between the two hypotheses. Up to the present, all the evidence is in favour of the neutron, while the quantum hypothesis can only be upheld if the conservation of energy and momentum be relinquished at some point.

J. CHADWICK.

Cavendish Laboratory,

# Discovery of the neutron

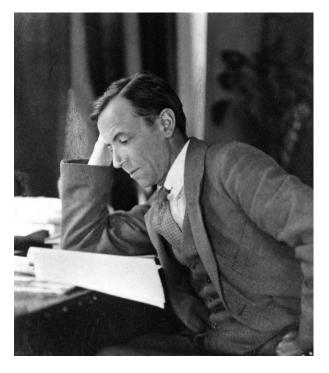
I. Curie and F. Joliot, 1932



Chadwick, 1932

? = neutrons

 ${}^{9}\text{Be} + \alpha = {}^{12}\text{C} + \text{n}$ 



James Chadwick (1891-1974) Nobel Prize in Physics (1935)

### Neutrons are particles and waves!

• matter wave:

- Kinematics (de Broglie):
  - momentum:  $p = h/\lambda$  (h = Planck's constant)  $\mathbf{p} = \hbar \mathbf{k}$  ( $\mathbf{k}$  = wavevector,  $k = 2\pi/\lambda$ ) ( $\hbar = h/2\pi$ )

kinetic energy:  $E = \frac{1}{2}m_{\mathrm{n}}v^2 = p^2/2m_{\mathrm{n}}$ 

$$\rightarrow E = h^2/(2m_{\rm n}\lambda^2) = \hbar^2 k^2/2m_{\rm n}$$

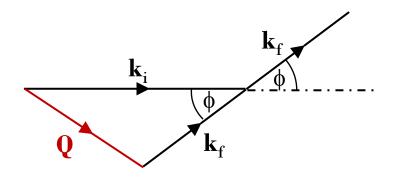
## Scattering kinematics

momentum transfer

 $\hbar \mathbf{Q} = \hbar \mathbf{k}_{\rm i} - \hbar \mathbf{k}_{\rm f}$ 

• energy transfer

 $\hbar \omega = E_{\rm i} - E_{\rm f}$ =  $\hbar^2 (k_{\rm i}^2 - k_{\rm f}^2)/2m_{\rm n}$ 

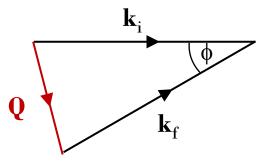


scattering triangle

• a scattering event is characterised by ( $\mathbf{Q}, \omega$ )

Elastic scattering (diffraction):  $\hbar \omega = 0$ Inelastic scattering (spectroscopy):  $\hbar \omega \neq 0$  Elastic scattering (diffraction):  $\hbar \omega = 0$ 

• Scattering vector  $\mathbf{Q} = \mathbf{k}_{i} - \mathbf{k}_{f}$ 



BUT sometimes written

$$\mathbf{Q} = \mathbf{k}_{\mathrm{f}} - \mathbf{k}_{\mathrm{i}}$$

isosceles scattering triangle

## Beware of crystallographers!

- Mass =  $1.675 \times 10^{-27} \text{ kg} (= 1.001 m_{\text{p}})$
- Charge = 0
- Mean lifetime ≈ 15 min
- Spin =  $\frac{1}{2}$

• Magnetic moment = 1.91  $\mu_N$  (~0.001 $\mu_B$ )

- Mass =  $1.675 \times 10^{-27} \text{ kg} (= 1.001 m_{\text{p}})$
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• Magnetic moment = 1.91  $\mu_N$  (~0.001 $\mu_B$ )

• Mass =  $1.675 \times 10^{-27} \text{ kg} (= 1.001 m_{\text{p}})$ 

kinetic energy:  $E = \frac{1}{2}m_{\rm n}v^2$ 

$$= p^2/2m_n$$
 ( $p$  = momentum)

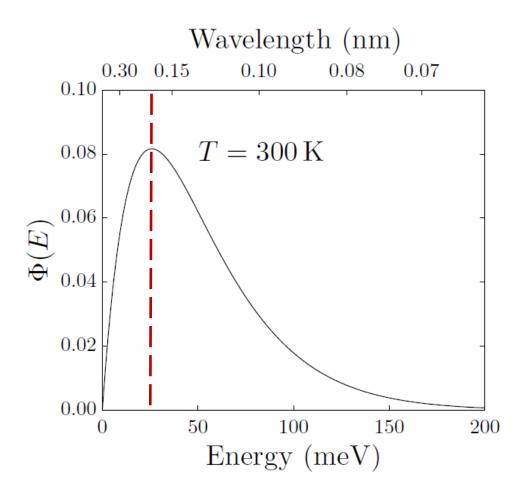
de Broglie:  $p = h/\lambda$  (h = Planck's constant)

 $\rightarrow E = h^2/(2m_{\rm n}\lambda^2)$ 

hence  $E = 25 \text{ meV} \rightarrow v = 2,200 \text{ ms}^{-1}$  and  $\lambda = 0.18 \text{ nm}$ 

 $E \sim$  energy of thermal excitations in condensed matter  $\lambda \sim$  inter-atomic distances in condensed matter

### **Thermal neutrons**



neutrons are well matched to probe the atomic-scale structure and dynamics of substances at normal operating temperatures

### Comparison with x-rays

X-rays are electromagnetic radiation (photons)

Einstein: 
$$E = hf$$
 ( $f =$  frequency)  
=  $hc/\lambda$  ( $c = f\lambda =$  speed of light  
=  $3.0 \ge 10^8 \text{ ms}^{-1}$ )

hence  $E = 6.9 \,\mathrm{keV} \longrightarrow \lambda = 0.18 \,\mathrm{nm}$  (x-rays)

 $E = 25 \text{ meV} \rightarrow \lambda = 0.18 \text{ nm}$  (neutrons)

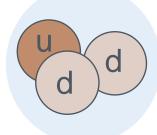
Neutrons better suited to low energy dynamics

- Mass =  $1.675 \times 10^{-27} \text{ kg} (= 1.001 m_{\text{p}})$
- Charge = 0
- Mean lifetime ≈ 15 min
- Spin =  $\frac{1}{2}$

• Magnetic moment = 1.91  $\mu_N$  (~0.001 $\mu_B$ )

Standard model: neutron composed of 3 quarks

Charge = 0



U:  $+\frac{2}{3}e$ d:  $-\frac{1}{3}e$ 

Electric dipole moment:  $d < 4.6 \times 10^{-47}$  C m (water molecule:  $d = 6.1 \times 10^{-30}$  C m)

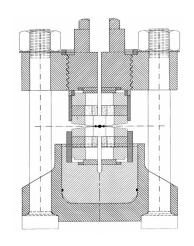
Neutrons do not feel electrostatic forces

 $\rightarrow$  neutrons are highly penetrating

Neutron penetration depth ~ 1 cm

- bulk probe
- → surface effects can usually be neglected (exception is neutron reflectometry)
- non-destructive
- can study samples in complex sample environments, e.g. cryostats, magnets, pressure cells, ...





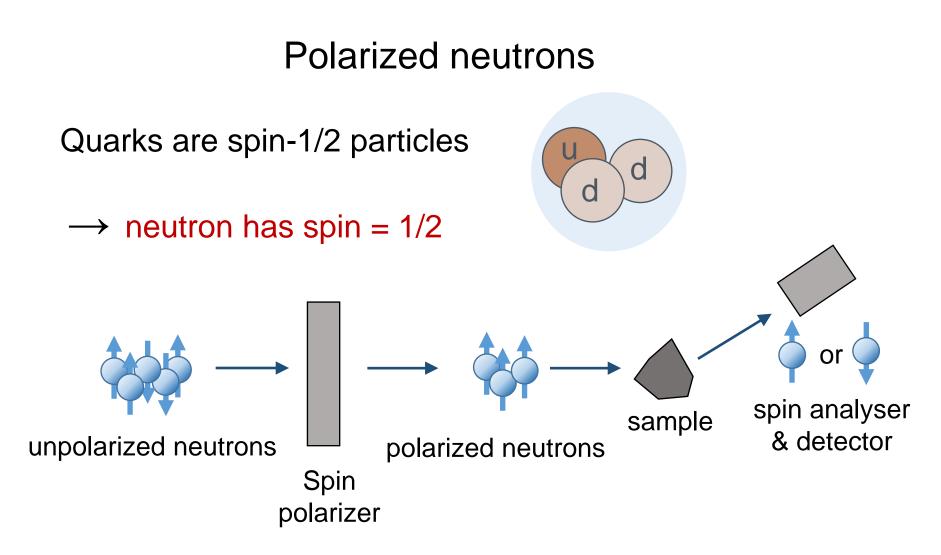


- Mass =  $1.675 \times 10^{-27} \text{ kg} (= 1.001 m_{\text{p}})$
- Charge = 0
- Mean lifetime ≈ 15 min
- Spin =  $\frac{1}{2}$
- Magnetic moment = 1.91  $\mu_N$  (~0.001 $\mu_B$ )

Neutrons last long enough to perform experiments

- Mass =  $1.675 \times 10^{-27} \text{ kg} (= 1.001 m_{\text{p}})$
- Charge = 0
- Mean lifetime ≈ 15 min
- Spin =  $\frac{1}{2}$
- Magnetic moment = 1.91  $\mu_N$  (~0.001 $\mu_B$ )

#### Neutrons can be spin-polarized



measure polarization-dependent scattering intensities

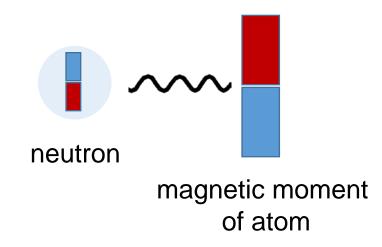
 $\rightarrow$  neutron polarization analysis

- Mass =  $1.675 \times 10^{-27} \text{ kg} (= 1.001 m_{\text{p}})$
- Charge = 0
- Mean lifetime ≈ 15 min
- Spin = 1/2

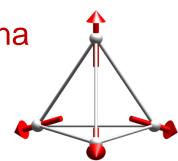
• Magnetic moment = 1.91  $\mu_N$  (~0.001 $\mu_B$ )

### → Magnetic neutron scattering

### Magnetic scattering



Can study a wide range of magnetic phenomena e.g. magnetic structures



Neutron interactions with matter

Neutrons interact with

Atomic nuclei (strong nuclear force — short-range)
Magnetic fields from unpaired electrons
In both cases, the interaction is very weak
Weaknesses:

Strengths:

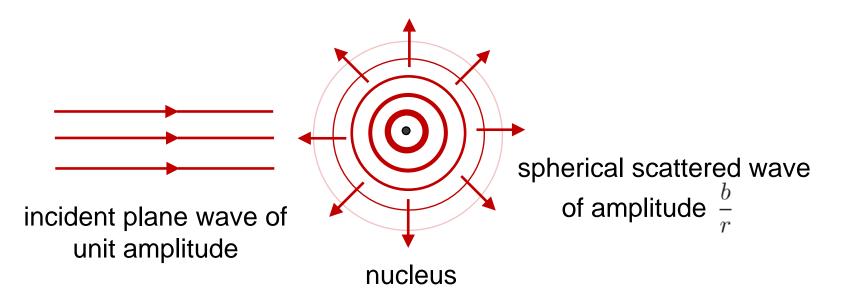
- 1. Neutrons probe the bulk (~1 cm)
- 2. Neutrons are non-destructive
- 3. Intensity can be calibrated
- 4. Theory is quantitative

Sample size is important:

single crystals: ~1 mm<sup>3</sup> (diffraction) ~1 cm<sup>3</sup> (spectroscopy)

powders/fluids: ~1 g (diffraction) ~10 g (spectroscopy)

## Scattering from bound nuclei



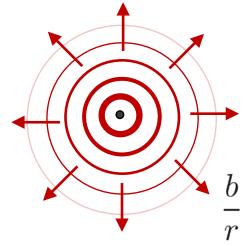


### Nuclear scattering length

Nucleus is very small ~  $10^{-15}$  m

Formal theory uses a pseudopotential:

$$V_{\rm N}(\mathbf{r}) = \frac{2\pi\hbar^2}{m_{\rm n}} \, b \, \delta(\mathbf{r})$$



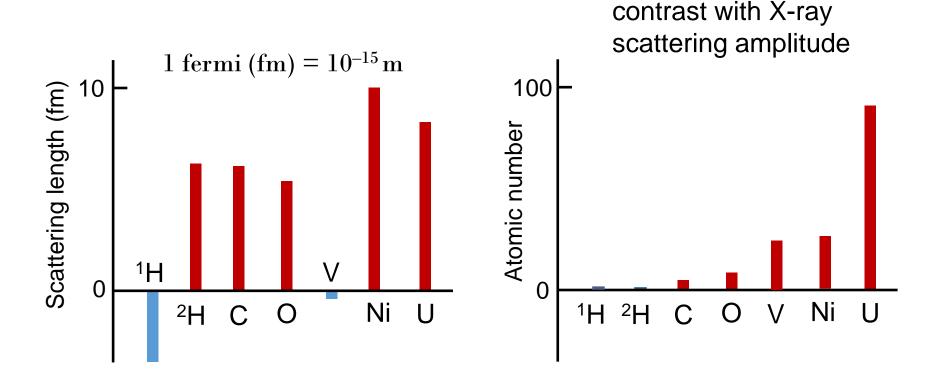
b = nuclear scattering length

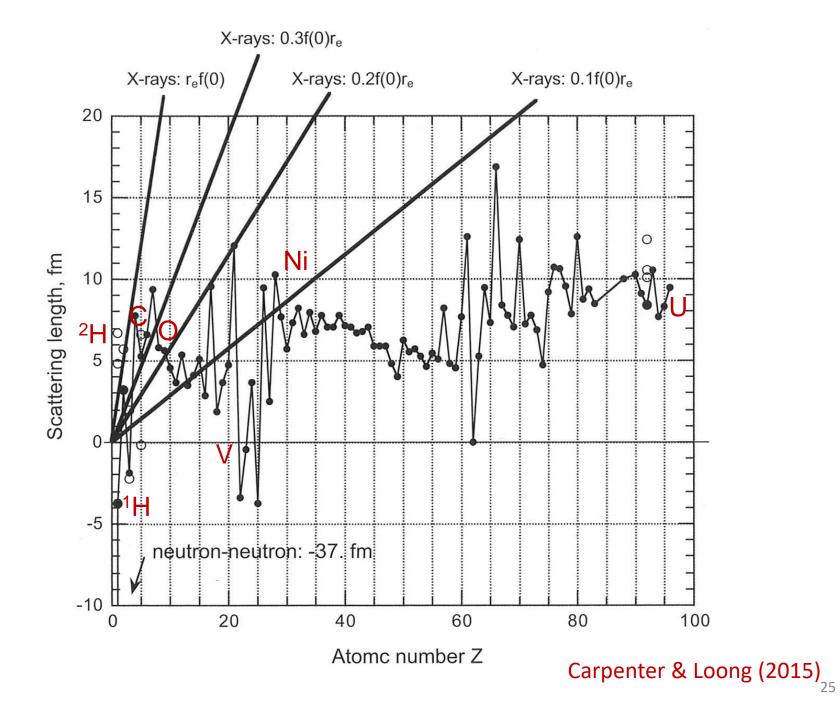
*b* measures both the amplitude of scattered wave and the strength of the nuclear potential

### Nuclear scattering length

**b** varies irregularly with atomic number

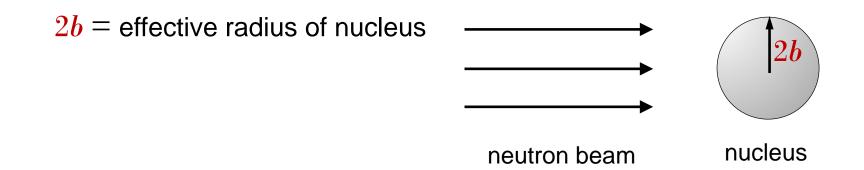
**b** also depends on isotope and on nuclear spin orientation





### **Cross-sections**

scattering amplitude ~ b [length] scattering intensity ~  $b^2$  [area]

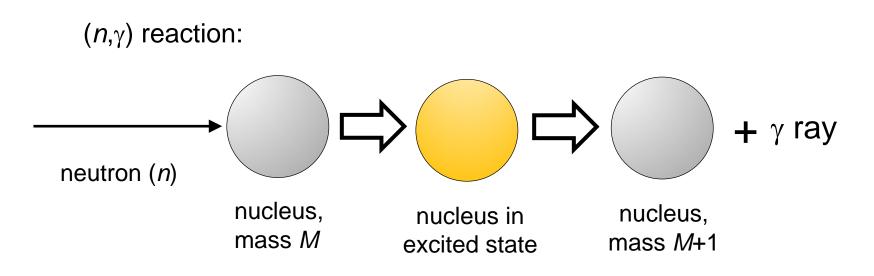


 $\sigma_s = 4\pi b^2$  = effective cross-sectional area of nucleus

Scattering probabilities are expressed as cross-sections

unit: 1 barn (b) =  $10^{-28}$  m<sup>2</sup>

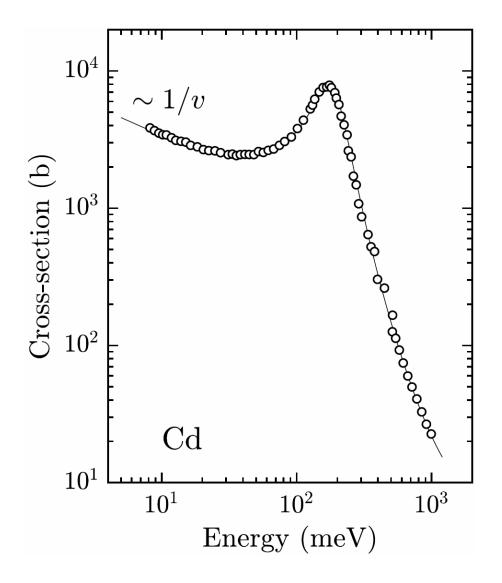
### Absorption cross-sections



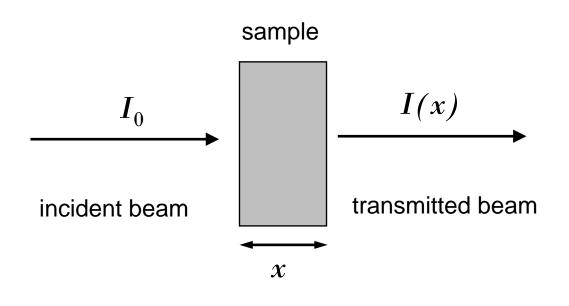
Strongly absorbing nuclei:

<sup>3</sup>He, <sup>10</sup>B, <sup>113</sup>Cd, <sup>149</sup>Sm, <sup>151</sup>Eu, <sup>155</sup>Gd, <sup>157</sup>Gd

Reaction probability is described by the absorption cross-section  $\sigma_a$ , which is dependent on neutron energy



### neutron transmission



• Beer-Lambert law:

$$I(x) = I_0 \exp(-n\sigma x)$$
  
transmission

n = number density  $\sigma = \sigma_{a} + \sigma_{s}$ 

# Comparison with x-rays

#### neutrons

- good for low energies
- bulk probe: mm to cm
- large samples required
- non-destructive
- good for light elements
- scattering is isotope-dependent
- can obtain cross-sections quantitatively
- theory is accurate and understood

#### X-rays

- good for high energies
- $\bullet$  probes nm to  $\mu m$
- can measure small samples
- X-rays can affect sample
- good for heavy elements
- scattering depends on atomic number
- difficult to be quantitative
- theory still being developed

### Strength of neutrons

It might well be said that, if the neutron did not exist, it would need to be invented!

(B. N. Brockhouse, 1983)