

# Neutrons in soft matter

Lecture 1 - Structure

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# Outline

## Lecture 1 – Structure & kinetics – SANS

### Introduction

soft matter & relevance of neutron scattering

Single objects: spheres, coils, rods...

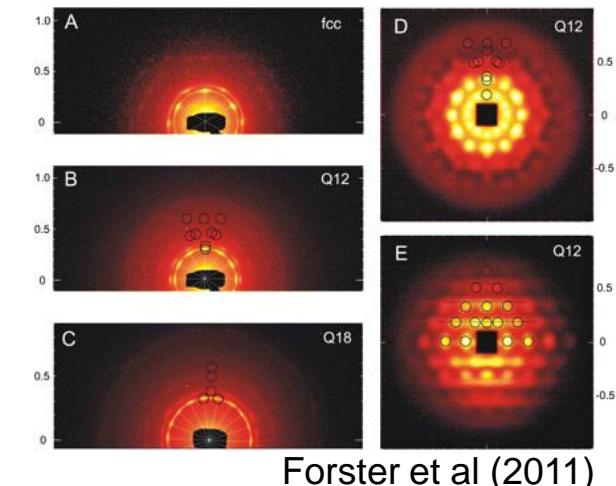
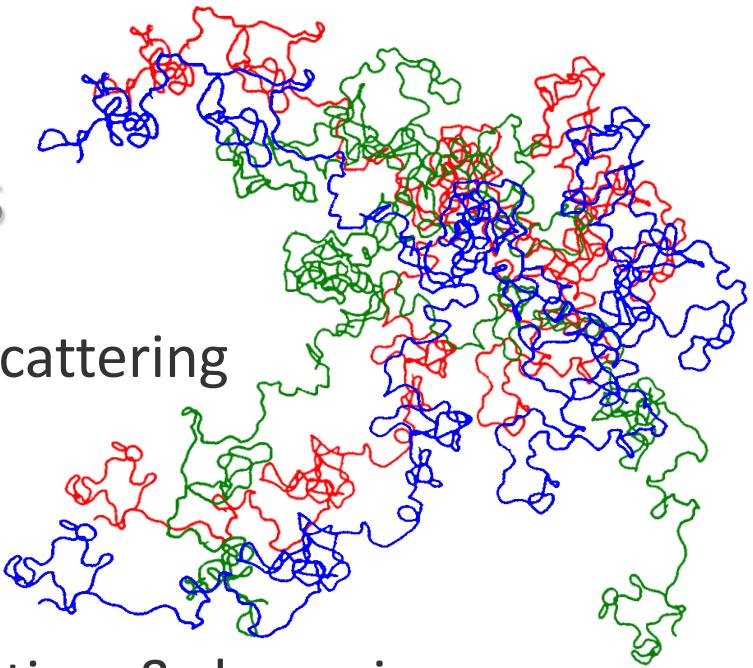
Single chain polymer conformation  
(solution and blends)

Polymer blends: interactions, conformation & dynamics  
(equilibrium and phase separation)

## Lecture 2 – Interfaces and dynamics

Reflectivity and diffusion

Dynamics in soft matter, QENS, BS, Spin-echo

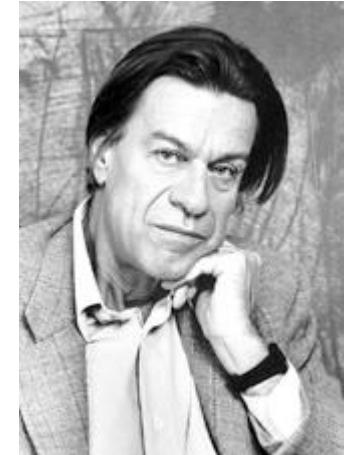


Forster et al (2011)

# Soft Matter

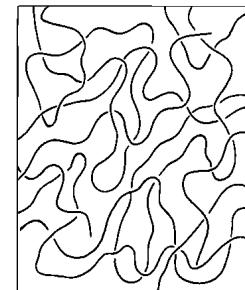
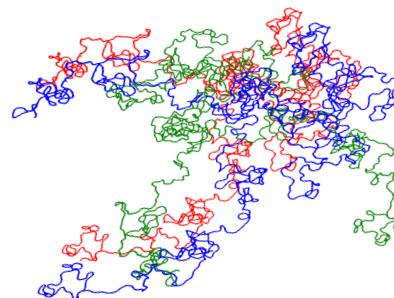
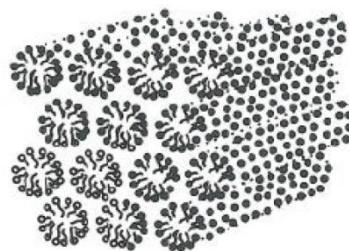
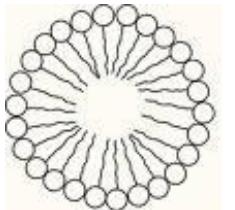
“molecular systems giving a strong response to very weak command signal”

Condensed matter: states are easily deformed by small external fields, including thermal stresses and thermal fluctuations.



deGennes (1991)

Relevant energy scale comparable with room temperature thermal energy.



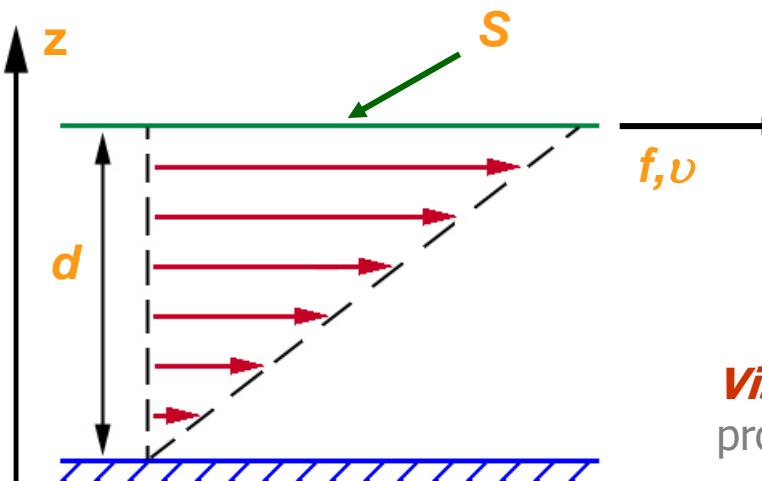
Complex fluids: including colloids, polymers, surfactants, foams, gels, liquid crystals, granular and biological materials.



Movie: complex fluids are generally non-Newtonian... and structured

# Viscosity

Simplest setup  
to measure  
viscosity

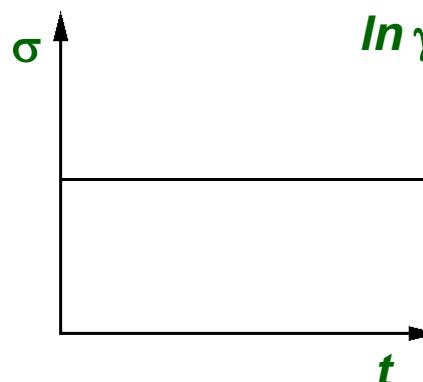


$$f = \eta \frac{Sv}{d}$$

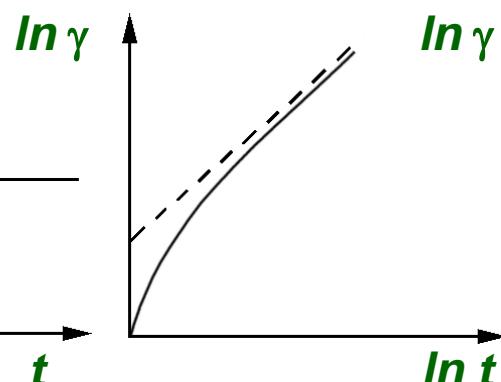
$$\sigma = \frac{f}{S} = \eta \frac{dv}{dz}$$

**Viscosity:**  
proportionality coefficient  $\eta$

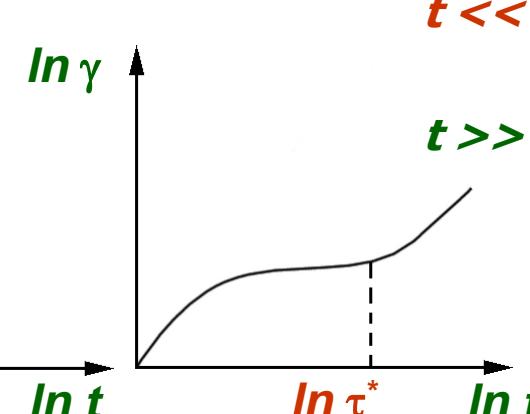
# Viscoelasticity?



Step-wise stress  
starting at  $t=0$



normal liquid



polymer

$t \ll \tau^*$ : **elastic response**

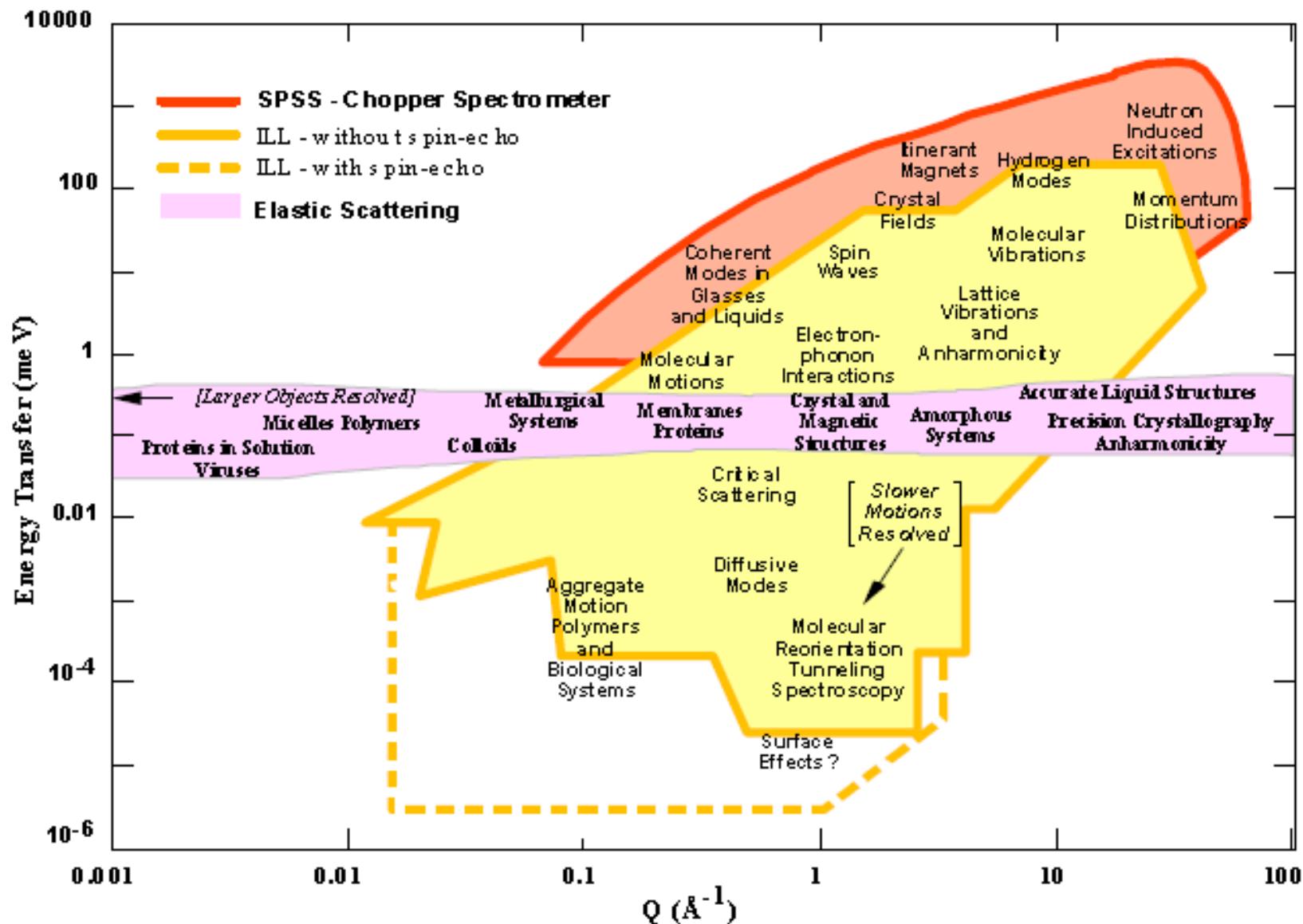
$$\gamma \approx \sigma/E$$

$t >> \tau^*$ : **viscous response**

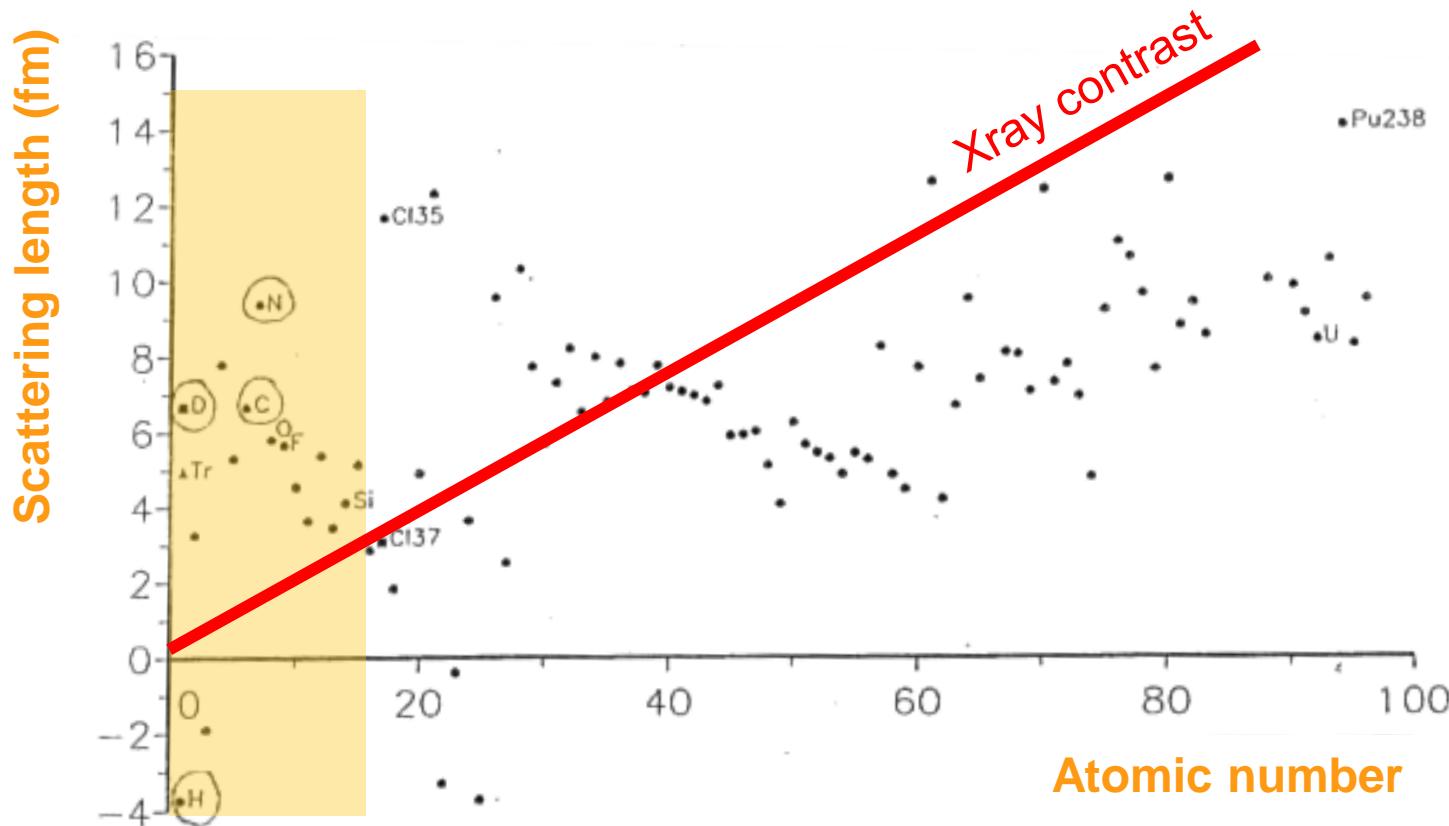
$$\gamma \propto \sigma \cdot t / \eta$$

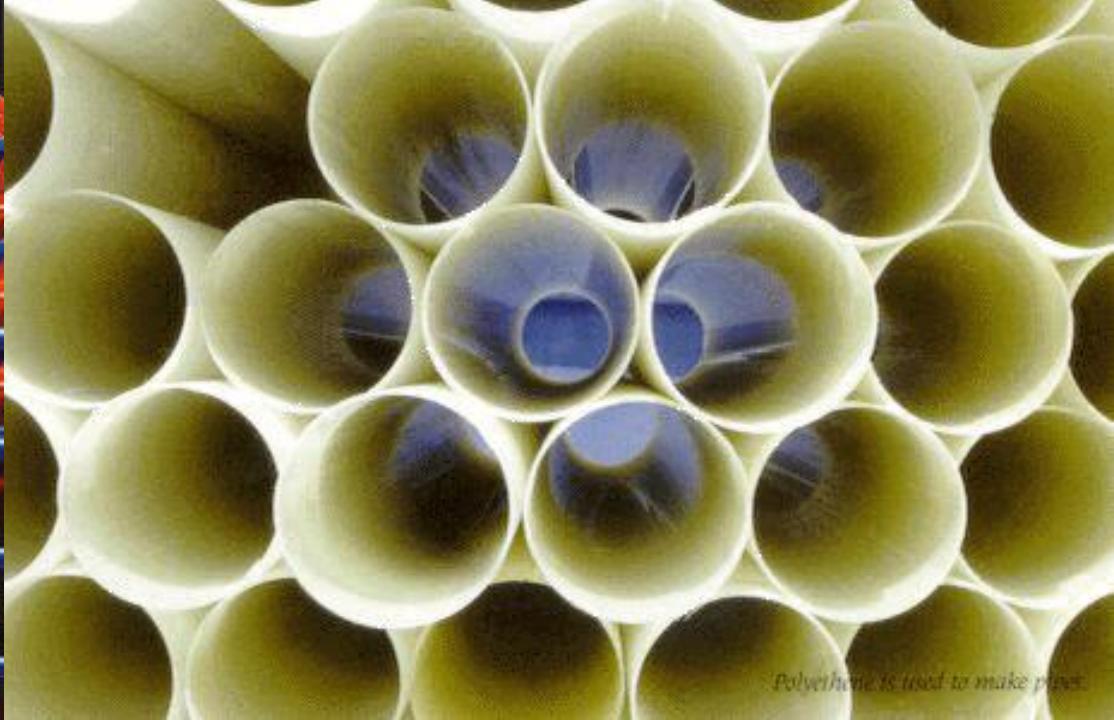
$\gamma$  shear  
angle

# Neutron scattering is key in soft condensed matter



# Neutron scattering is key in soft condensed matter





*Polythene is used to make pipes*

## Common soft matter



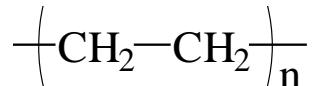
Household products  
packaged in plastic  
containers.  
*Courtesy of British Plastics  
Federation.*



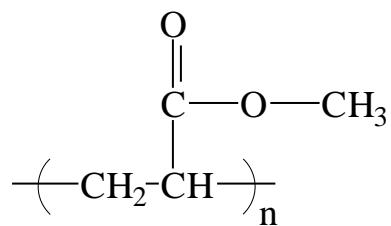
# Speciality polymers

# Common polymers

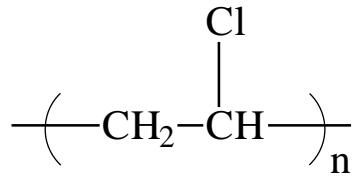
PE poly(ethylene)



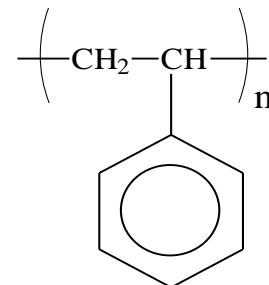
PMMA  
poly(methylmethacrylate)



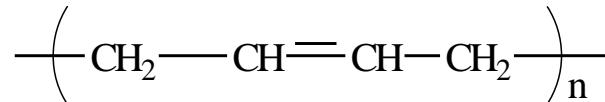
PVC poly(vinylchloride)



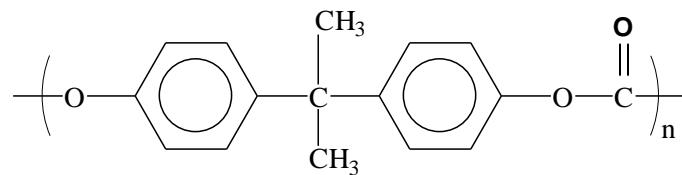
PS poly(styrene)



PB poly(butadiene)



BPA-PC bisphenol-A  
polycarbonate



# Polymer key properties

Molecular weight, N



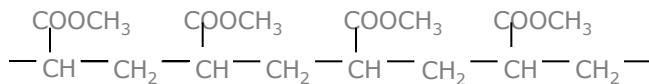
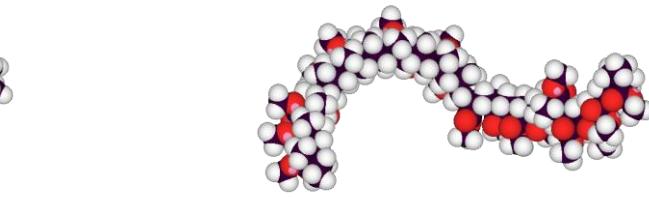
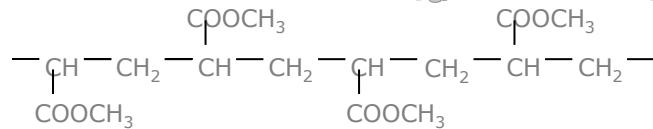
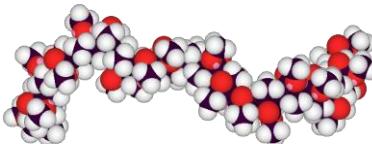
(size)

Polydispersity



(distribution of sizes)

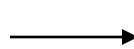
Tacticity



Glass transition (solid-liquid)

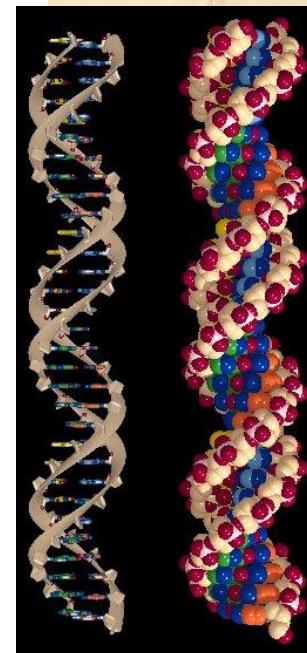
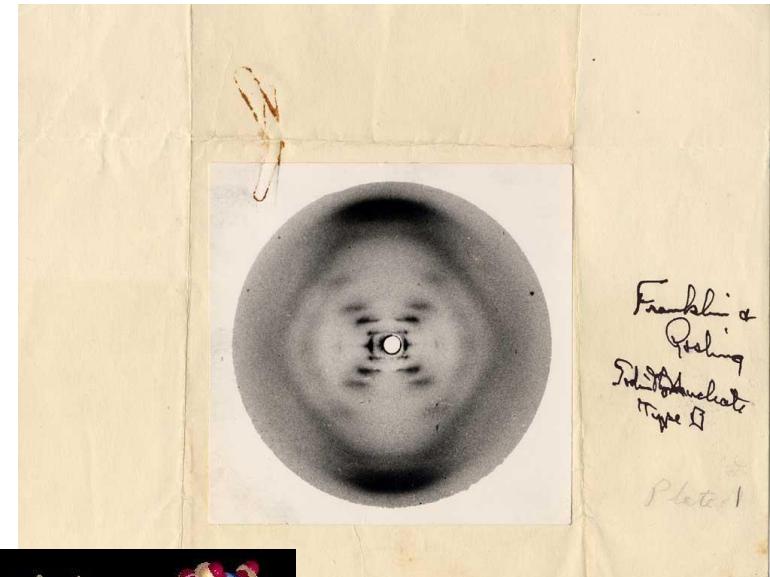
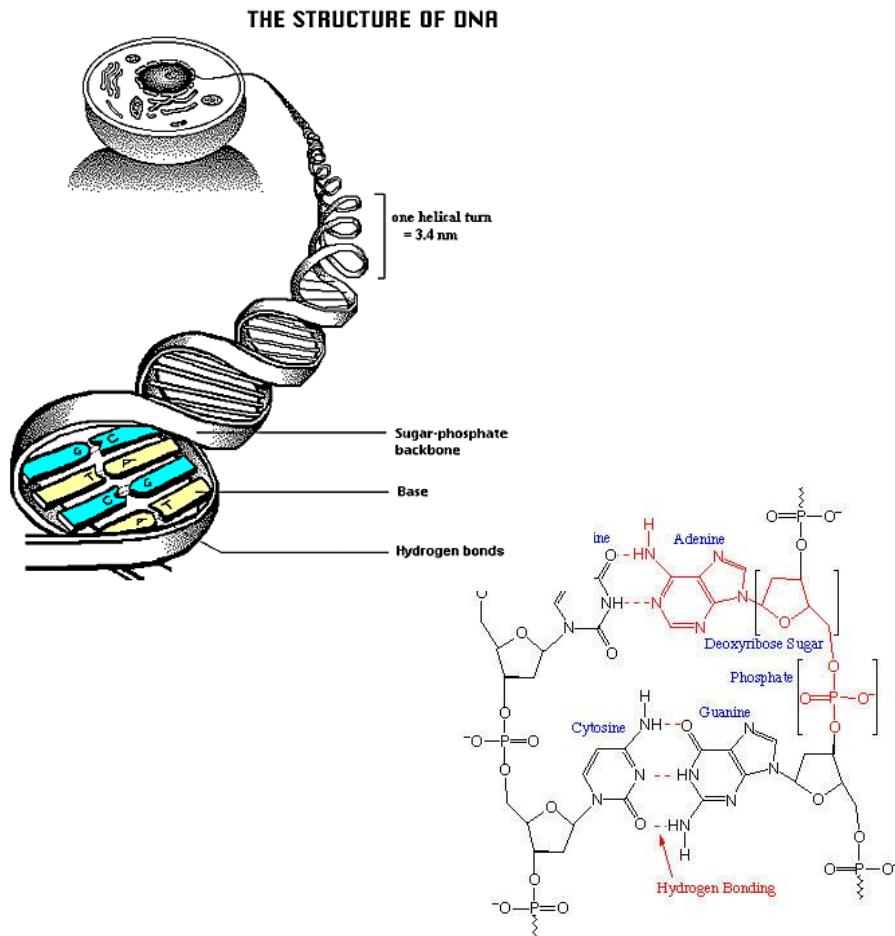
Crystallinity

Interaction parameter  $\chi$

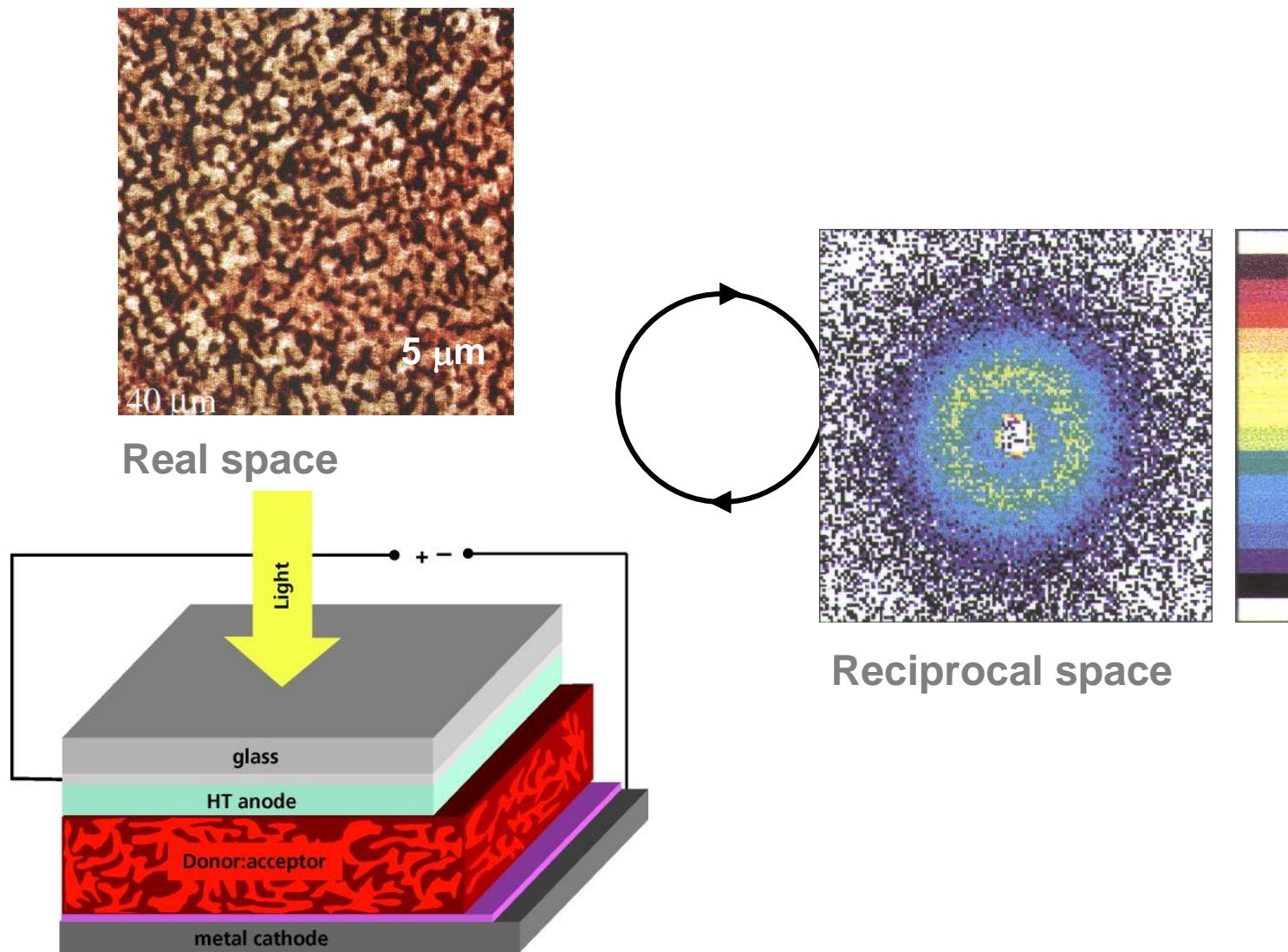


Combine properties to make new materials!

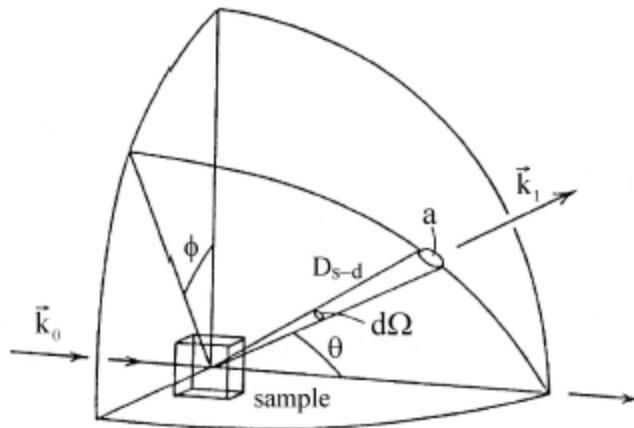
# Soft matter: DNA



# Soft Matter: membranes, photovoltaics (BHJ)



# Scattering theory reminder



## Scattering cross section

$$\frac{d^2\sigma}{d\Omega dE} = \left( \frac{d^2\sigma}{d\Omega dE} \right)_{coh} + \left( \frac{d^2\sigma}{d\Omega dE} \right)_{inc}$$

**coherent    incoherent**

$$\begin{aligned} \left( \frac{d^2\sigma}{d\Omega dE} \right)_{coh} &= \frac{1}{2\pi\hbar} \frac{k_1}{k_0} \frac{\sigma_{coh}}{4\pi} \int_{-\infty}^{+\infty} \sum_{i,j} \left\langle e^{-i\mathbf{q}\cdot\mathbf{R}_i(0)} e^{i\mathbf{q}\cdot\mathbf{R}_j(t)} \right\rangle e^{-i\omega t} dt \\ \left( \frac{d^2\sigma}{d\Omega dE} \right)_{inc} &= \frac{1}{2\pi\hbar} \frac{k_1}{k_0} \frac{\sigma_{inc}}{4\pi} \int_{-\infty}^{+\infty} \sum_i \left\langle e^{-i\mathbf{q}\cdot\mathbf{R}_i(0)} e^{i\mathbf{q}\cdot\mathbf{R}_i(t)} \right\rangle e^{-i\omega t} dt \end{aligned}$$

## Dynamic structure factor

$$\text{FT } (t, \omega) \quad \Updownarrow \quad S(\mathbf{q}, \omega) = \frac{1}{2\pi\hbar} \int_{-\infty}^{+\infty} I(\mathbf{q}, t) e^{-i\omega t} dt.$$

## Intermediate scattering function

$$\text{FT } (r, q) \quad \Updownarrow \quad I_s(\mathbf{q}, t) = \frac{1}{N} \sum_i \left\langle e^{-i\mathbf{q}\cdot\mathbf{R}_i(0)} e^{i\mathbf{q}\cdot\mathbf{R}_i(t)} \right\rangle e^{-i\omega t}.$$

## Pair correlation function

$$G(\mathbf{r}, t) = \frac{1}{(2\pi)^3} \int I(\mathbf{q}, t) e^{-i\mathbf{q}\cdot\mathbf{r}} d\mathbf{q}.$$

$$\int d\omega$$

## Elastic structure factor

$$\int_{-\infty}^{+\infty} S(\mathbf{q}, \omega) |_{\mathbf{q}=Const.} d\omega = S(\mathbf{q})$$

$$S(\mathbf{q})$$

$$S(q) = Nz^2 P(q) + N^2 z^2 Q(q)$$

## Form factor

$$P(q) = \frac{1}{z^2} \sum_{i=1}^z \sum_{j=1}^z \langle e^{-i\mathbf{q}\cdot\mathbf{r}_{ij}} \rangle$$

## Structure factor

$$Q(q) = \frac{1}{z^2} \sum_{i_p=1}^z \sum_{j_q=1}^z \langle e^{-i\mathbf{q}\cdot\mathbf{r}_{ipjq}} \rangle$$

# Reminder: Fourier Transforms

$$H(f) = \int_{-\infty}^{\infty} h(t) e^{2\pi i f t} dt$$

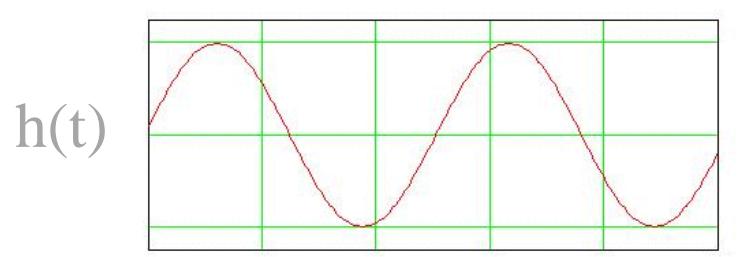
$$h(t) = \int_{-\infty}^{\infty} H(f) e^{-2\pi i f t} df$$



Fourier  
transform:

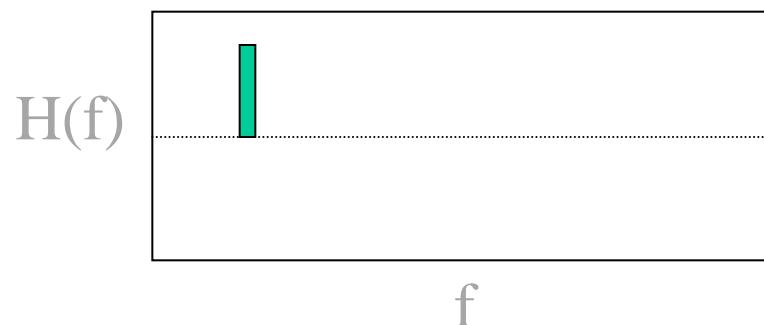
$$\Phi(H(f), t) = h(t)$$

$$\Phi^{-1}(h(t), f) = H(f)$$



$$h(t) = A e^{i(t+\phi)}$$

$$H(f) = \int_{-\infty}^{\infty} h(t) e^{2\pi i f t} dt$$



$$H(f) = 0 \quad \text{if } (ft \neq 1)$$

$$H(f) = A e^{i\phi} \quad \text{if } (ft = 1)$$

# Reminder: Fourier Transforms

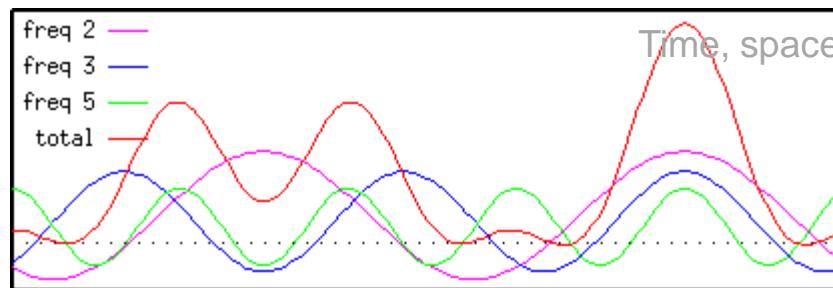
$$H(f) = \int_{-\infty}^{\infty} h(t) e^{2\pi i f t} dt$$

$$h(t) = \int_{-\infty}^{\infty} H(f) e^{-2\pi i f t} df$$

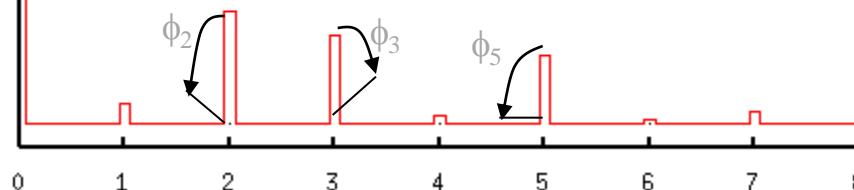
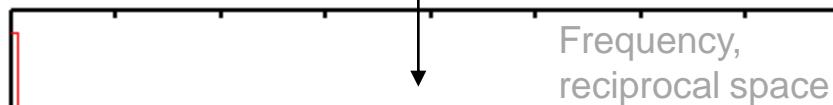
Fourier  
transform:

$$\Phi(H(f), t) = h(t)$$

$$\Phi^{-1}(h(t), f) = H(f)$$

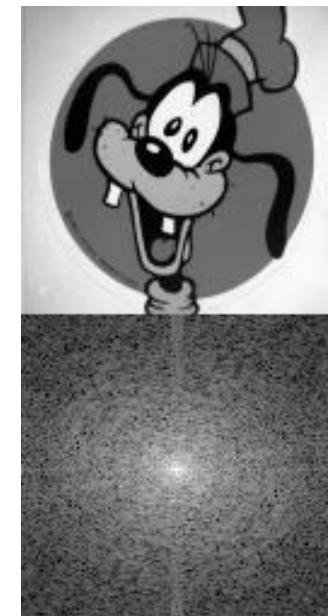
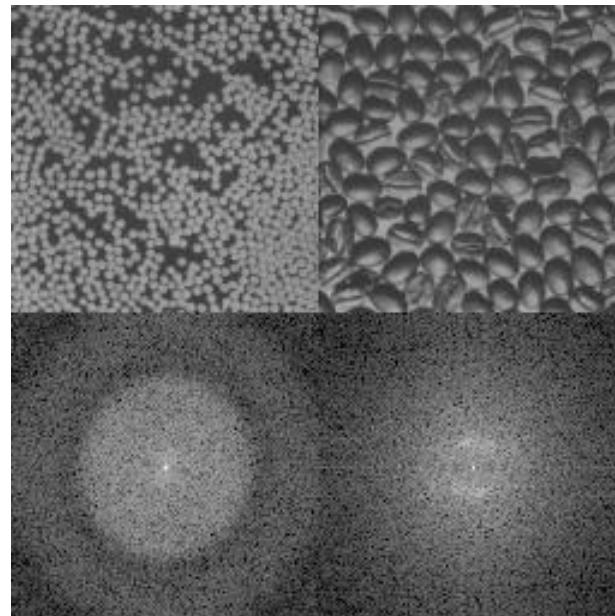
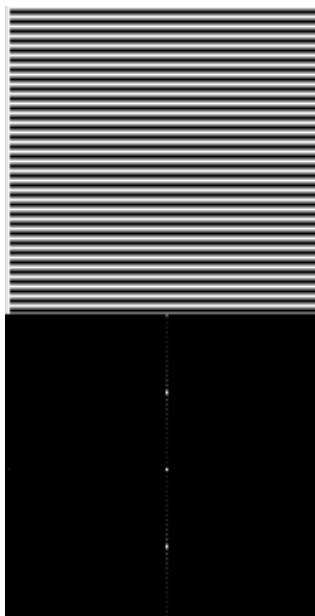
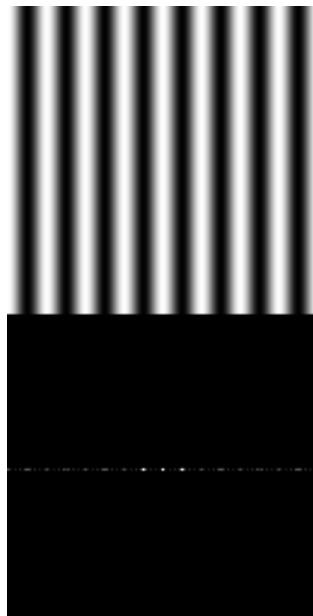


Fourier Transform



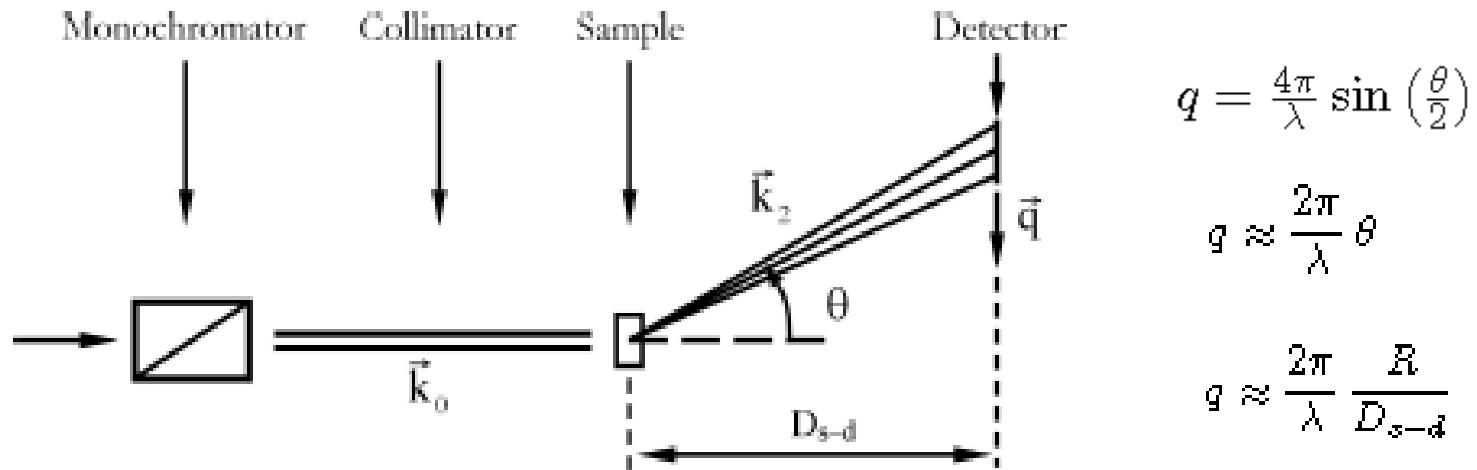
# Reminder: Fourier Transforms

Real space



Reciprocal space

# SMALL-ANGLE NEUTRON SCATTERING



Absolute scattering intensity [cm<sup>-1</sup>]

$$\frac{\partial \sigma}{\partial \Omega}(Q) = N_p V_p^2 (\Delta \delta)^2 P(Q) S(Q) + B_{in} \quad \text{incoherent background}$$

$\diagup$ number density	$\diagdown$ volume contrast	$\mid$ form factor	$\mid$ structure factor	<b>Scattering length density</b>
-----------------------------	--------------------------------	-----------------------	----------------------------	----------------------------------

$$\delta = \sum_i b_i \cdot \frac{D N_A}{M_w}$$

# Relationship between $q$ $\lambda$ $\theta$ and $d$

$$q = \frac{4\pi}{\lambda} \sin(\theta/2)$$

$$d = \frac{2\pi}{q}$$

small  $q \sim$  large  $d$       [large  $q \sim$  small  $d$ ]  
small  $\lambda \sim$  large  $q \sim$  small  $d$   
[large  $\lambda \sim$  small  $q \sim$  large  $d$ ]

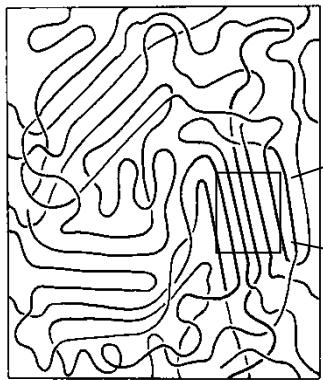
Radiation	Wavelength
light	~ 500 nm
X-rays	~ 1 Å
neutrons	~ 5 Å

Ångstrom:  
1 Å =  $10^{-10}$  m  
1 nm = 10 Å

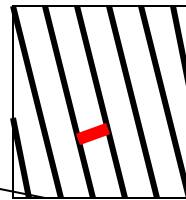
## Bottom line:

radiation of small wavelength  $\lambda$  can ‘see’ smaller sample features  $d$   
(provided that contrast is sufficient).

# Example: crystalline structure of polymer



Semi-crystalline poly(ethylene) PE



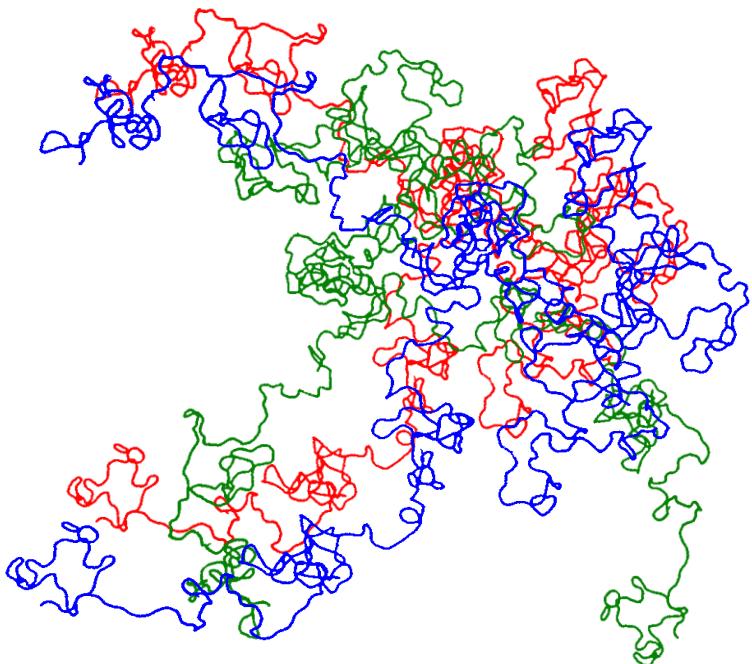
lamella spacing is  $d \sim 20$   
nm  
 $q = 2\pi/d \sim 0.3 \text{ nm}^{-1}$

$$q = \frac{4\pi}{\lambda} \sin(\theta/2)$$

Radiation	Wavelength	Scattering angle
X-rays	$\sim 1 \text{ \AA}$	$\theta \sim 0.27 \text{ degrees}$
neutrons	$\sim 5 \text{ \AA}$	$\sin(\theta/2) \sim 0.013, \theta \sim 0.7 \text{ degrees}$
light	$\sim 500 \text{ nm}$	$\sin(\theta/2) > 1$ impossible!

the smallest dimension probed by wavelength  $\lambda$  corresponds to largest angle  $\theta=180$  degrees (backscattering). For light  $d_{min} \sim 0.25 \mu\text{m}$ , for X-rays or neutrons  $d_{min} \sim 0.5$  to  $2.5 \text{ \AA}$ .  
In typical experiments, scattering angles range from  $0.1 < \theta < 70$  degrees

# A brief history of polymers



1920 Staudinger proposed plastics are long molecules with covalent bonds.

1940 Kuhn established they are usually flexible

1953 Flory described the shape and size of individual polymer molecules in solutions and melts.

1967 Edwards – polymer molecules in rubbers and glasses are trapped by neighbours in a “tube”. Chain confirmation  $\leftrightarrow$  trajectories quantum particles.

1971 de Gennes – molecules in tube move like snakes and eventually escape

# Main factors governing polymer physical properties :

1. Number of monomer units in the chain, N is large  $N \gg 1$ .

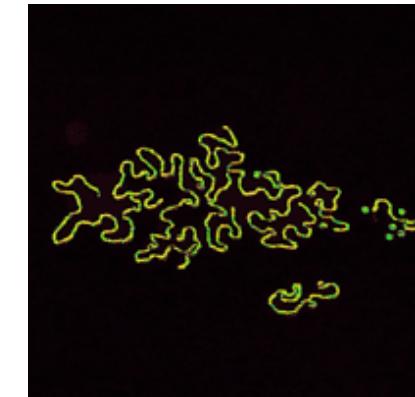
synthetic polymers

$$N \approx 10^2 - 10^4$$

DNA

$$N \approx 10^9 - 10^{10}$$

2. Monomer units are connected in the chain.  $\Rightarrow$  no freedom of independent motion (unlike systems of disconnected particles, e.g. low molecular gases and liquids).  $\Rightarrow$  Polymer systems are poor in entropy.



3. Polymer chains are generally *flexible*.

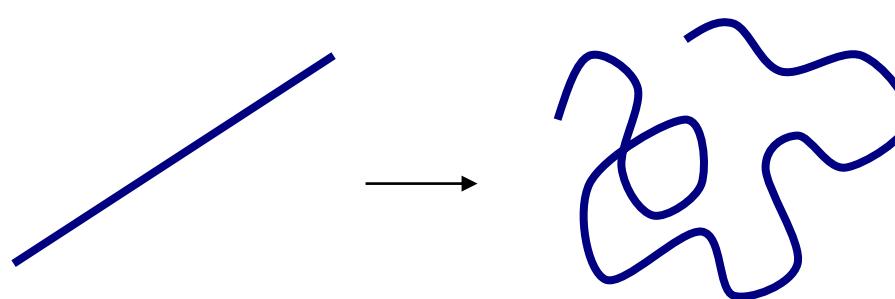
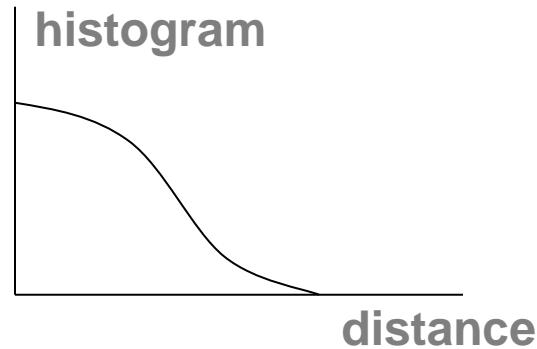
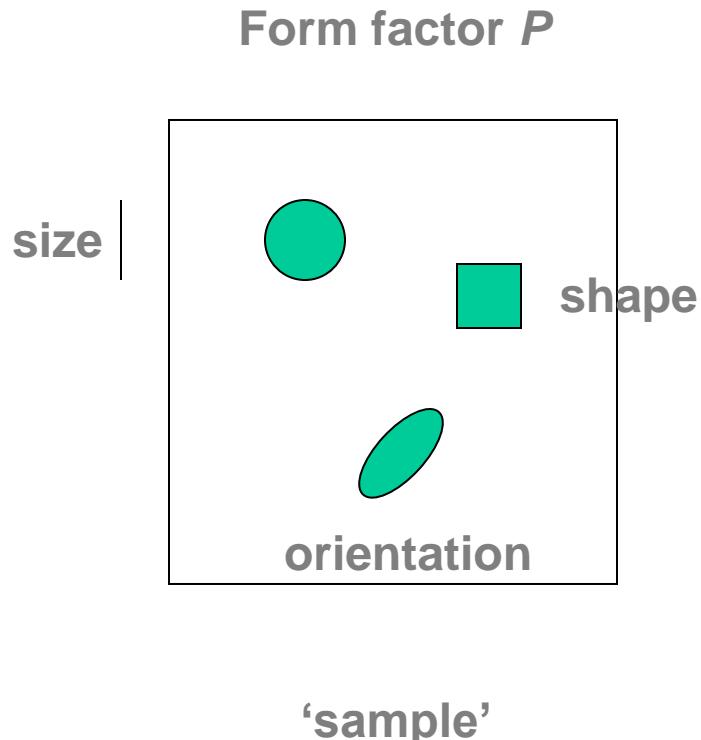


Image single molecule (AFM)  
Magonov

# Form factor: P

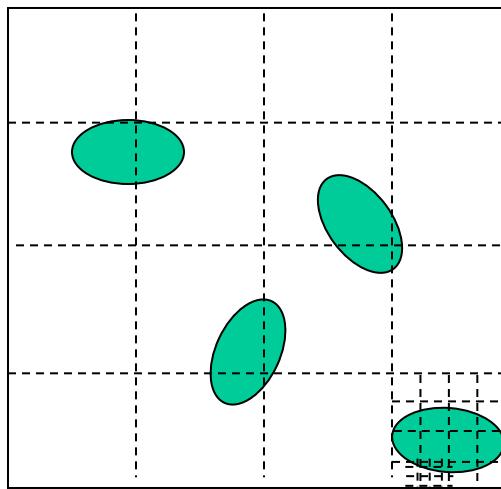
*Self-correlations*



$$P(Q) = \frac{1}{V_p^2} \left| \int_0^{V_p} \exp [i f(Q \alpha)] dV_p \right|$$

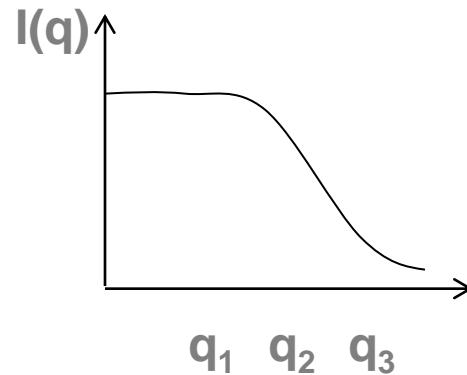
Interference between scattered radiation  
from different parts of the same object  
(analytical solutions for common shapes)

# Multiple lengthscales



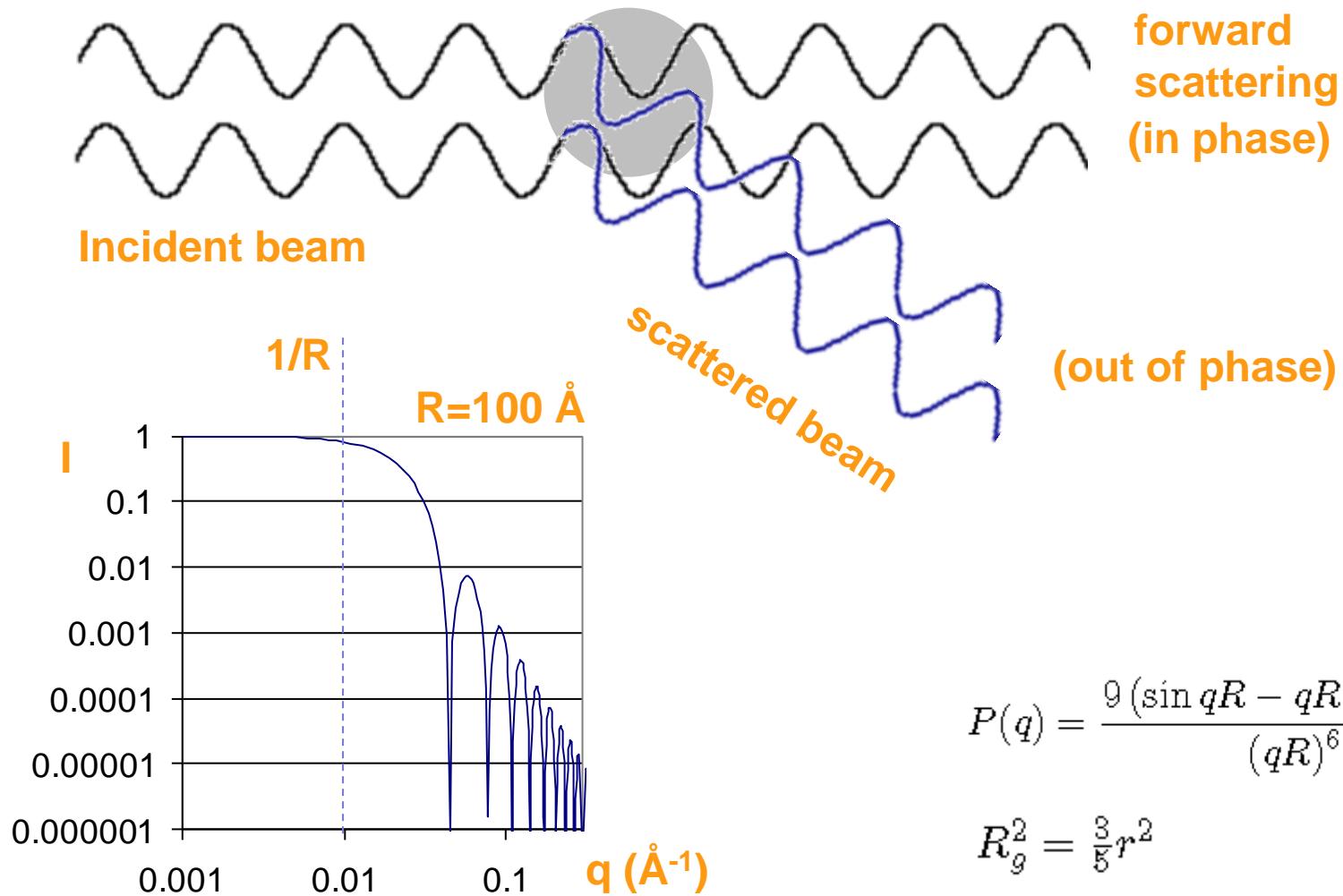
$$2\pi/q_1$$

$$2\pi/q_3 \quad 2\pi/q_2$$

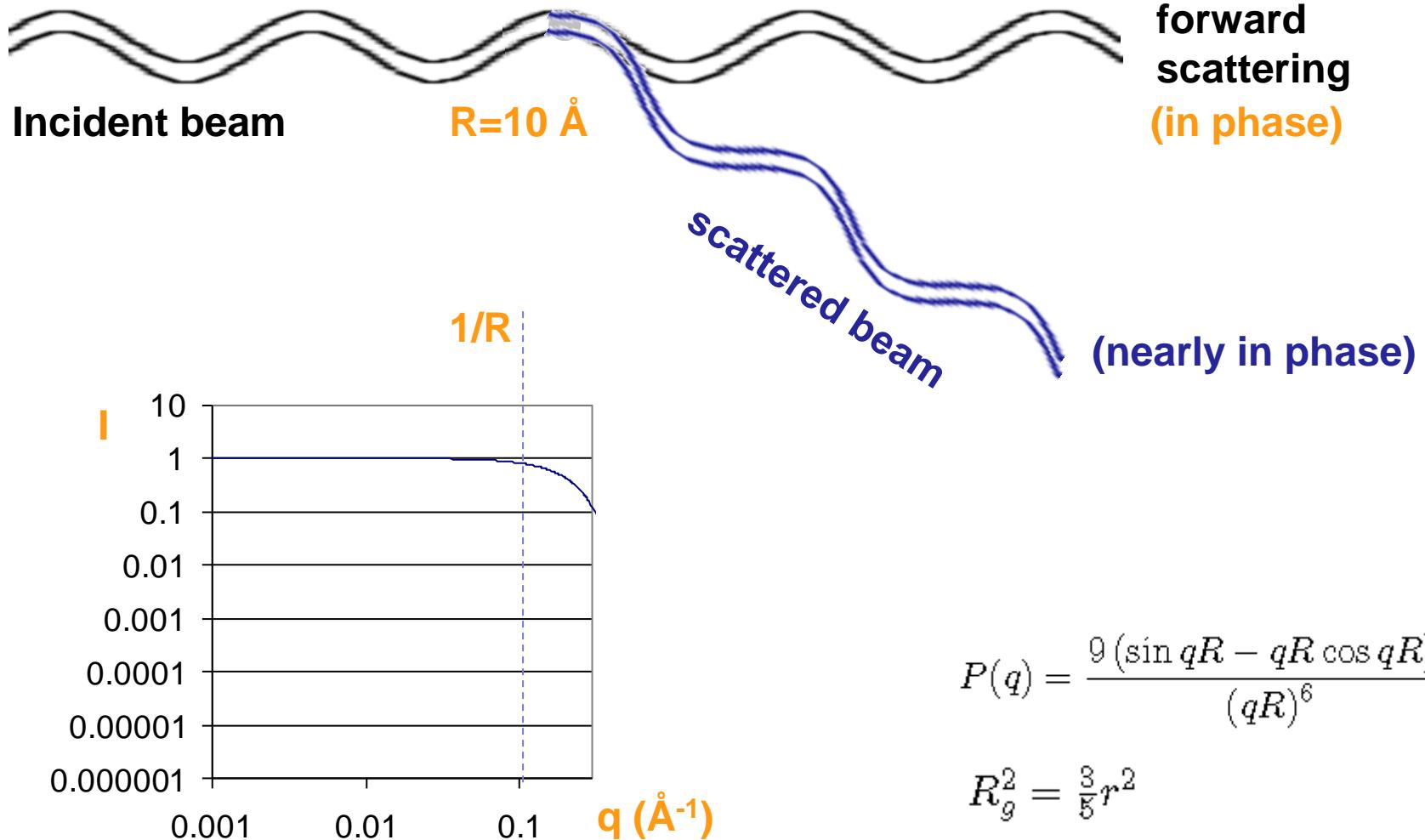


scattering spectrum corresponds to different “magnifications”, thus several approximations may be relevant

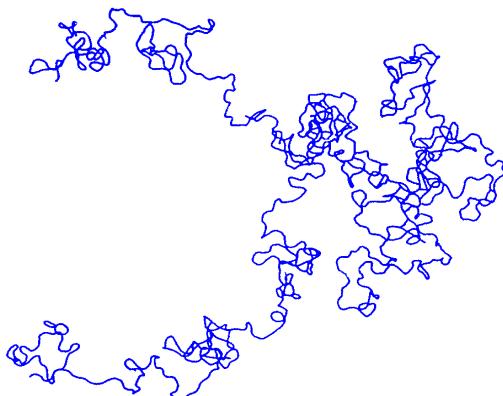
# Scattering from a sphere



# Scattering from a (tiny) sphere



# Scattering from a random coil



Debye form factor

$$g_D(x) = \frac{2}{x^2} (x - 1 + e^{-x})$$

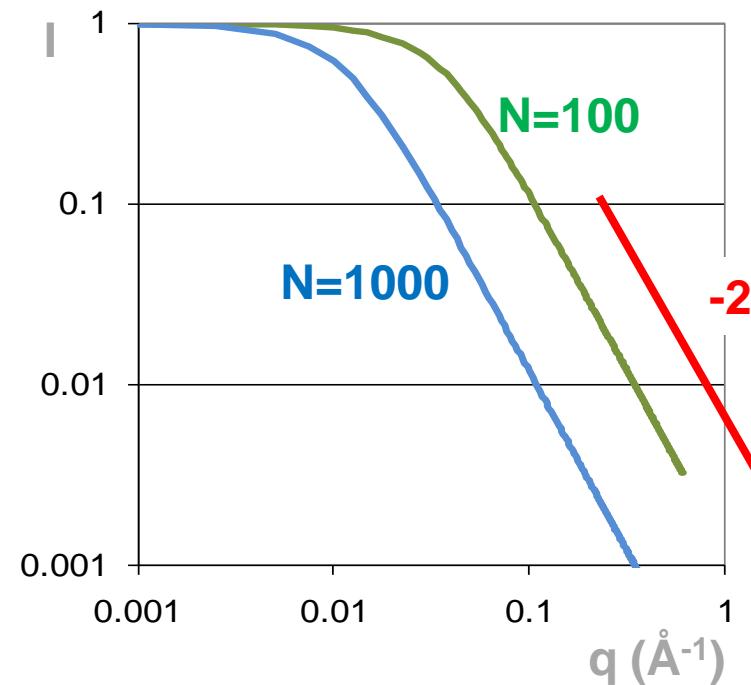
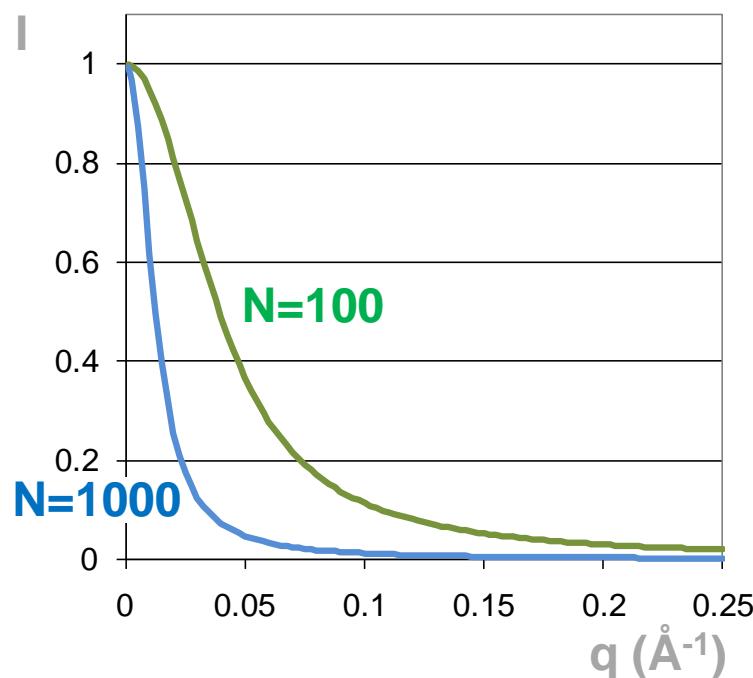
$$x \equiv q^2 R_g^2$$

$$R_g^2 = \frac{Na^2}{6}$$

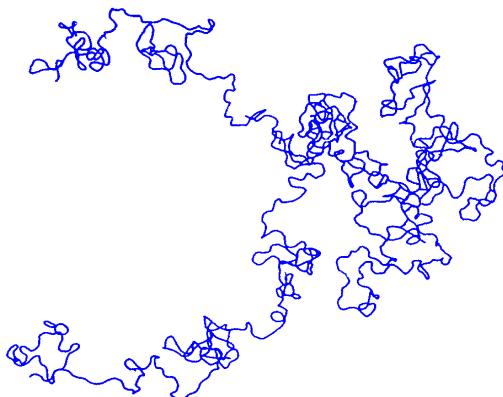
$a = 10 \text{ \AA}$

$N=100 \rightarrow R_g \approx 4 \text{ nm}$

$N=1000 \rightarrow R_g \approx 13 \text{ nm}$



# Scattering from a random coil



Debye form factor

$$g_D(x) = \frac{2}{x^2} (x - 1 + e^{-x})$$

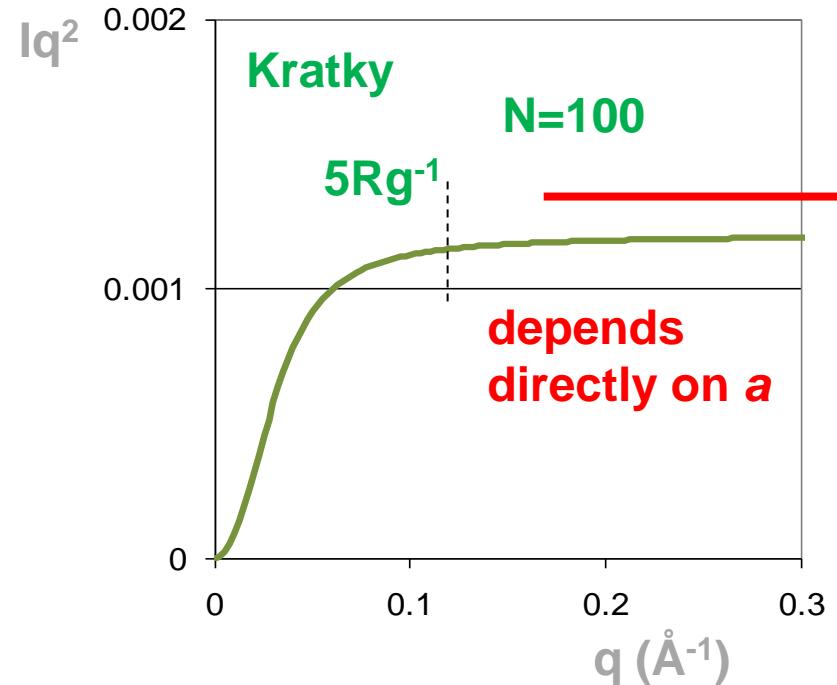
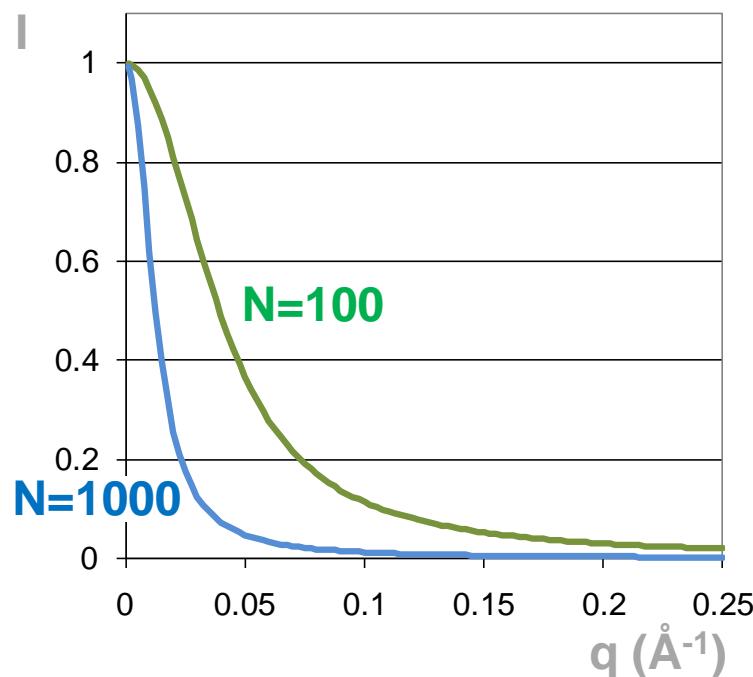
$$x \equiv q^2 R_g^2$$

$$R_g^2 = \frac{Na^2}{6}$$

$$a=10 \text{ \AA}$$

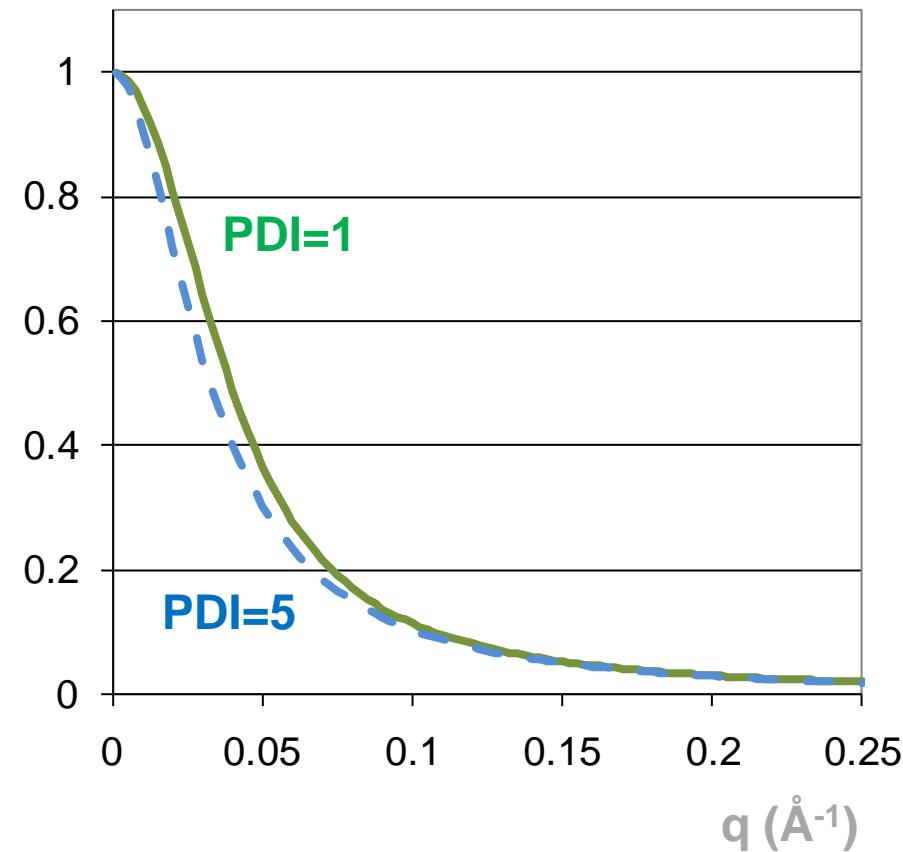
$$N=100 \rightarrow R_g \approx 4 \text{ nm}$$

$$N=1000 \rightarrow R_g \approx 13 \text{ nm}$$



# Polydisperse random coils

$M_n/M_w \ g_D$



Polydisperse debye form factor

$$g_D(x) = \frac{2}{(1 + 1/h)x^2} \left[ \left(1 + \frac{x}{h}\right)^{-h} - 1 + x \right]$$

$$x \equiv q^2 \langle R_g^2 \rangle_n \equiv \frac{q^2 \langle R_g^2 \rangle_z}{1 + 2/h}$$

$$h = \left( \frac{M_w}{M_n} - 1 \right)^{-1}$$

(normalised to PDI)

$a=10 \text{ \AA}$

$N=100 \rightarrow Rg \approx 4 \text{ nm}$

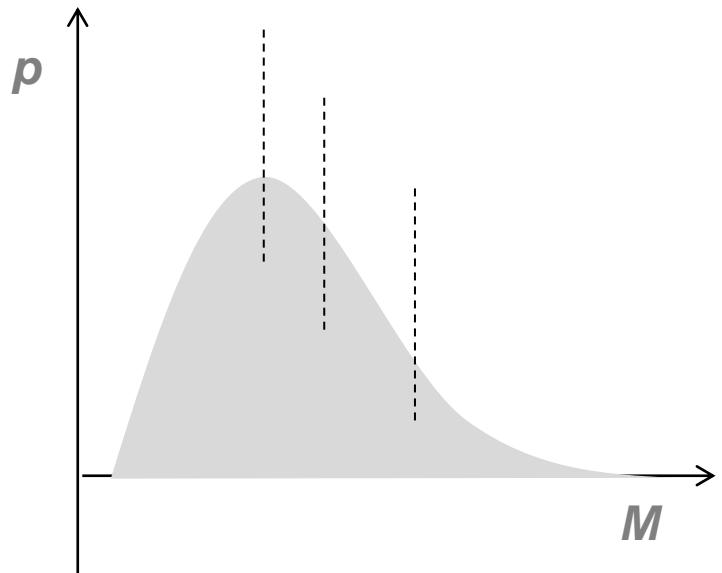
$N_w=100, N_w/N_n=5$

# A polydispersity model

(Schultz-Zimm)

$$p(M) = \frac{h^h}{\Gamma(h)} \left( \frac{M}{M_n} \right)^h e^{-h(\frac{M}{M_n})}$$

$$h = \left( \frac{M_w}{M_n} - 1 \right)^{-1}$$

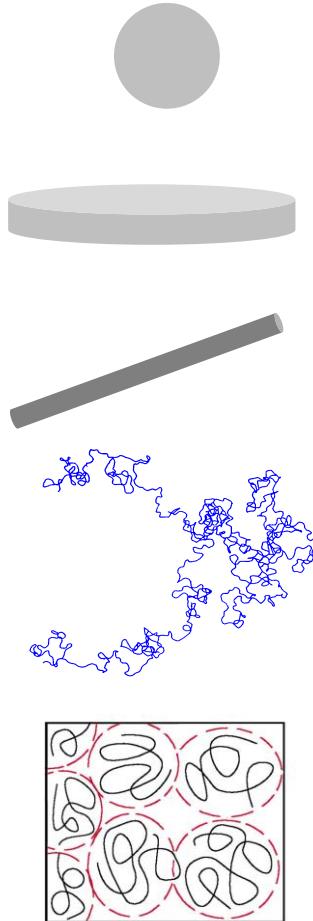


Number average  $M_n = \int p(M) M dM$

Weight average  $M_w = \frac{\int p(M) M^2 dM}{\int p(M) M dM} \equiv \frac{\langle M \rangle_2}{\langle M \rangle_1}$

Z-average  $M_z = \frac{\int p(M) M^3 dM}{\int p(M) M^2 dM} \equiv \frac{\langle M \rangle_3}{\langle M \rangle_2}$

# Useful form factors

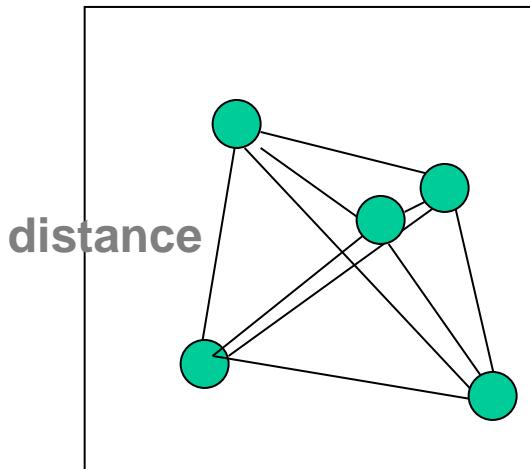


Sphere of radius $R_p$	$P(Q) = \left[ \frac{3(\sin(QR_p) - QR_p \cos(QR_p))}{(QR_p)^3} \right]^2$
Disc of negligible thickness and radius $R_p$ ( $J_1$ is a first-order Bessel function)	$P(Q) = \frac{2}{(QR_p)^2} \left[ 1 - \frac{J_1(2QR_p)}{QR_p} \right]$
Rod of negligible cross-section and length $L$ ( $S_i$ is the Sine integral function)	$P(Q) = \frac{2S_i(QL)}{QL} - \frac{\sin^2(QL/2)}{(QL/2)}$
Gaussian random coil with z-average radius of gyration $R_g$ , polydispersity ( $Y+1$ ) and $U = \frac{(QR_g)^2}{(1+2Y)}$	$P(Q) = \frac{2[(1+UY)^{-\frac{1}{2}} + U - 1]}{(1+Y) U^2}$
Concentrated polymer solution with screening length $\xi$ where $\zeta = R_p \left( \frac{\phi}{\phi^*} \right)^{\frac{1}{1-\phi}}$	$P(Q) = P(0) \left[ \frac{1}{1 + (Q\xi)^2} \right]$

S King

<http://www.ncnr.nist.gov/resources/>

# Structure factor: S

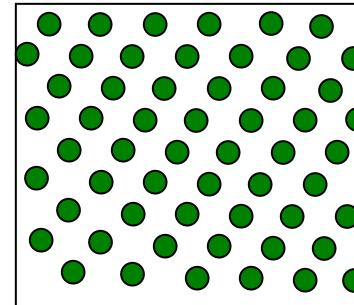


Interference between radiation scattered by distinct objects

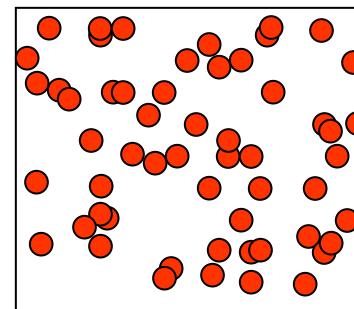
$$S(Q) = 1 + \frac{4\pi N_p}{QV} \int_0^\infty [g(r) - 1] r \sin(Qr) dr$$

$$G(r) = \frac{4\pi N_p r^2}{V} g(r)$$

Radial distribution function,  
provides information about their  
relative position

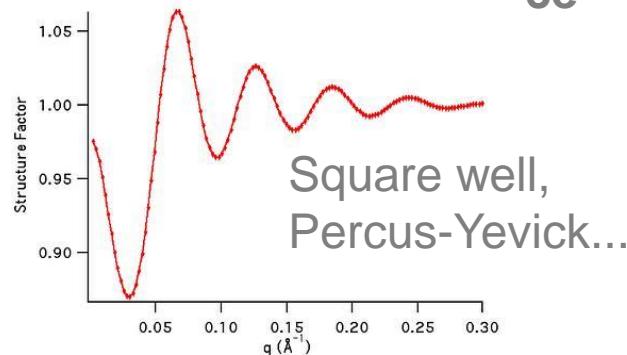
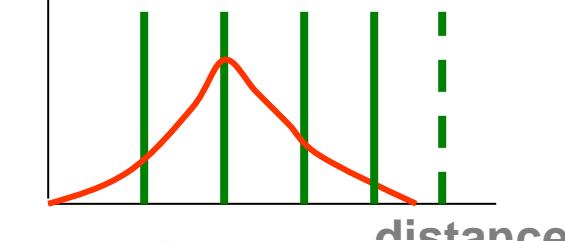


Ordered structure  
'crystal'



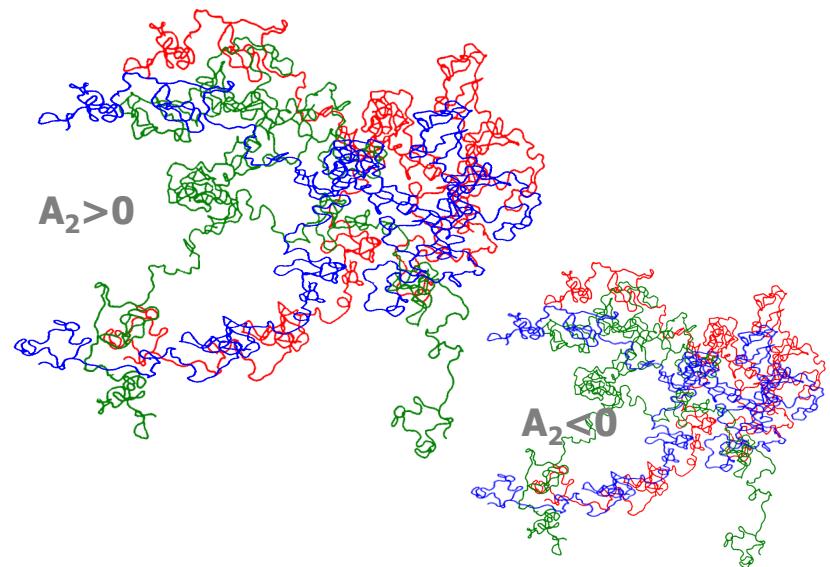
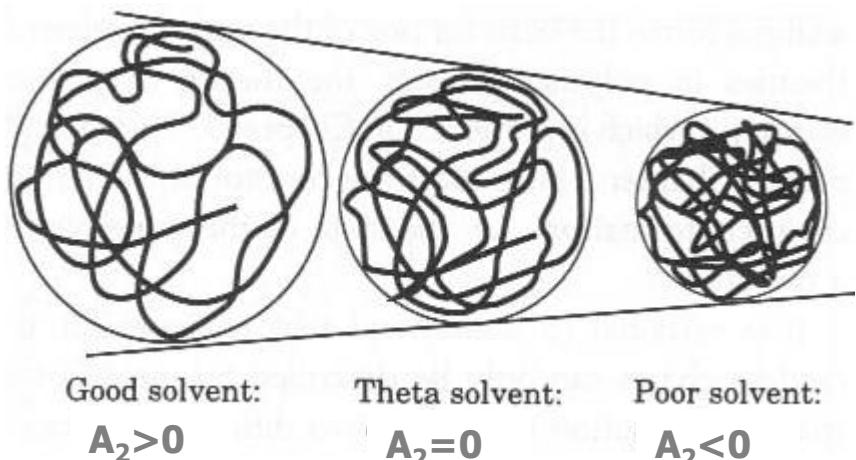
Disordered structure

histogram



# Interactions: Polymers in solution and melt

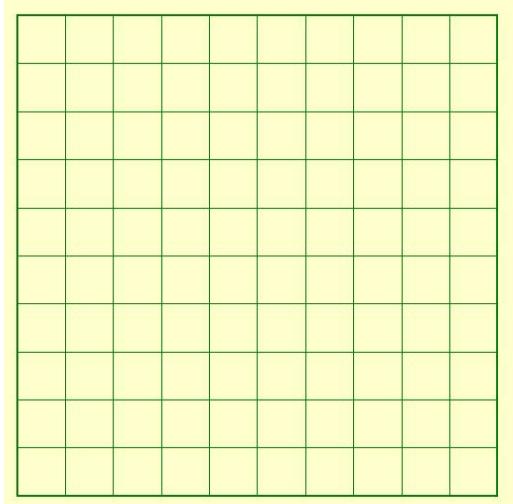
1953 Flory described the shape and size of individual polymer molecules in solutions and melts.



Ideal chains occur in  $\theta$ -solutions (ie, neutral solvent,  $A_2=0$ ) or in melt. Chains expand or contract depending on interactions:  $A_2$  (Second Virial coeff, for solutions) or  $\chi$  (for polymer mixtures)

# Polymer miscibility (1)

**Flory-Huggins  
lattice**



Binary mixture

Thermodynamics       $\Delta G_m = \Delta H_m - T\Delta S_m$

Combinatorial entropy

$$-\frac{\Delta S}{R} = n_A \ln \phi_A + n_B \ln \phi_B \quad \Omega \text{ Boltzmann law}$$

Enthalpy     $\Delta H_m = K_B T \phi_A \phi_B \chi_{AB}$

$$\frac{\Delta G_m}{K_B T} = \frac{\phi_A}{N_A} \ln \phi_A + \frac{\phi_B}{N_B} \ln \phi_B + \phi_A \phi_B \chi_{AB}$$

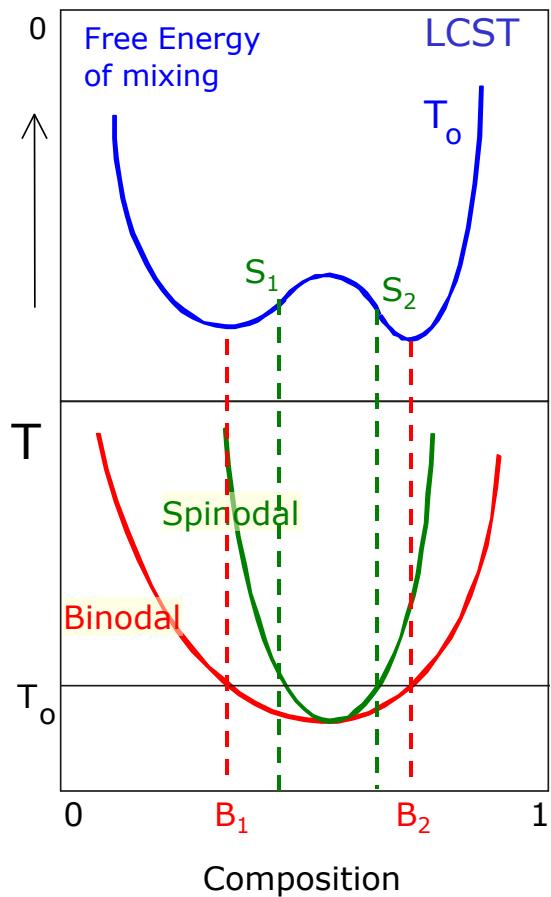
Combinatorial  
entropy

Enthalpy

$\chi < 0$  mixing occurs  $\forall T, \phi$

$\chi > 0$   $\Delta G_m = \Delta H_m - T\Delta S_m$  only at high T

# Polymer miscibility (2)



Thermodynamics

$$\Delta G_m = \Delta H_m - T\Delta S_m$$

$$\frac{\Delta G_m}{k_B T} = \frac{\phi}{v_A N_A} \ln \phi + \frac{(1-\phi)}{v_B N_B} \ln(1-\phi) + \frac{\phi(1-\phi)}{v} \chi$$

Combinatorial entropy

Enthalpy

Phase boundaries?

Binodal

$$\left\{ \begin{array}{l} \frac{\partial \Delta G_m}{\partial \phi_{B1}} = \frac{\partial \Delta G_m}{\partial \phi_{B2}} \equiv \mu \\ \Delta G_m(\phi_{B1}) + \Delta G_m(\phi_{B2}) = \min \end{array} \right.$$

*'minima'*

Spinodal

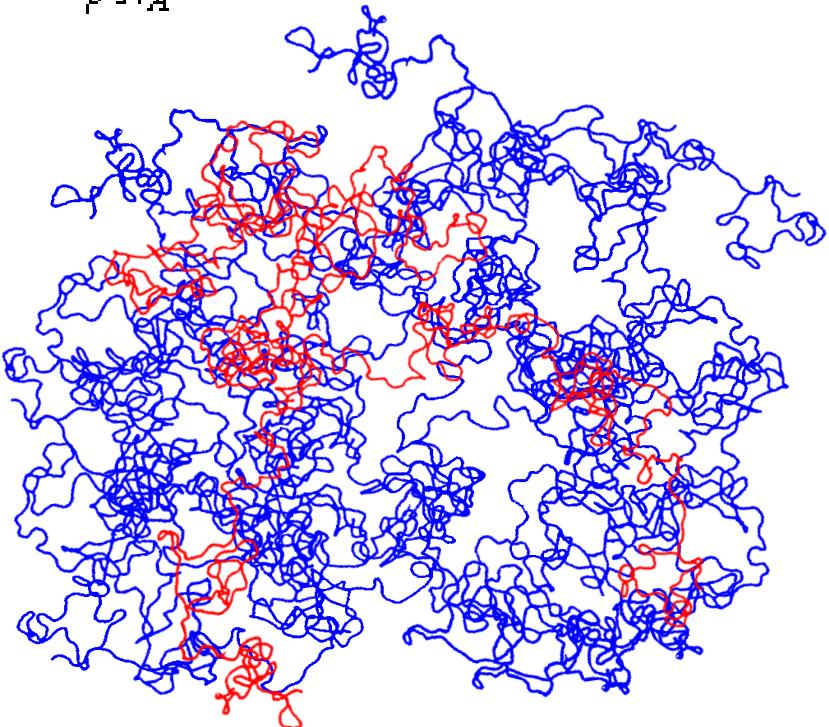
$$\frac{\partial^2 \Delta G_m}{\partial \phi^2} = 0$$

*inflection points*

# Isotopic polymer mixture

$$\frac{\partial \sigma}{\partial \Omega}(q) = (b_D - b_H)^2 S_{DD}(q) = (b_D - b_H)^2 \phi(1 - \phi) N_A^2 P(q)$$

$$\begin{aligned} \left. \frac{1}{V} \frac{d\sigma}{d\Omega}(q) \right|_{\text{eqm}} &= (b_D - b_H)^2 \phi(1 - \phi) \langle M' \rangle_w \frac{\rho N_A}{m^2} P(q) \\ &= (b_D - b_H)^2 \phi(1 - \phi) \langle M' \rangle_w \frac{(\Delta \phi)^2}{\rho N_A} P(q) \end{aligned}$$



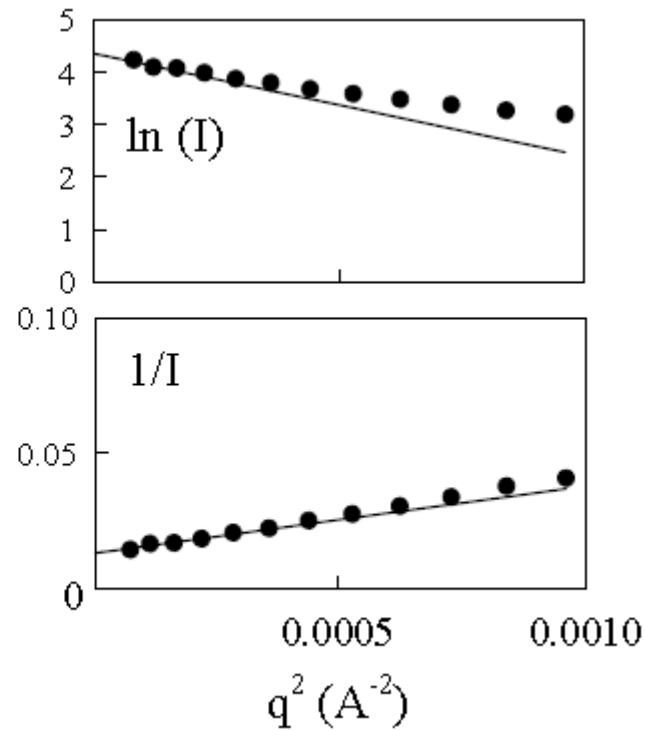
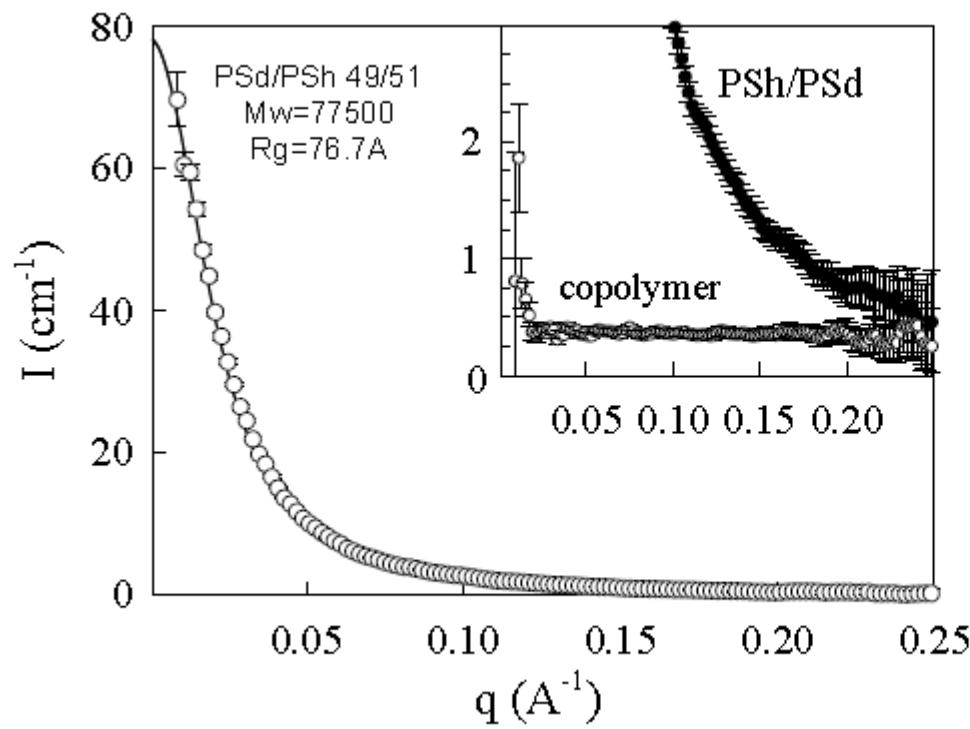
# Approximations: Guinier & Zimm

$$\text{Guinier : } \frac{d\Sigma(q)}{d\Omega} \approx \frac{d\Sigma(0)}{d\Omega} \left( -\frac{(qR_g)^2}{3} \right)$$

where  $\frac{d\Sigma(0)}{d\Omega} = \frac{\phi(1-\phi) M_w (\Delta\rho)^2}{N_A \rho}$

$$\text{Zimm : } \left[ \frac{d\Sigma(q)}{d\Omega} \right]^{-1} \approx \left[ \frac{d\Sigma(0)}{d\Omega} \right]^{-1} \left[ 1 + \frac{(qR_g)^2}{3} \right]$$

for a polymer coil



# Interacting polymer mixtures

$$\frac{1}{V} \frac{d\sigma}{d\Omega}(q) \Big|_{cm=1} = N_A \left( \frac{b_1}{v_1} - \frac{b_2}{v_2} \right)^2 S(q)$$

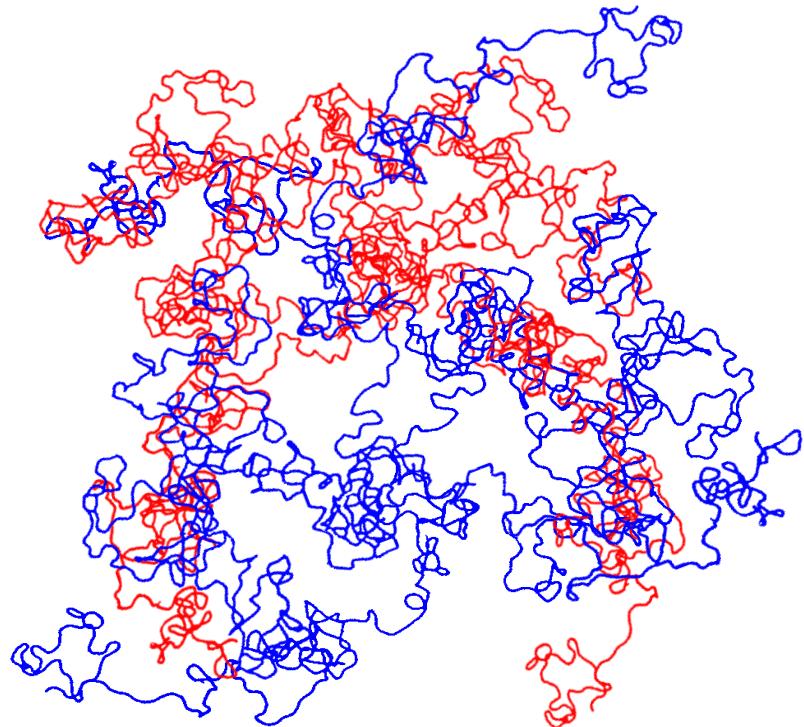
$$\frac{1}{S(q)} = \frac{1}{S_1(q)} + \frac{1}{S_2(q)} - 2 \frac{\tilde{\chi}_{12}}{v_0}$$

$$S(q) = \phi_1 v_1 \langle N_1 \rangle_\omega \langle g_D(R_{gq}^2 q^2) \rangle_\omega$$

**Zimm**  $S(q) \approx \phi_1 v_1 \langle N_1 \rangle_\omega \left( 1 - \frac{1}{3} \langle R_{gq}^2 \rangle_z q^2 \right)$

$$\frac{1}{S(q)} = \frac{1}{S(0)} \left[ 1 + \frac{1}{3} R_{\text{eff}}^2 q^2 \right] \quad \text{where} \quad R_{\text{eff}}^2 = \left( \frac{\langle R_{g1}^2 \rangle_z}{\phi_1 v_1 \langle N_1 \rangle_\omega} + \frac{\langle R_{g2}^2 \rangle_z}{\phi_2 v_2 \langle N_2 \rangle_\omega} \right) S(0)$$

$$\frac{1}{S(0)} = \frac{1}{\phi_1 v_1 \langle N_1 \rangle_\omega} + \frac{1}{\phi_2 v_2 \langle N_2 \rangle_\omega} - 2 \frac{\tilde{\chi}_{12}}{v_0}$$



# Random Phase Approximation

RPA (de Gennes, 1979):

$$\frac{1}{S(q)} = \frac{1}{S_1(q)} + \frac{1}{S_2(q)} - 2 \frac{\tilde{\chi}_{12}}{v_o} \quad 1/S(0) = G''$$

Ornstein-Zenike:

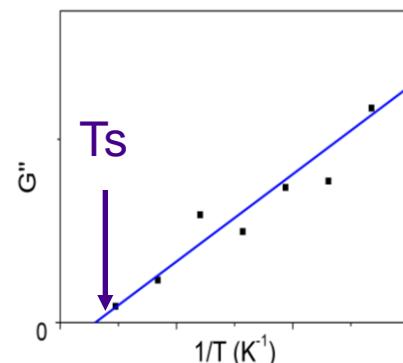
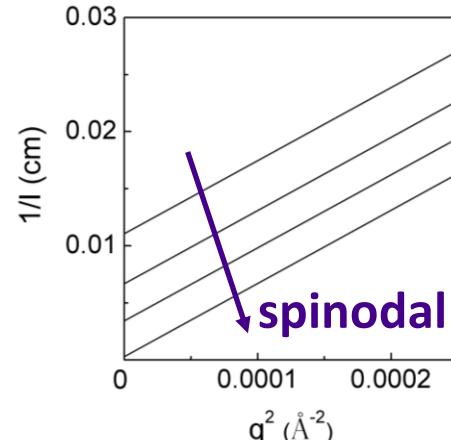
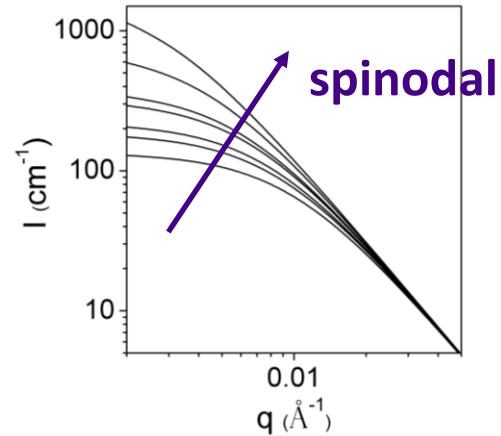
At low angle ( $qRg \ll 1$ ),

$$S(q) = \frac{S(0)}{1 + \xi^2 q^2}$$

Linear in the Zimm representation as:

$$\begin{aligned} 1/S(q) &= 1/S(0) + Aq^2 \\ \frac{1}{S(q)} &= 2(\chi_s - \chi_f) + \frac{\xi^2}{S(0)} q^2 \end{aligned}$$

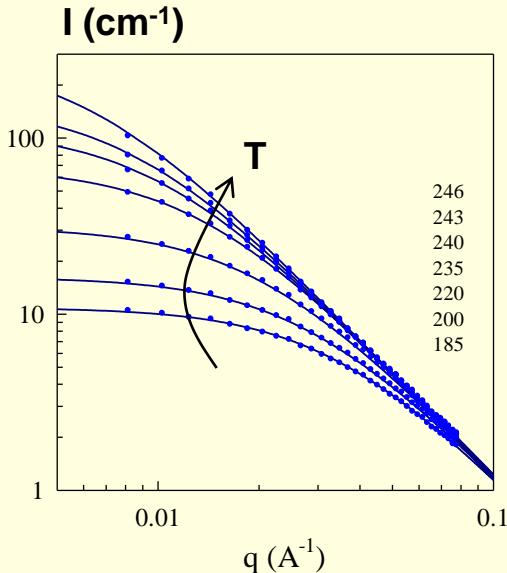
$T_s$  determined by extrapolation of  $G''$  to 0



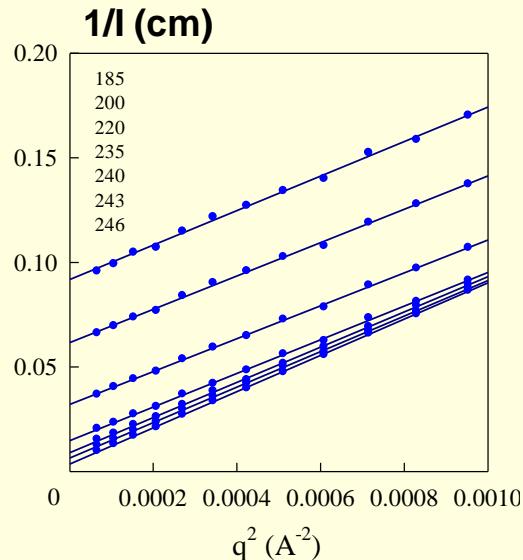
# Equilibrium: SANS

TMPC/PSd 50/50

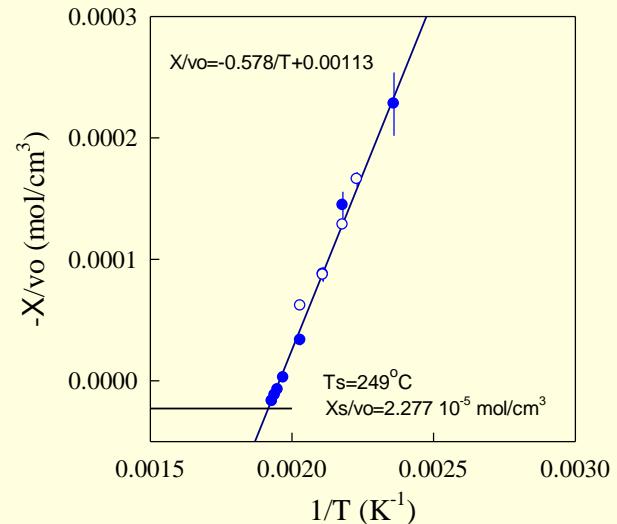
## 1-phase scattering



## Orstein-Zernike



## Interaction $\chi$ FH



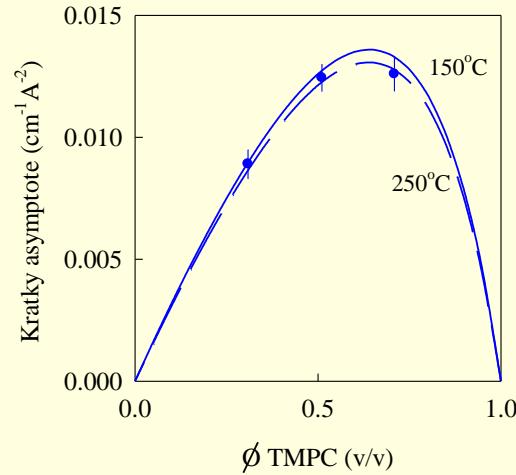
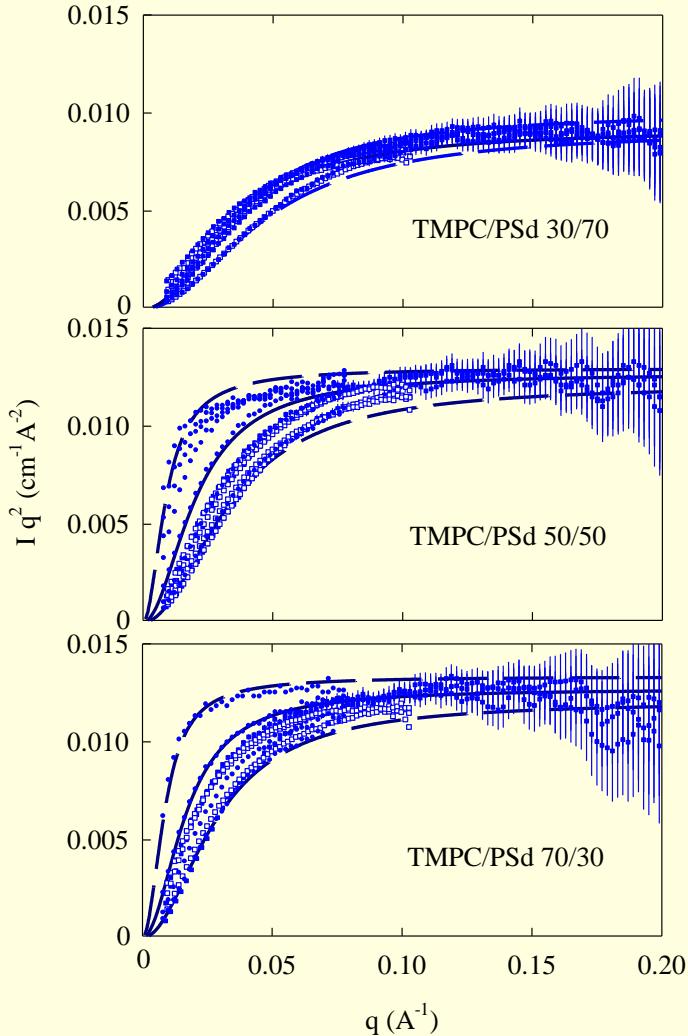
## Random Phase Approximation

$$\frac{1}{S_{AB}(q)} = \frac{1}{\phi_A v_A \langle N_A \rangle_n \langle g_D(q)_A \rangle_w} + \frac{1}{\phi_B v_B \langle N_B \rangle_n \langle g_D(q)_B \rangle_w} - 2 \frac{\tilde{\chi}_{AB}}{v_o}$$

$$\frac{1}{S(q)} = 2(\chi_s - \chi_f) + \frac{\xi^2}{S(0)} q^2$$

$$\xi^2 = \frac{v_0}{36(\tilde{\chi}_s - \tilde{\chi}_{AB})} \left( \frac{\langle N_A \rangle_z}{\langle N_A \rangle_w} \frac{a_A^2}{\phi_A v_A} + \frac{\langle N_B \rangle_z}{\langle N_B \rangle_w} \frac{a_B^2}{\phi_B v_B} \right)$$

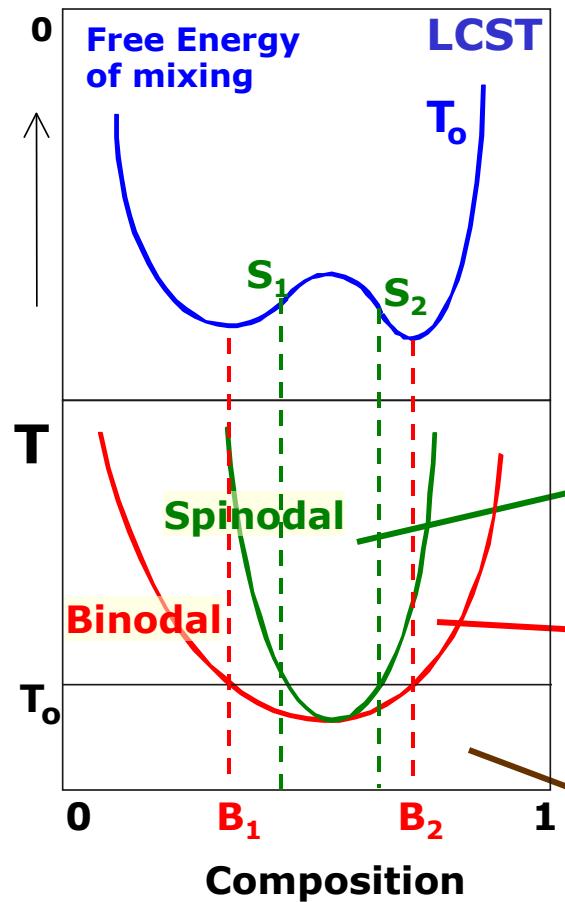
# Equilibrium: Kratky asymptote



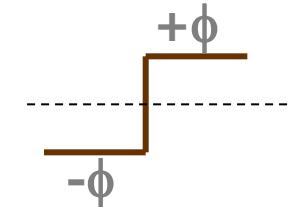
**Kratky asymptote and segment length**

$$S(q) \approx \frac{12\phi_1\phi_2}{q^2} \frac{v_o}{\hat{a}^2} \quad \frac{\hat{a}^2}{v_0} \equiv \phi_1\phi_2 \left( \frac{a_1^2}{\phi_1 v_1} + \frac{a_2^2}{\phi_2 v_2} \right)$$

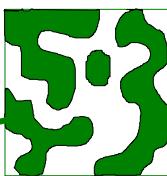
# Non-equilibrium: Fluctuations & Phase separation



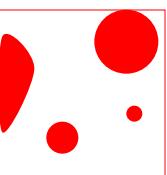
Concentration fluctuations



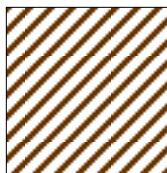
**Unstable:**  $G'' < 0$   
*spinodal decomposition*



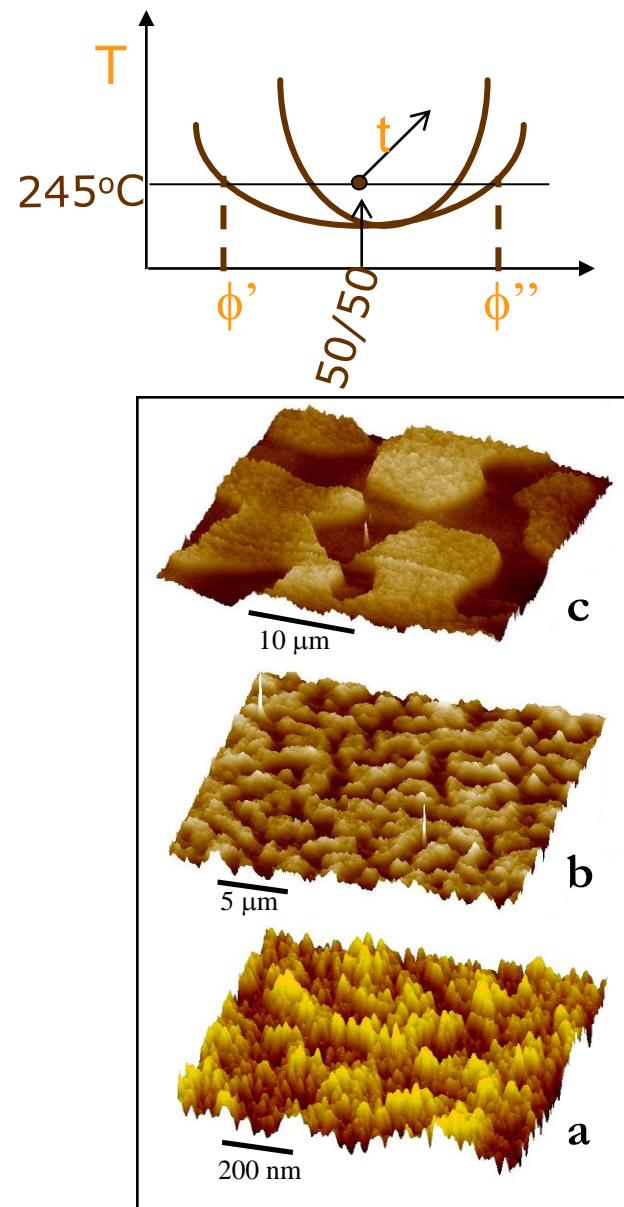
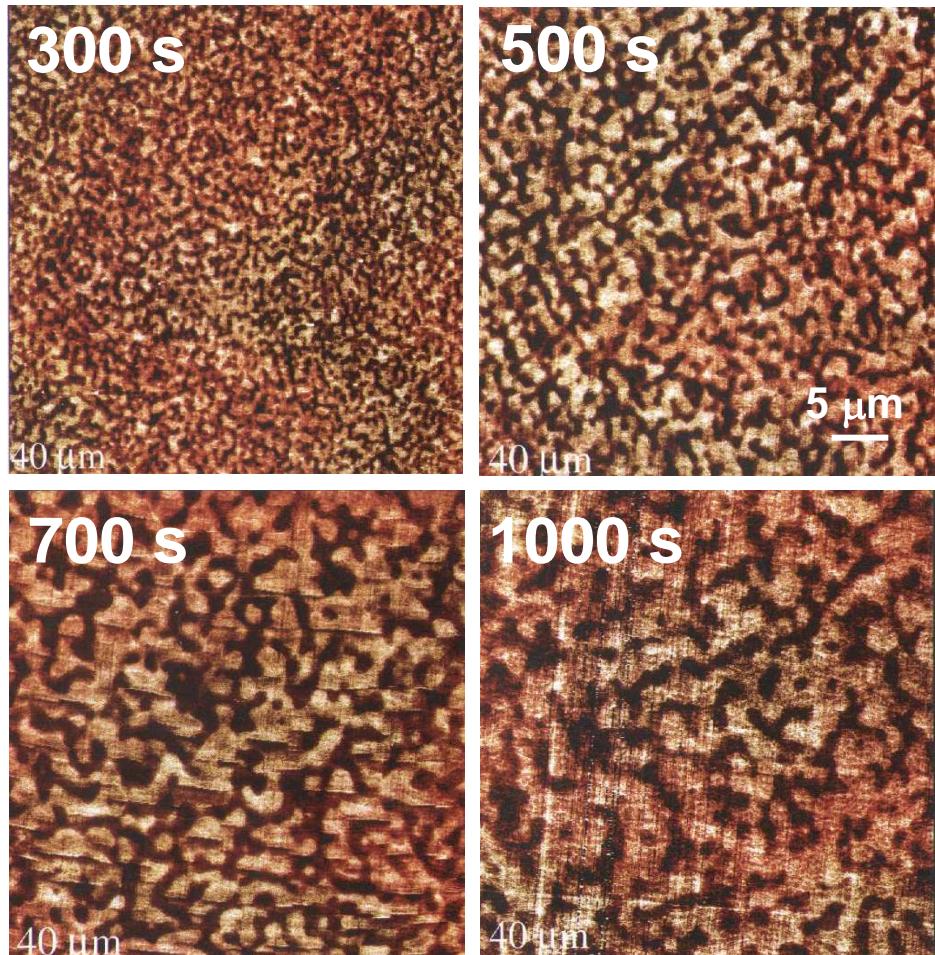
**Metastable:**  
*nucleation & growth*



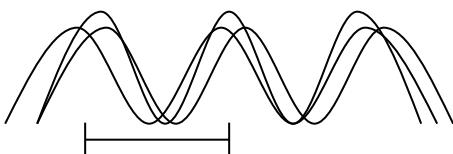
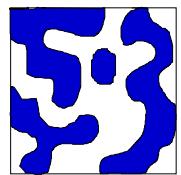
**Stable:**  
*equilibrium*



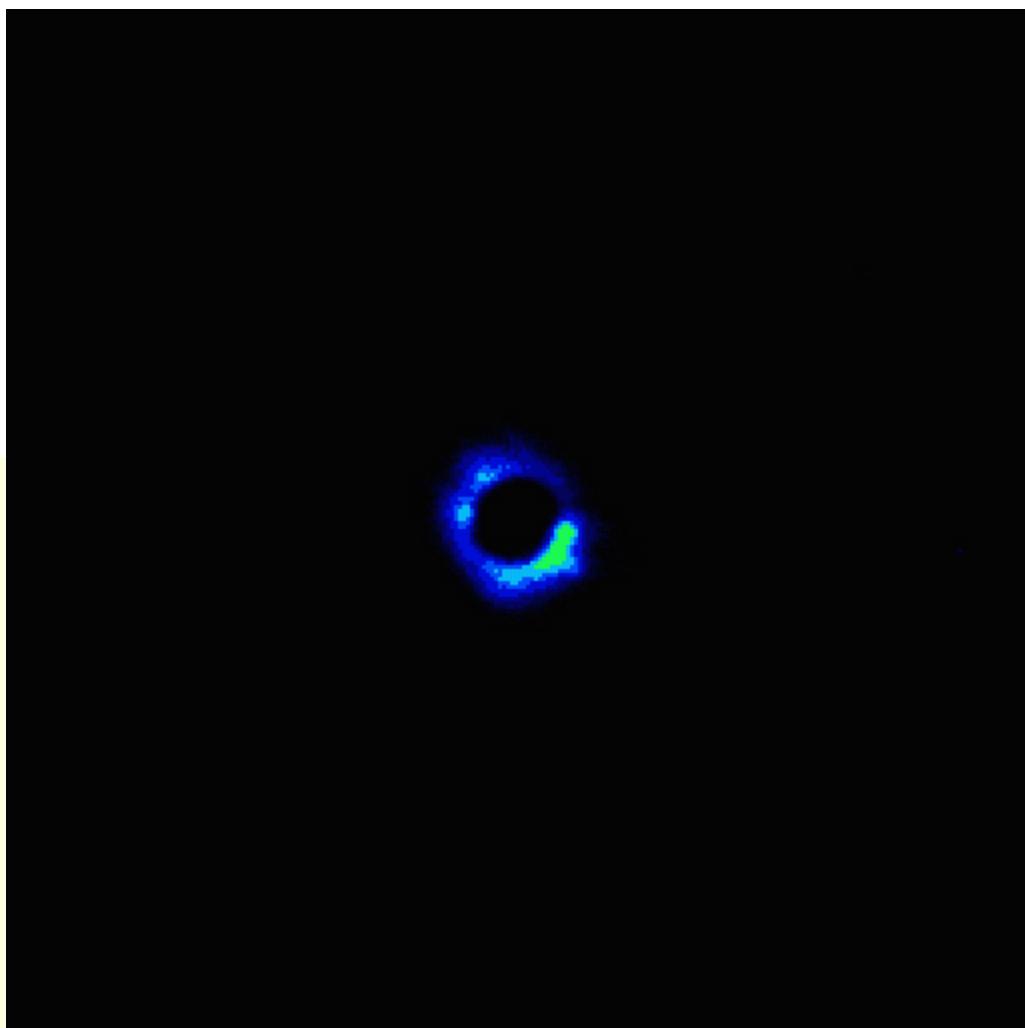
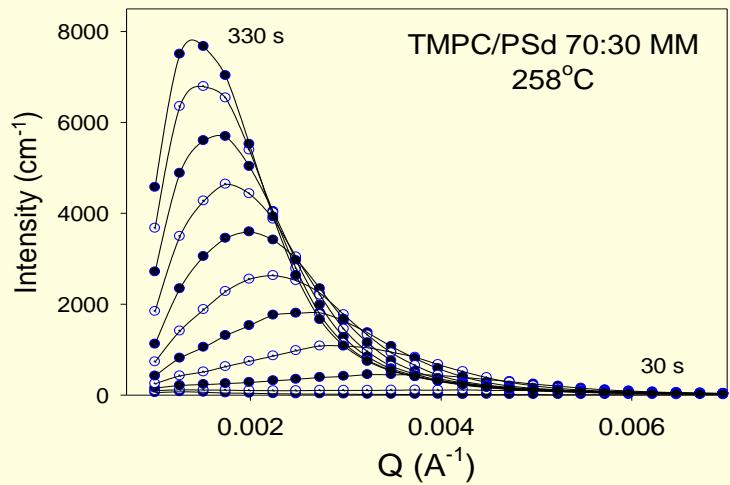
# Phase separation: spinodal decomposition



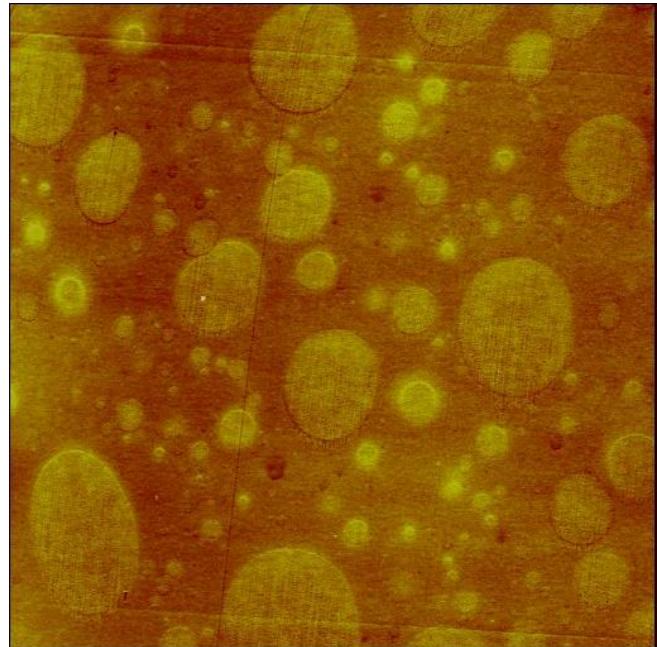
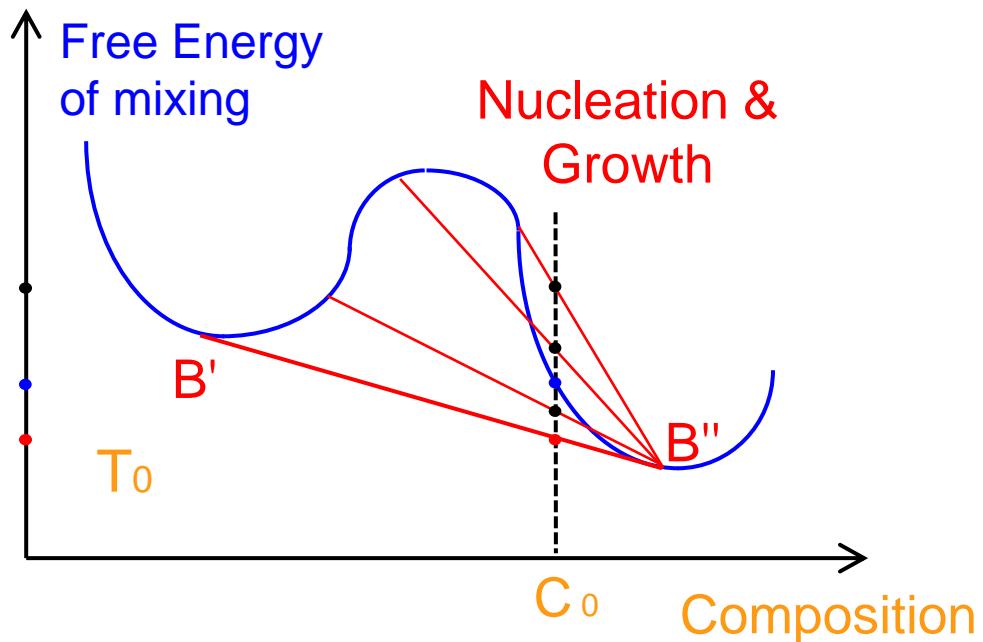
# Phase separation



$\Lambda_m$ : characteristic length of phase separation  $\sim 10\text{s}-100\text{s nm}$

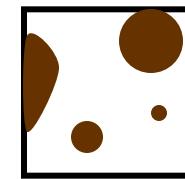
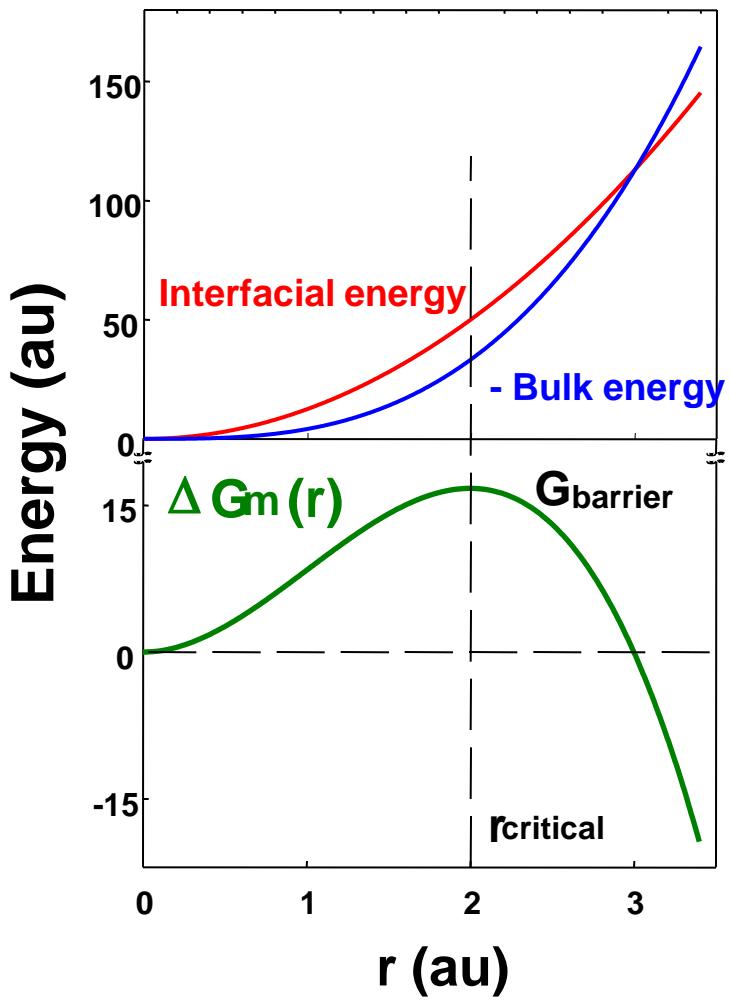


# Nucleation & Growth



Comparatively SLOW,  
since activated process

# Nucleation & Growth



## Energy balance

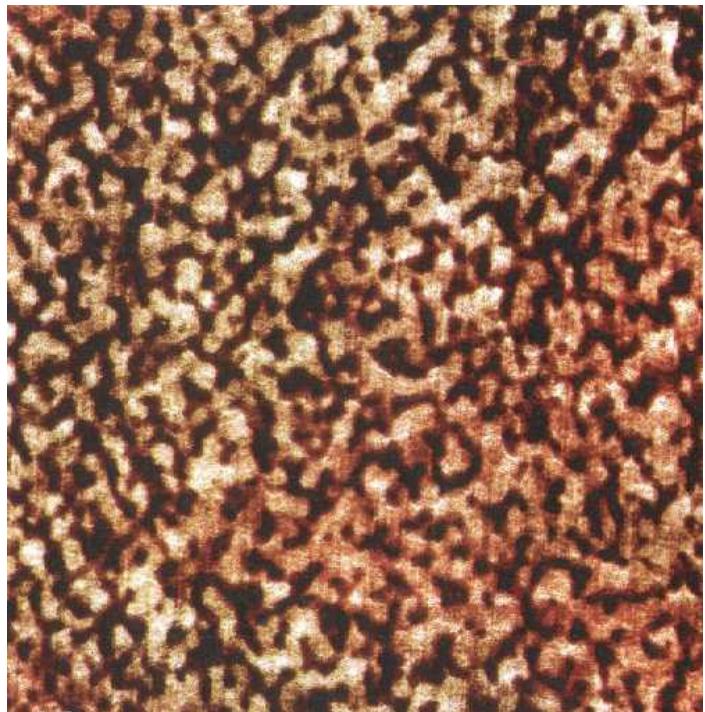
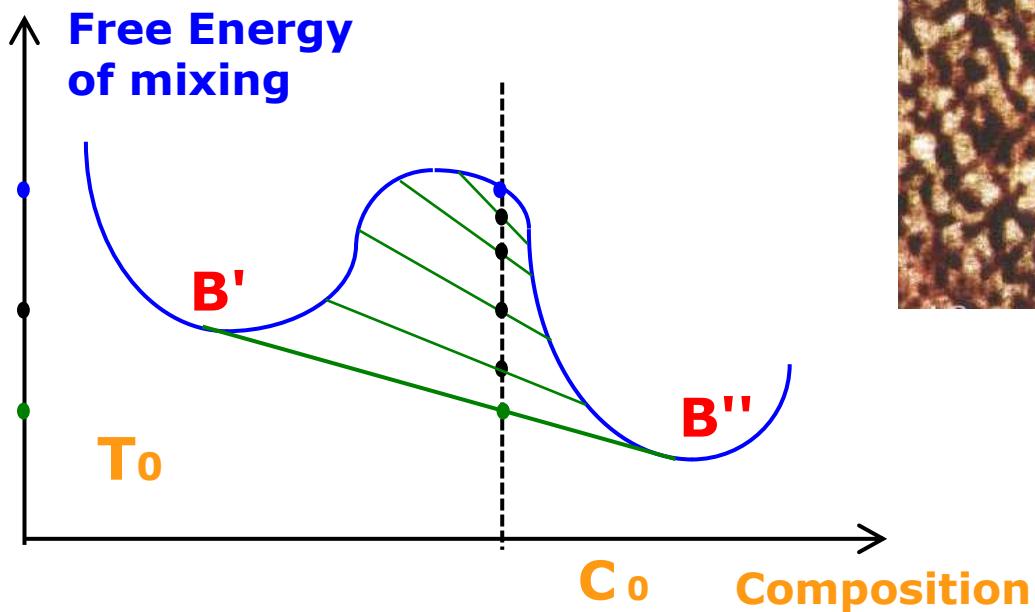
$$\Delta G(r) = -\frac{4\pi}{3} r^3 \Delta g + 4\pi r^2 \sigma$$

**bulk   interface**

→ Critical droplet

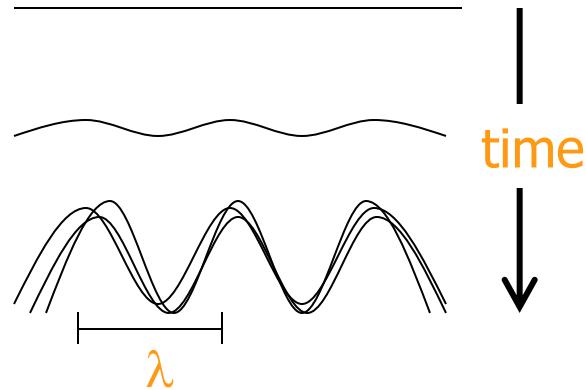
$$\Delta G_{barrier} = \frac{16\pi\sigma^3}{3\Delta g^2}$$
$$r_c = 2\sigma/\Delta g$$

# Spinodal decomposition



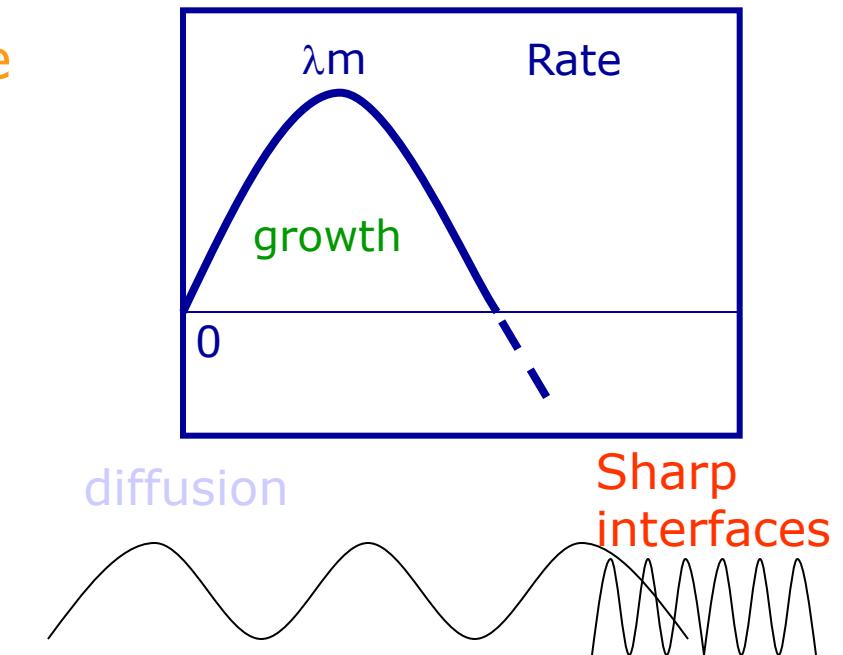
FAST, spontaneous

# Spinodal decomposition



concentration fluctuations

$$\phi = \phi_o + \tilde{\phi}$$



$$Rate(\lambda) = -A \frac{\partial^2 G}{\partial \Phi^2} (1/\lambda)^2 - B(1/\lambda)^4$$

# Cahn-Hilliard Cook theory

$$S(q, t) = S_T(q) + [S(q, 0) - S_T(q)] \exp^{-2\tau_q^{-1}t}$$

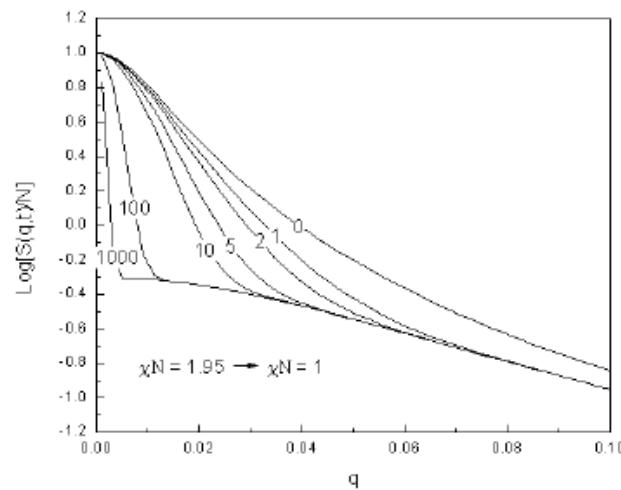
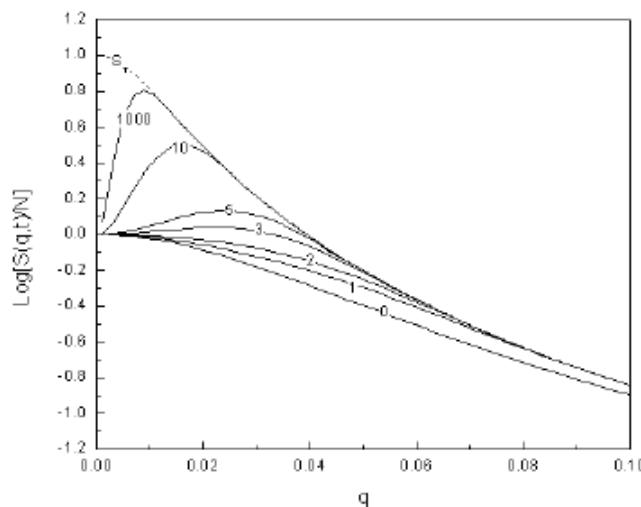
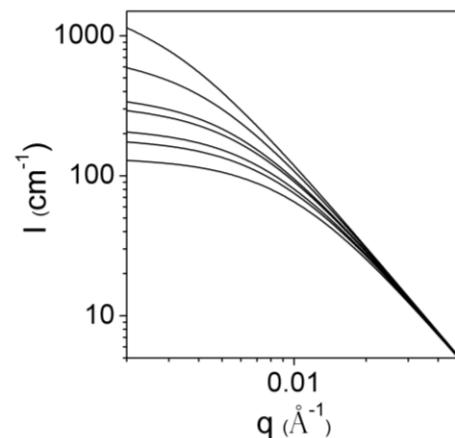
$$\tau_q^{-1} = q^2 \Lambda(q) S_T^{-1}(q)$$

$$\Lambda(q) = Wa^2 \phi(1 - \phi) g_D(q)$$

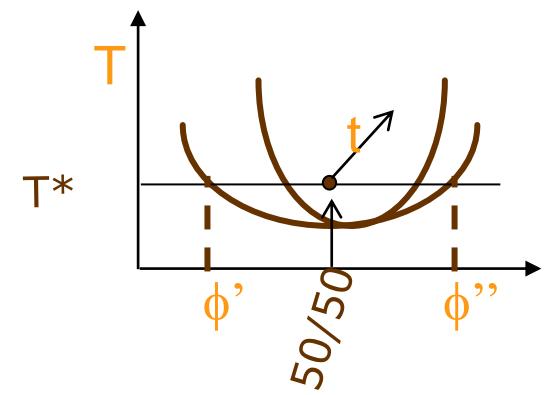
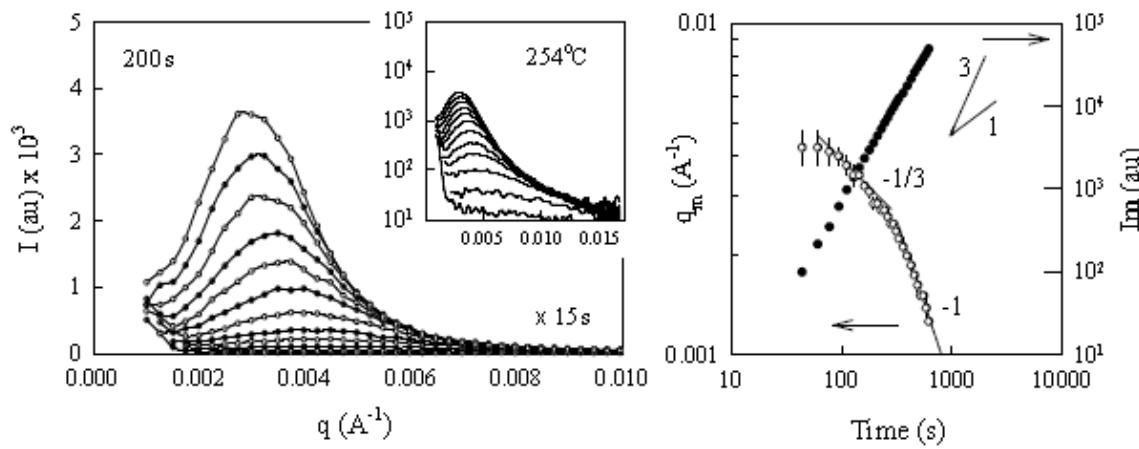
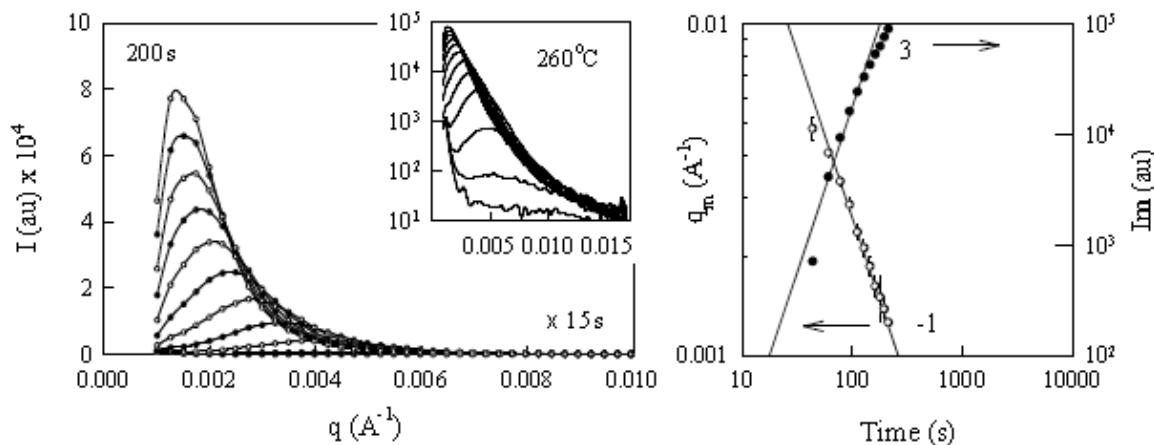
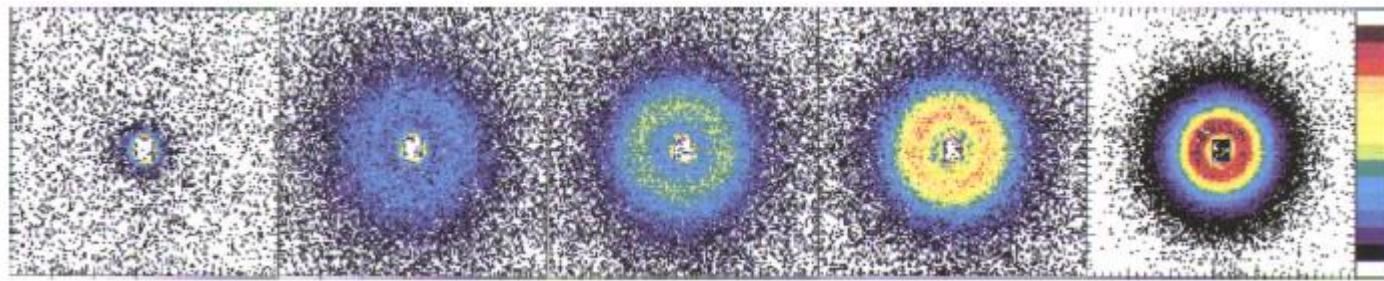
(Rouse Model)

$$\Lambda(q) = \frac{Wd^2}{Z} \phi(1 - \phi) g_D(q)$$

(Reptation Model)



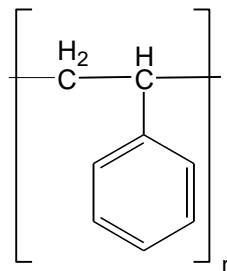
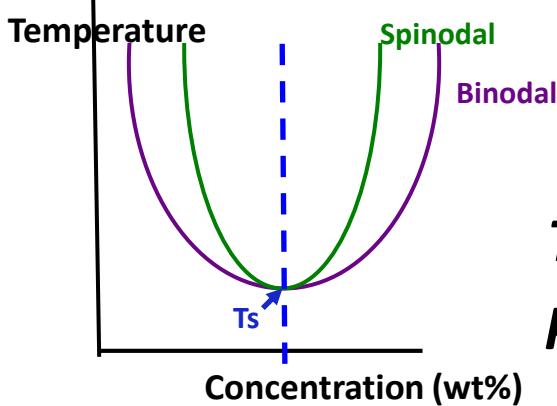
Spatially resolved “rheology”



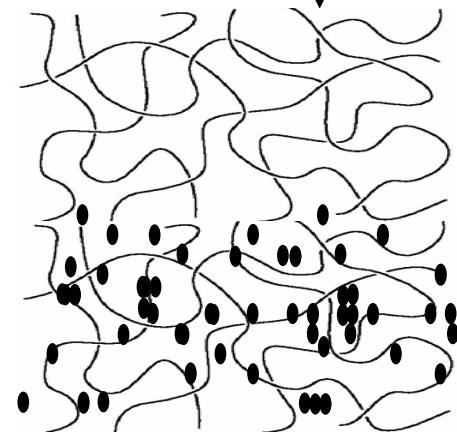
# Opportunities & recent developments

## nano*composites*

### structure



### dynamics

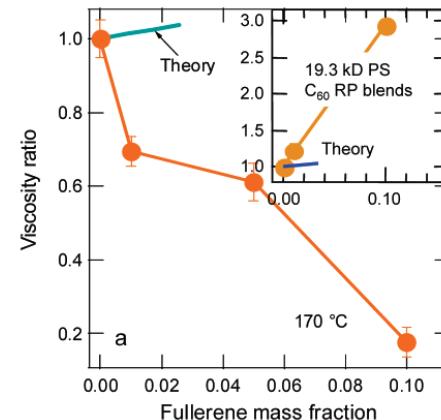
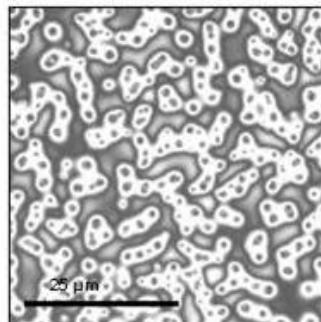


### Bulk

*Thermodynamics & phase separation*

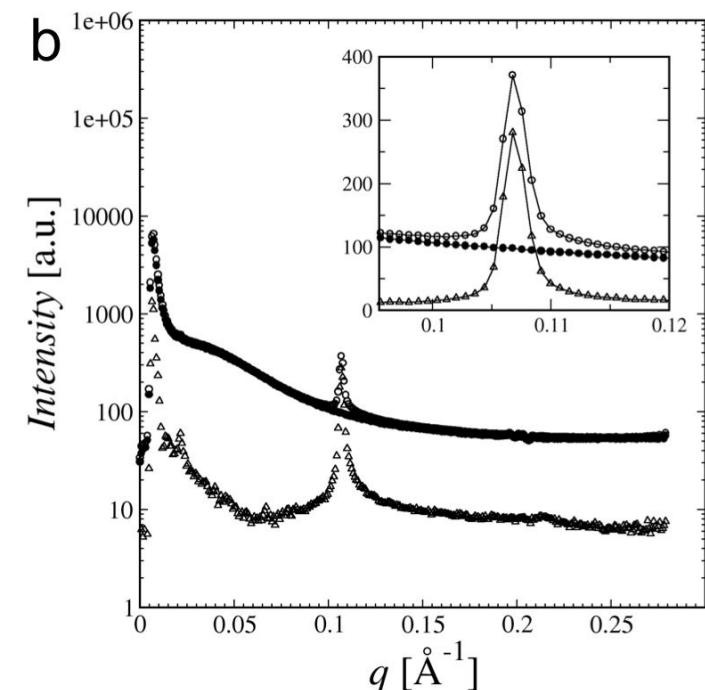
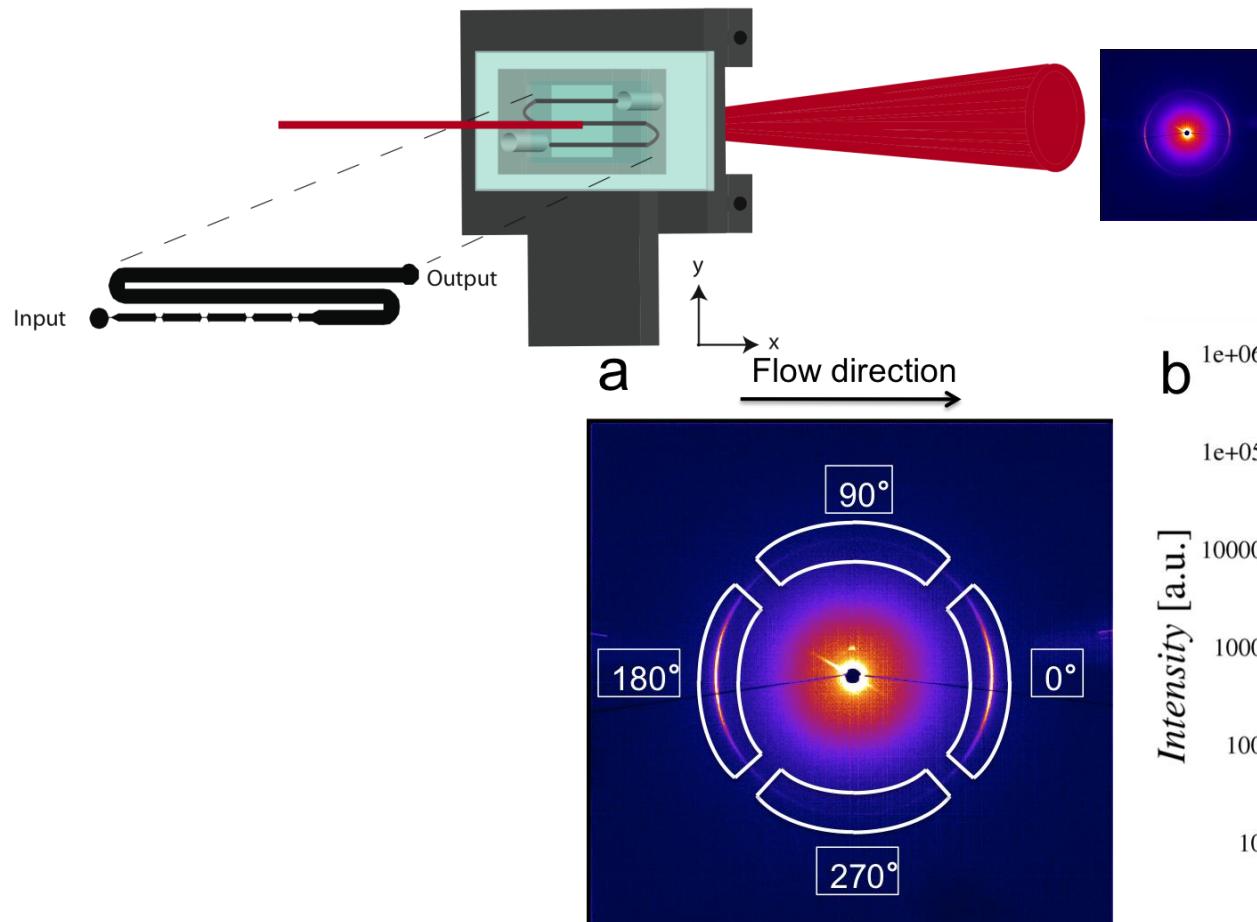
### Thin Films

### Morphology

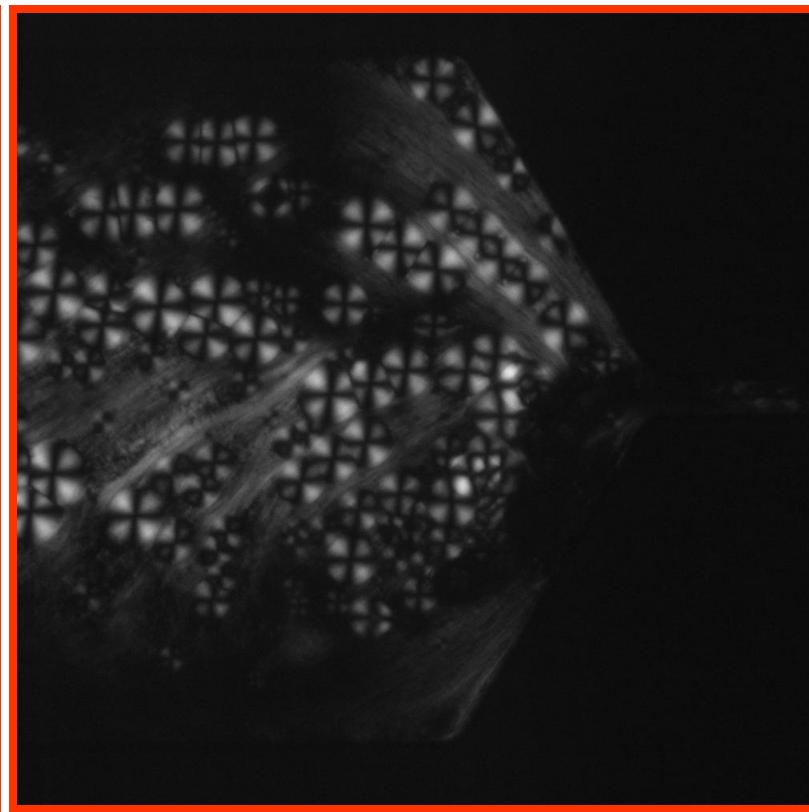
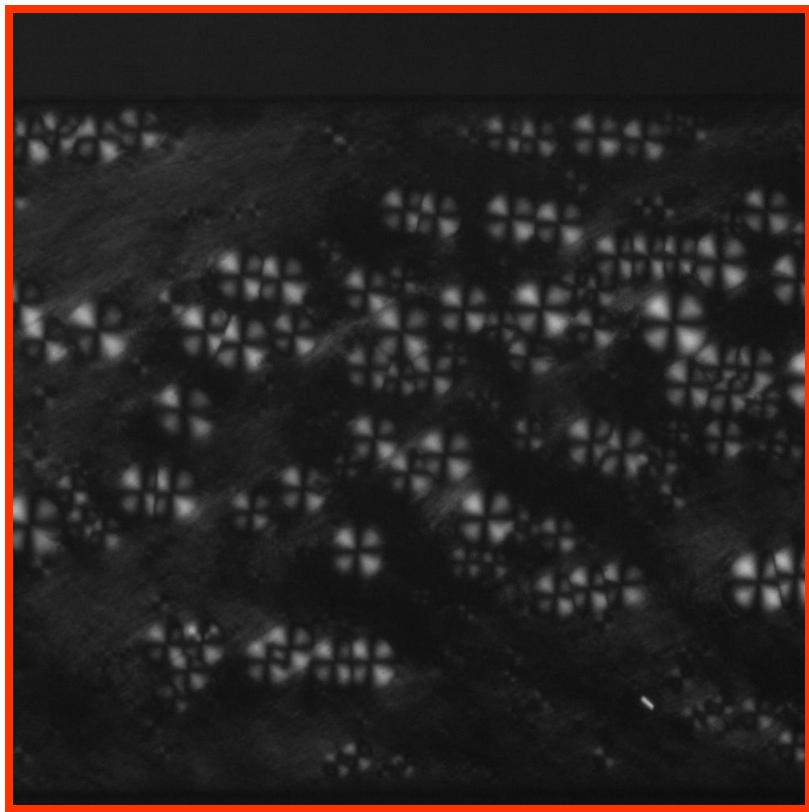


# Flow fields (and microfluidics?)

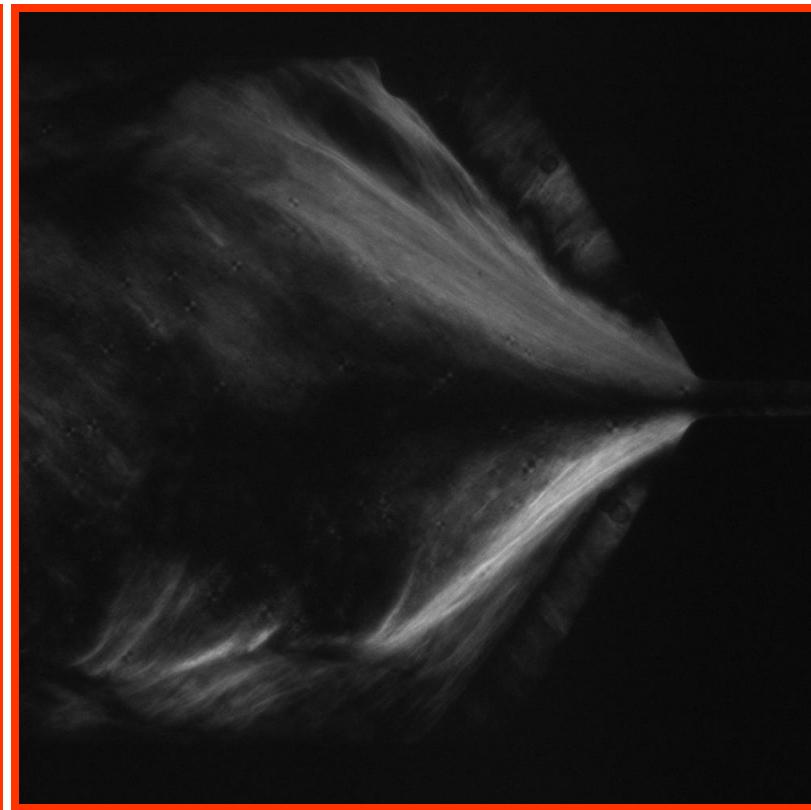
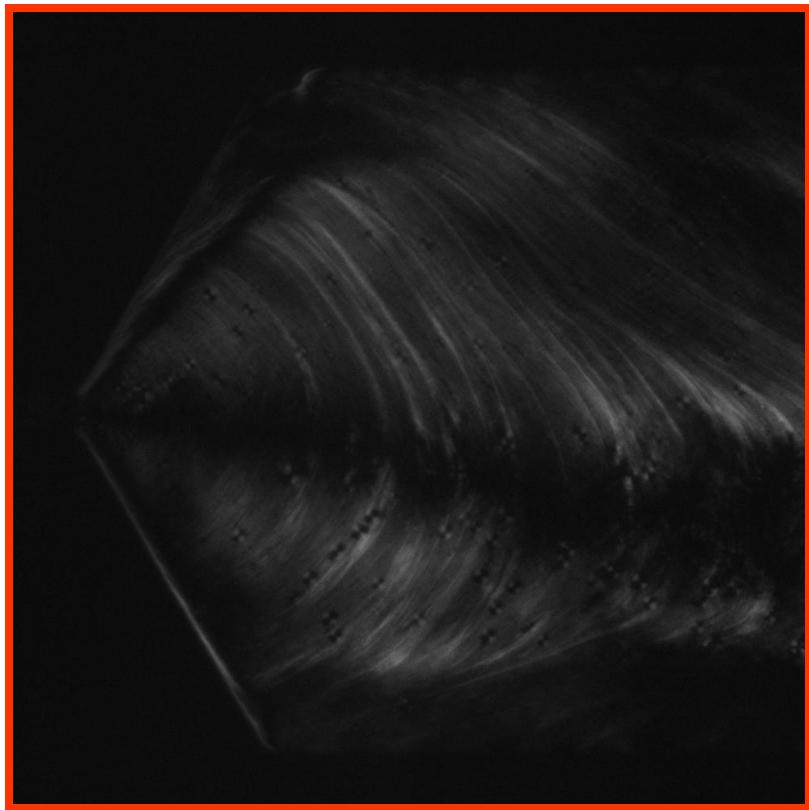
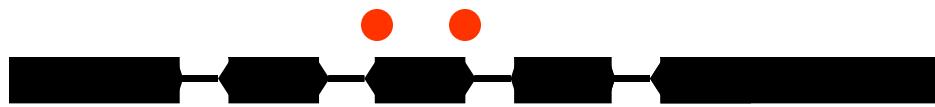
CTAC/ Pentanol/Water



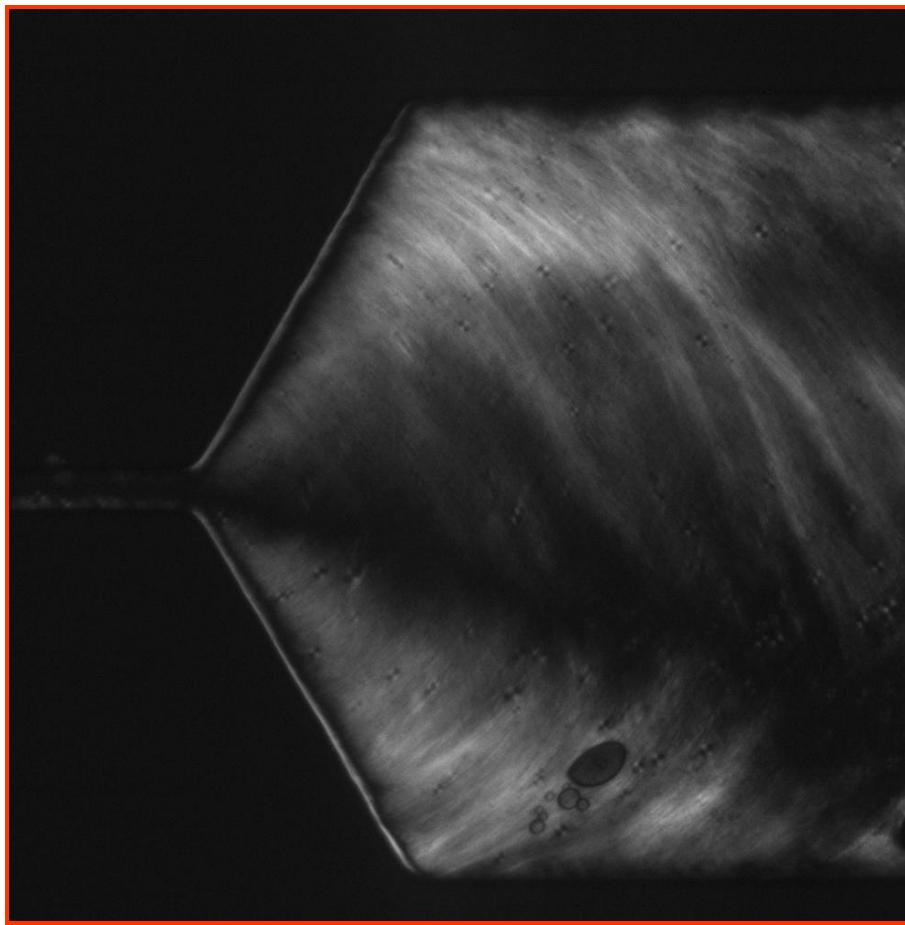
# Periodic constrictions: contraction/expansion flows



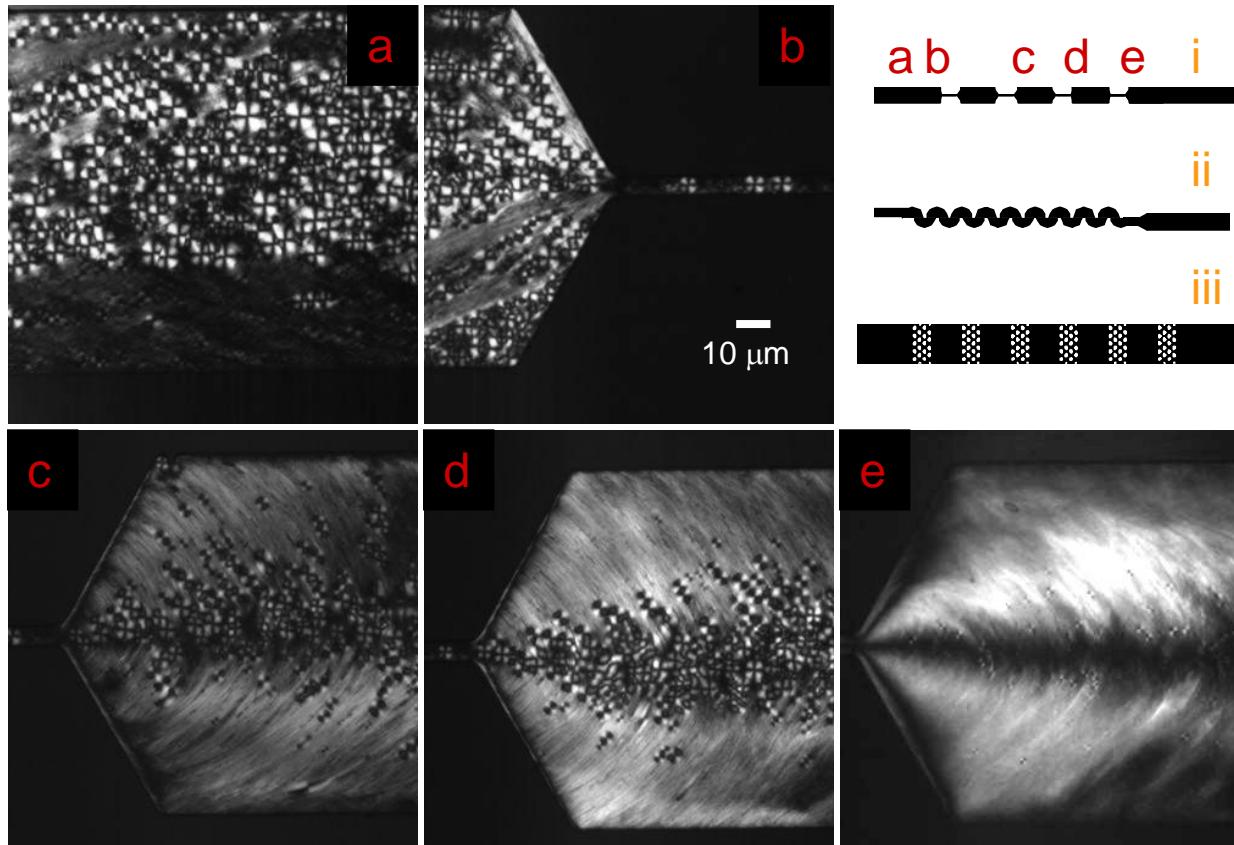
# Periodic constrictions: contraction/expansion flows



# Periodic constrictions: contraction/expansion flows

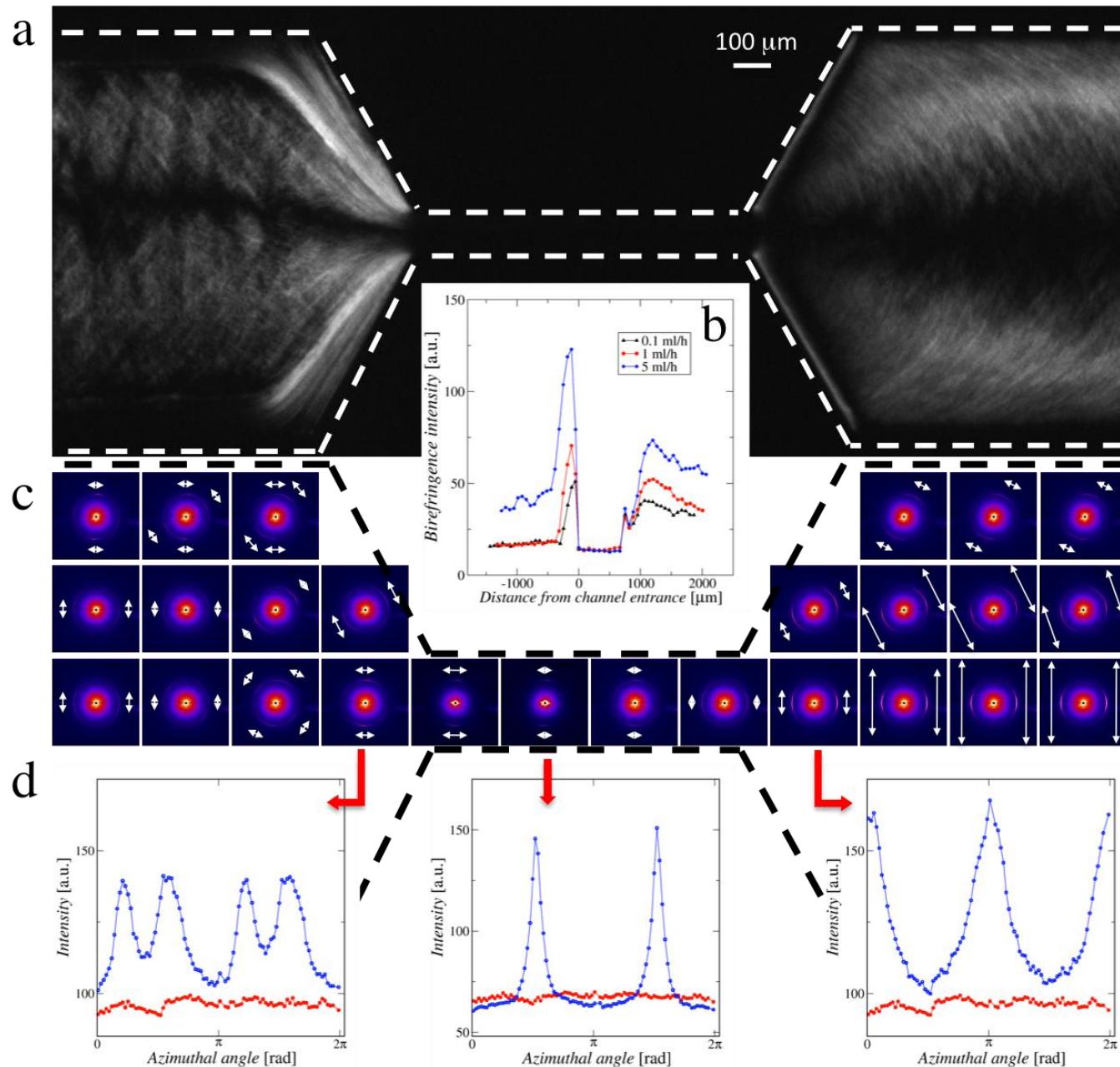


# Periodic constrictions: contraction/expansion flows



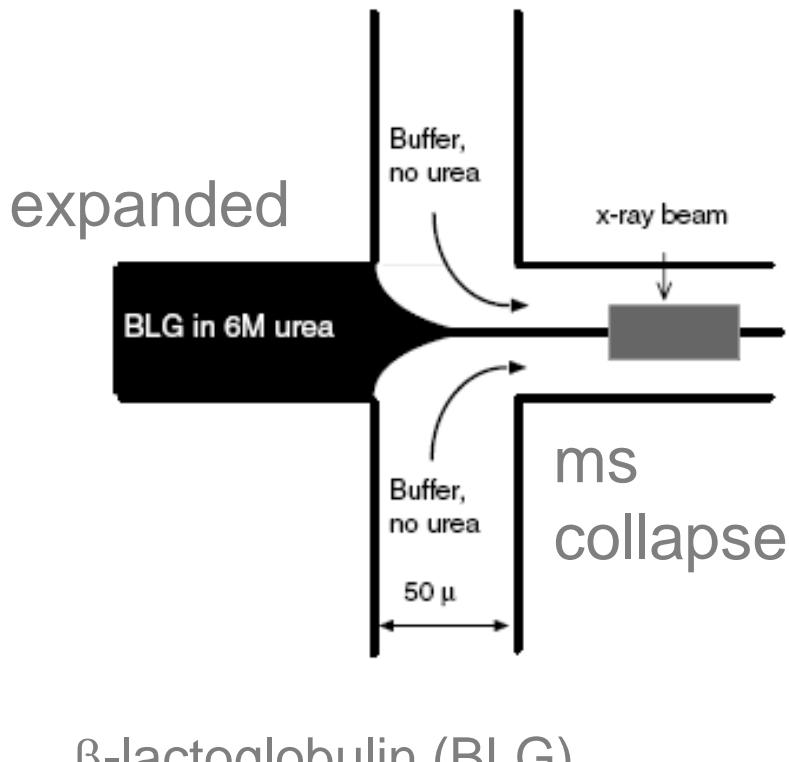
Optical micrographs: mixture of SDS / Octanol / Brine  
 $q = 1 \text{ ml/h}$ ; extension rate of  $10^4 \text{ s}^{-1}$ .

# Conclusion: Microstructure alignment & orientation

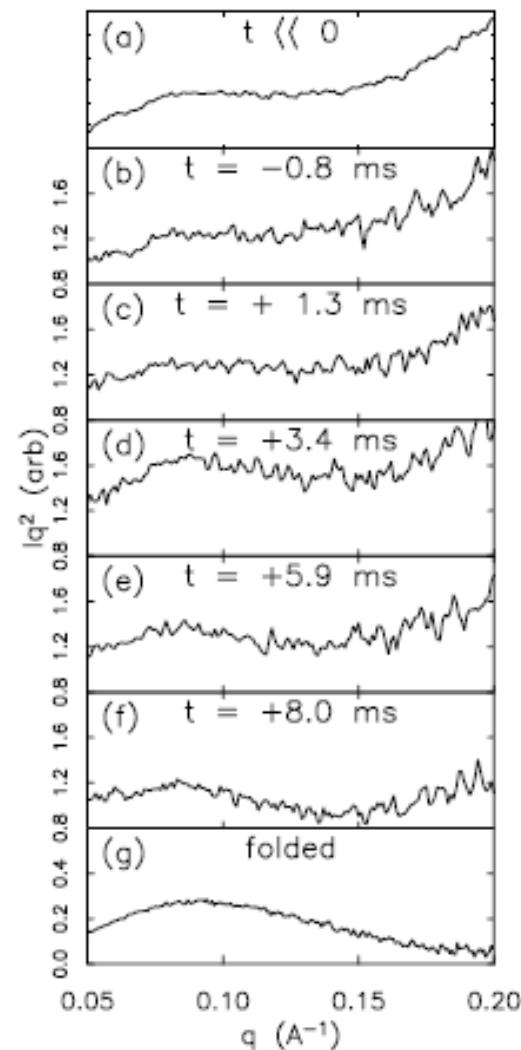


# Spatio-temporal mapping:

Coil-to-globule transition



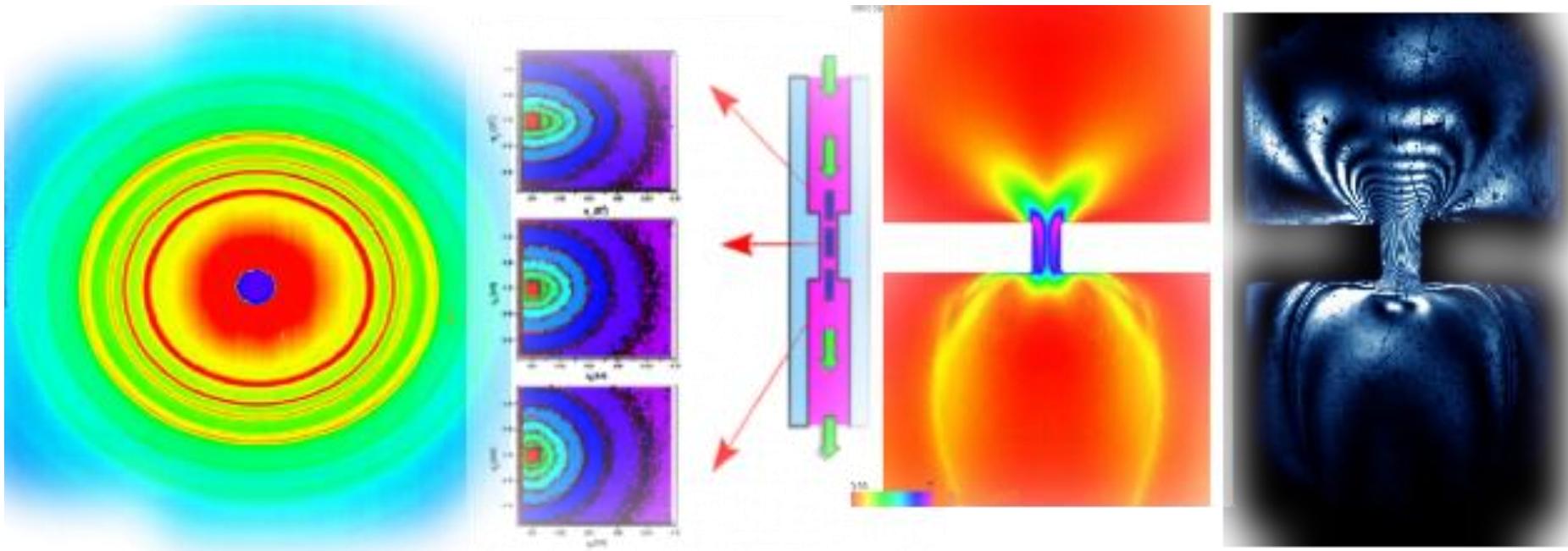
$\beta$ -lactoglobulin (BLG)



(Pollack, Austin, etc)

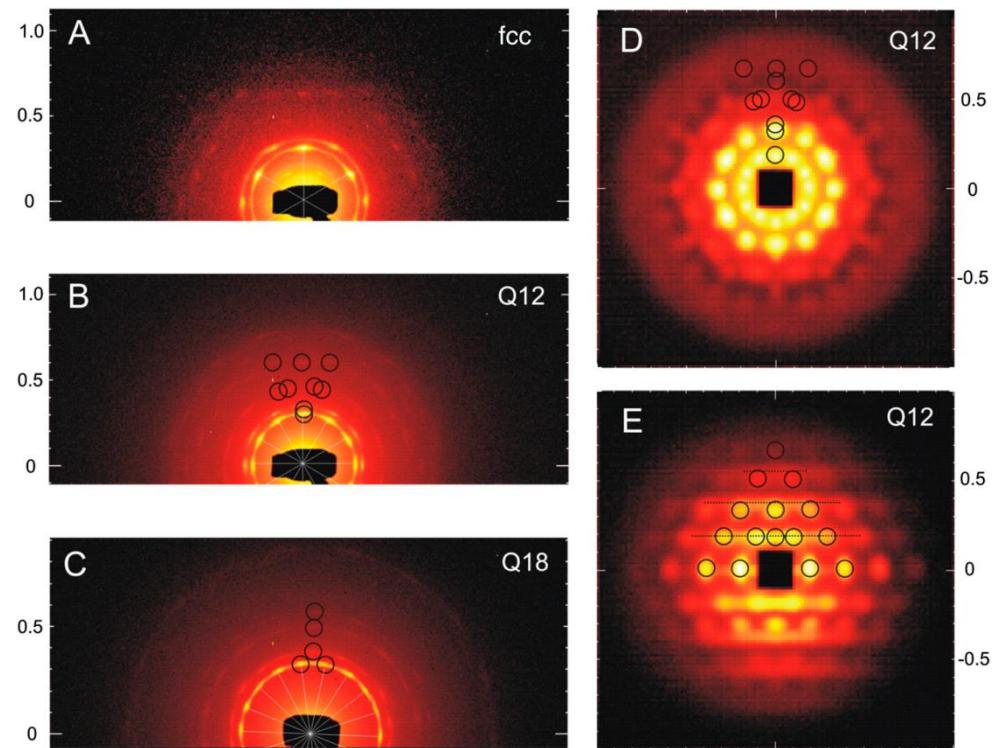
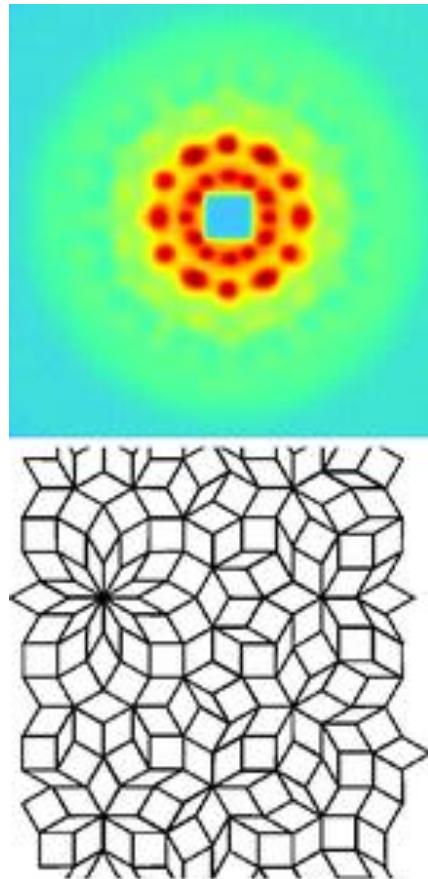
# Spatio-temporal mapping:

Entangled polymers under flow



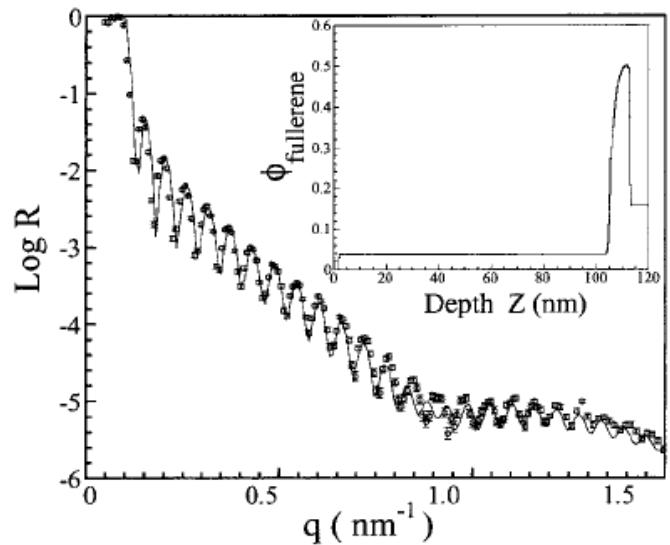
Neutron-Mapping Polymer Flow: Scattering, Flow Visualization, and Molecular Theory  
J. Bent, et al.  
*Science* 301, 1691 (2003)

# Soft colloids under flow

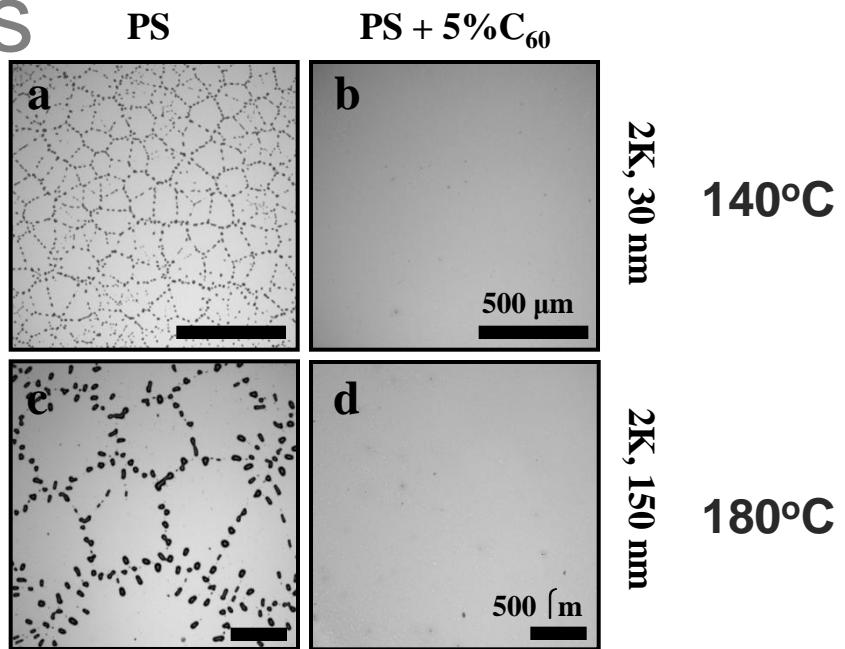


Forster et al. PNAS 2011

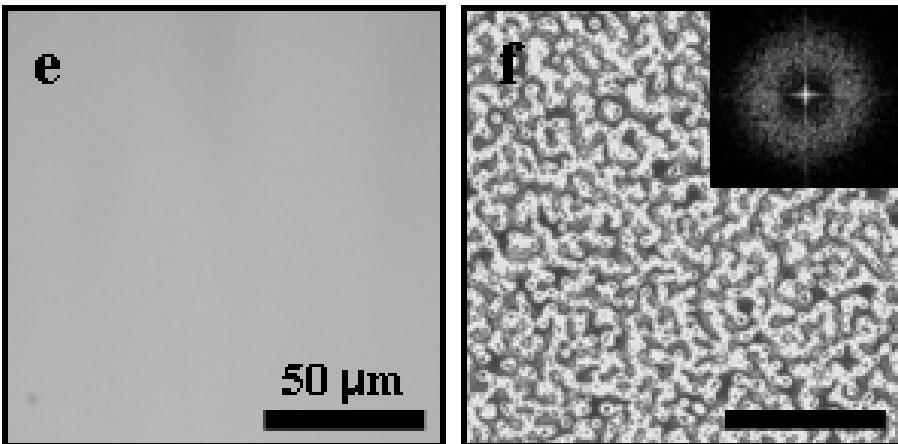
# Design 3D composites



**PS+5% C<sub>60</sub> h = 100nm**



270K, 150 nm  
180°C



**GISANS & reflecometry**

PRL 2010

**'Spinodal Clustering'**

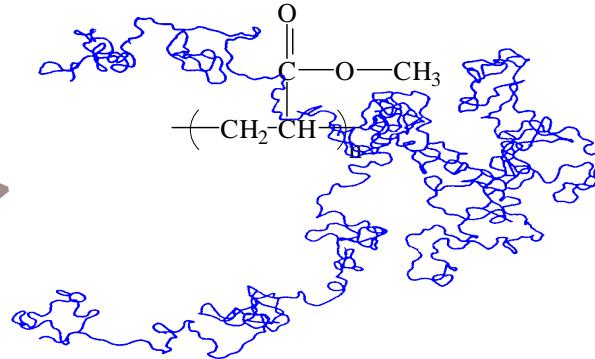
# Summary

## ① Intro



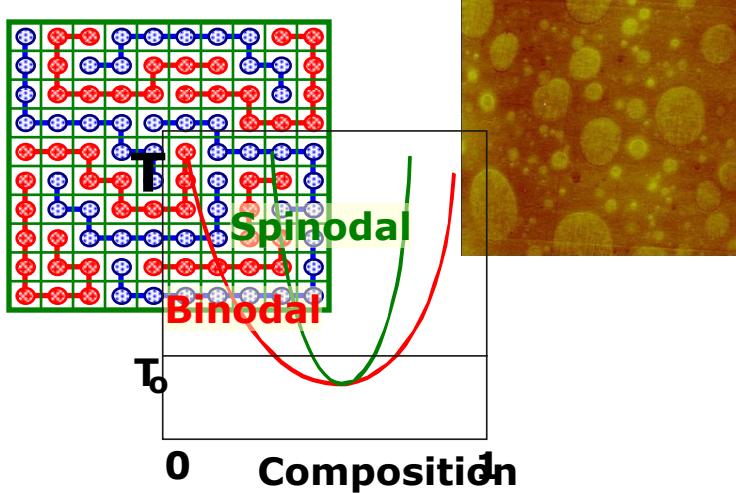
## History

## ② Soft matter



## ③ Form and structure

## ④ Mixtures & design



## ⑤ Outlook

