

Neutrons in soft matter

Lecture 2 – Reflectometry & Dynamics

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Outline

Lecture 1 – Structure & kinetics – SANS

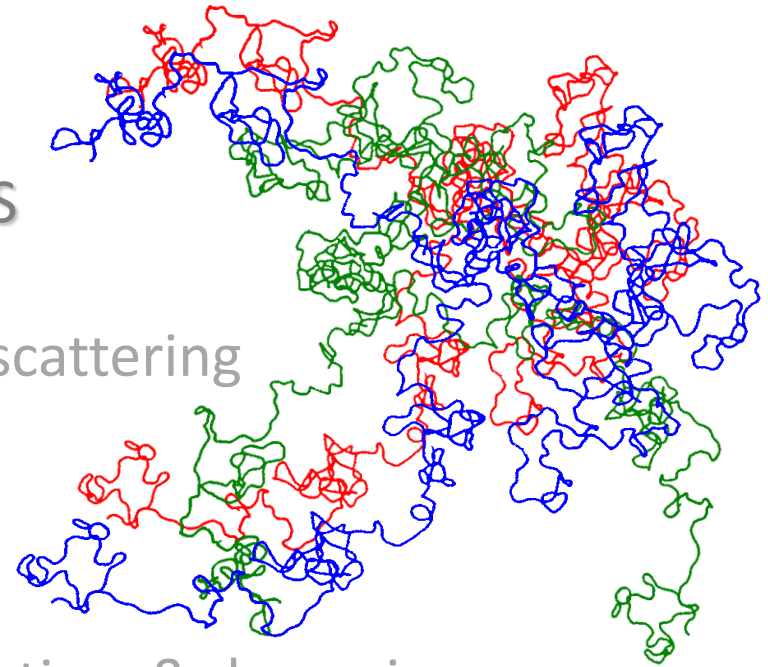
Introduction

soft matter & relevance of neutron scattering

Single objects: spheres, coils, rods...

Single chain polymer conformation
(solution and blends)

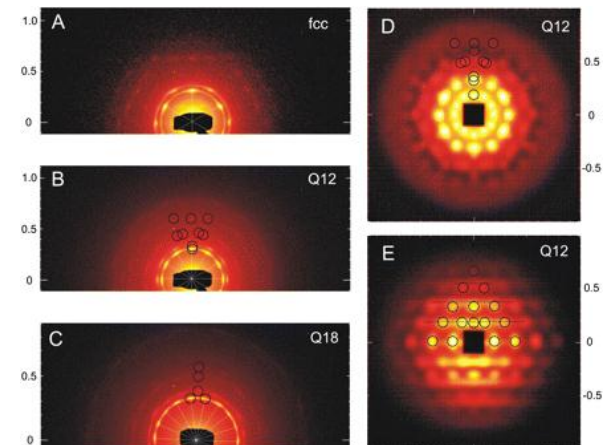
Polymer blends: interactions, conformation & dynamics
(equilibrium and phase separation)



Lecture 2 – Interfaces and dynamics

Reflectivity and diffusion

Dynamics in soft matter, QENS, BS, Spin-echo

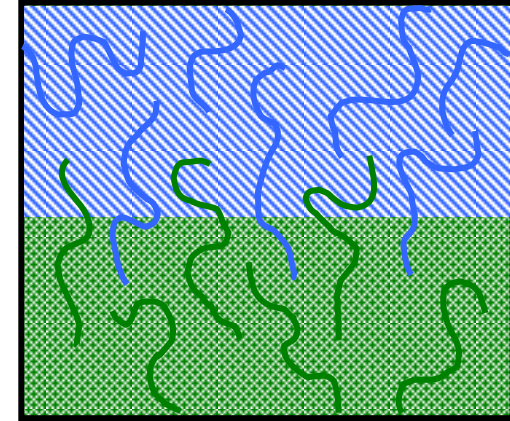


Forster et al (2011)

Reflectometry: study of interfaces

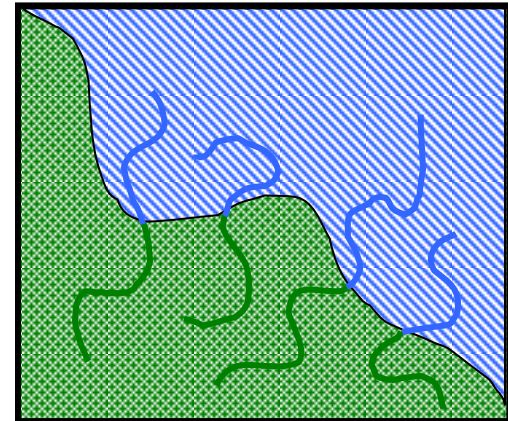
Miscible systems

- Interdiffusion, e.g., welding

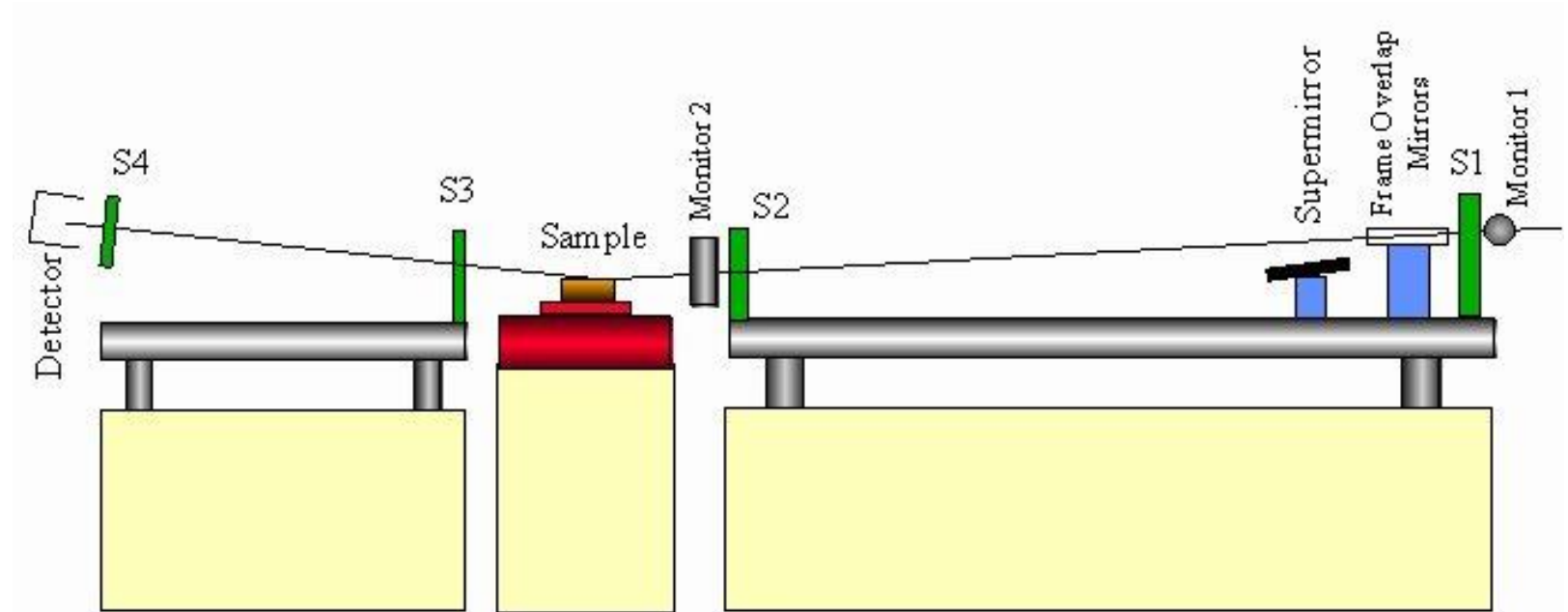


Immiscible systems

- Copolymers, e.g., di-blocks
- Reduce interfacial tension
→ smaller dispersed phase
- Entangle with homopolymers
→ **increase strength**



Reflectometry



CRISP (ISIS)



Significance of the interfacial width

Theoretical width

- Infinite molecular weight limit

$$w_t = \frac{2a}{(6\chi)^{0.5}}$$

E Helfand & AM Sapse
J Chem Phys 62 (1975) 1327

where a (statistical segment length)

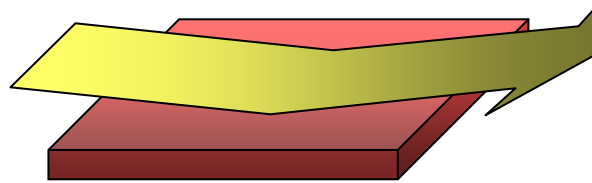
- Finite molecular weight limit

$$w_t = \frac{2a}{\sqrt{6}} \left(\chi - \frac{\pi^2}{6} (N_1^{-1} + N_2^{-1}) \right)^{-1/2}$$

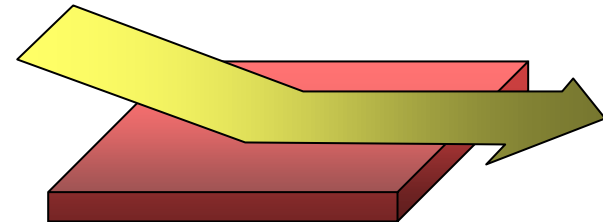
M Stamm & DW Schubert
Ann Rev Mater Sci
25 (1995) 325

⇒ Measure interfacial width to find χ

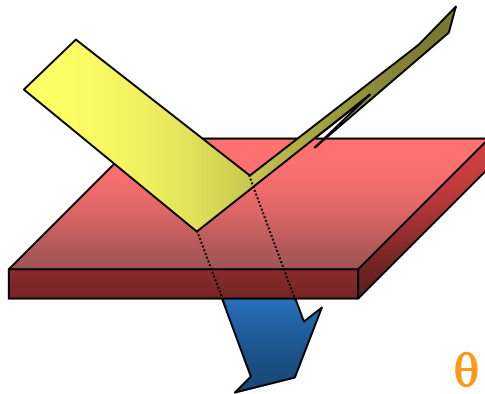
Basics of Reflectivity



$\theta < \theta_{\text{crit}}$
only reflection

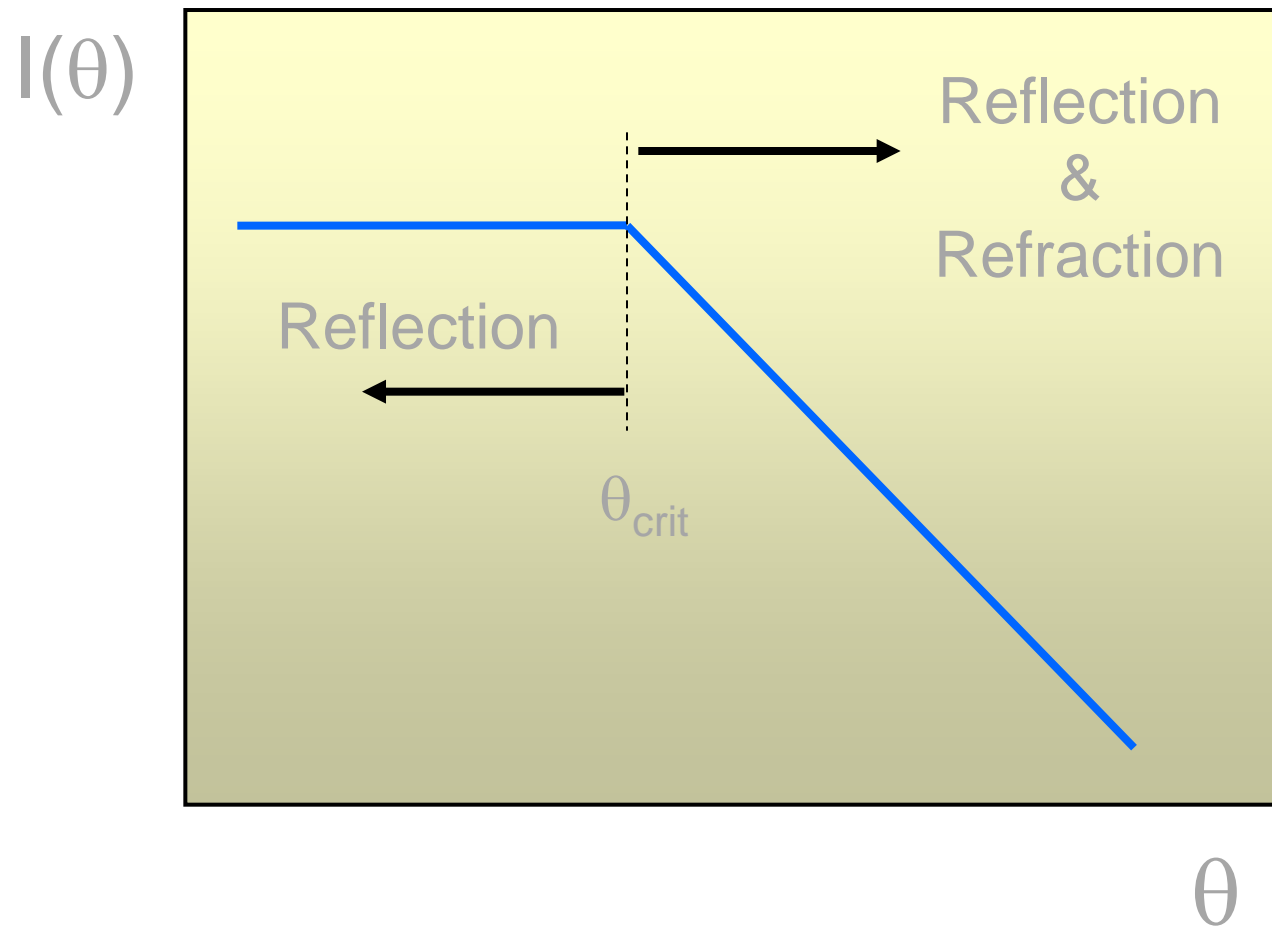


critical angle $\theta = \theta_{\text{crit}}$

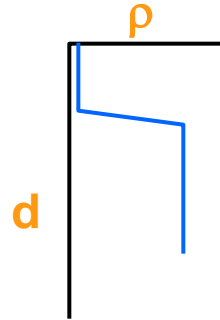
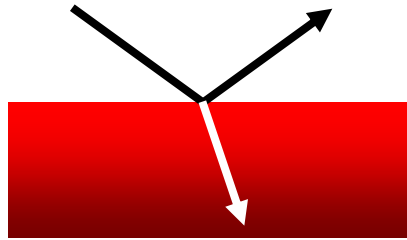


$\theta > \theta_{\text{crit}}$
reflection and refraction

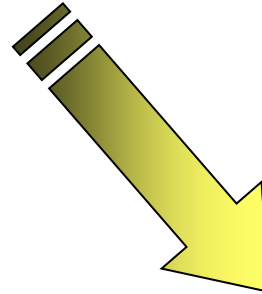
The Reflectivity Profile



Simplest Case



Information Content

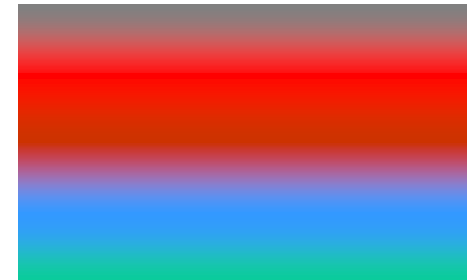
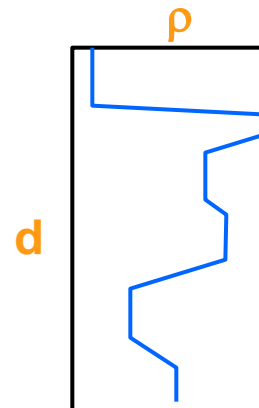


What about lateral information?

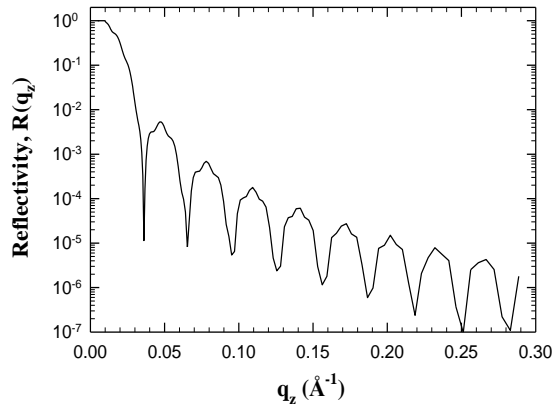


Off-specular !

Complex Case

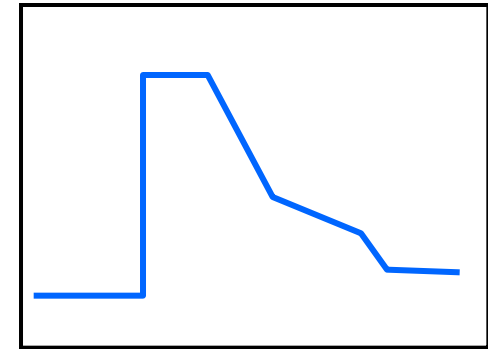


Evaluating Reflectivity Data



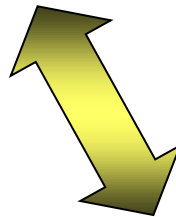
fft
➡
Ideal

ρ

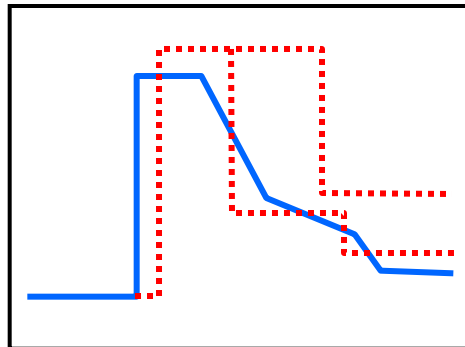


d

Real world

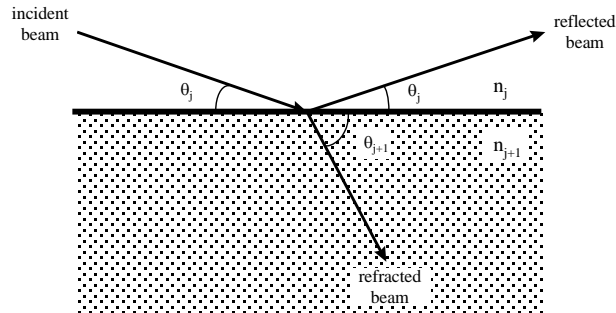


ρ



d

Single layers and bilayers



$$n_j = 1 - \frac{\lambda^2 N_d b}{2\pi} = 1 - \frac{\lambda^2 \rho_z}{2\pi}$$

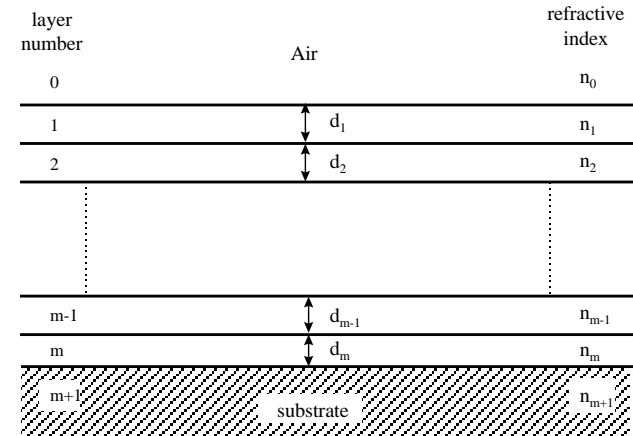
$$n_j \cos \theta_j = n_{j+1} \cos \theta_{j+1} \quad \text{Snell's Law}$$

$$r_{j,j+1} = \frac{n_j \sin \theta - n_{j+1} \sin \theta_{j+1}}{n_j \sin \theta + n_{j+1} \sin \theta_{j+1}} \quad \text{Fresnel's law}$$

$$q = 2k = \frac{4\pi}{\lambda} \sin \theta$$

$$r_{j,j+1} = \left(\frac{q_{z,j} - q_{z,j+1}}{q_{z,j} + q_{z,j+1}} \right)$$

$$R = r_{j,j+1} r_{j,j+1}^*$$



$$r'_{m-1,m} = \frac{r_{m-1,m} - r_{m,m+1} \exp(2i\beta_m)}{1 + r_{m-1,m} r_{m,m+1} \exp(2i\beta_m)}$$

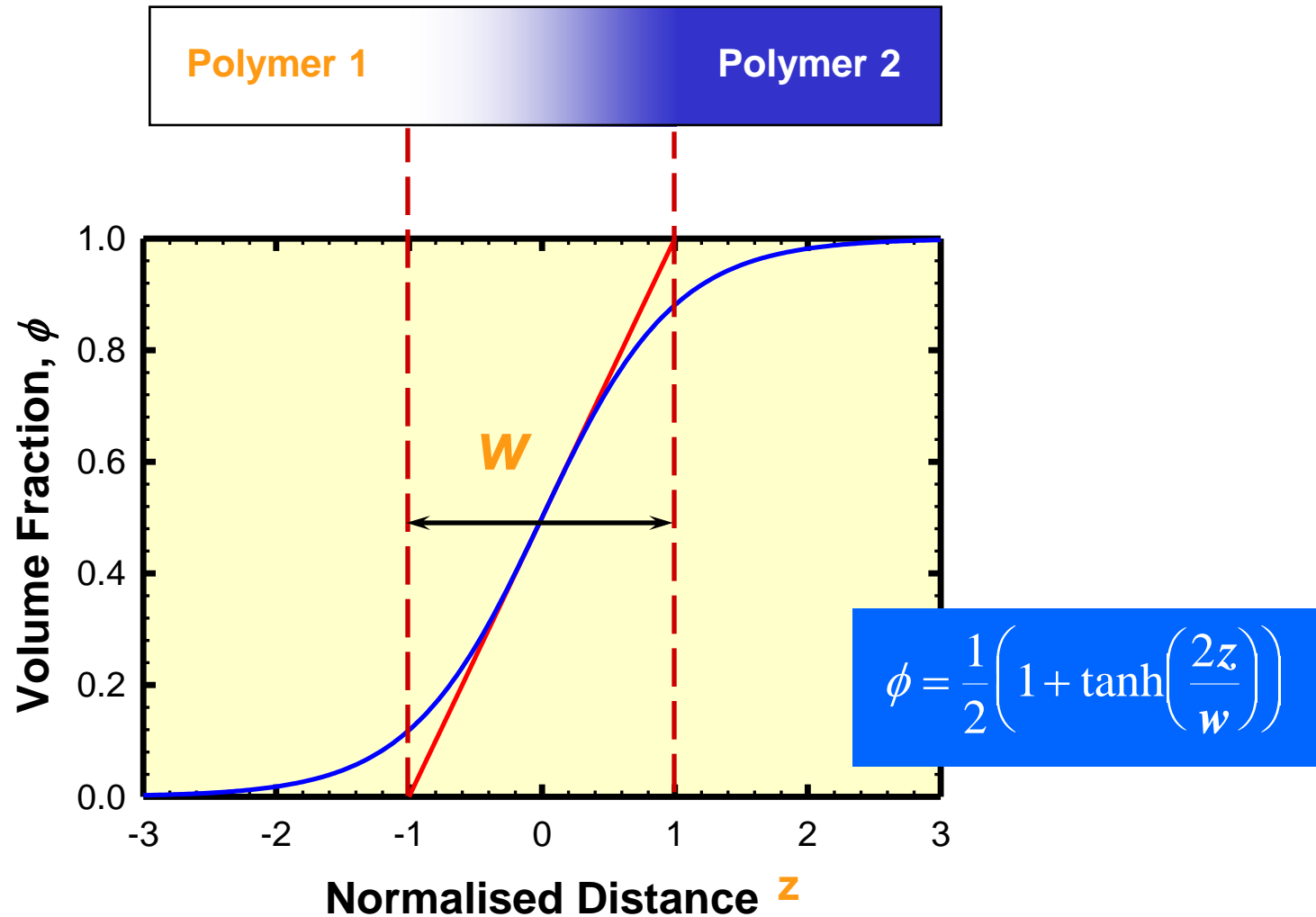
$$\beta_m = (2\pi/\lambda) n_m d_m \sin \theta$$

$$c_m = \begin{bmatrix} \cos \beta_m & -(i/\kappa_m) \sin \beta_m \\ -i\kappa_m \sin \beta_m & \cos \beta_m \end{bmatrix}$$

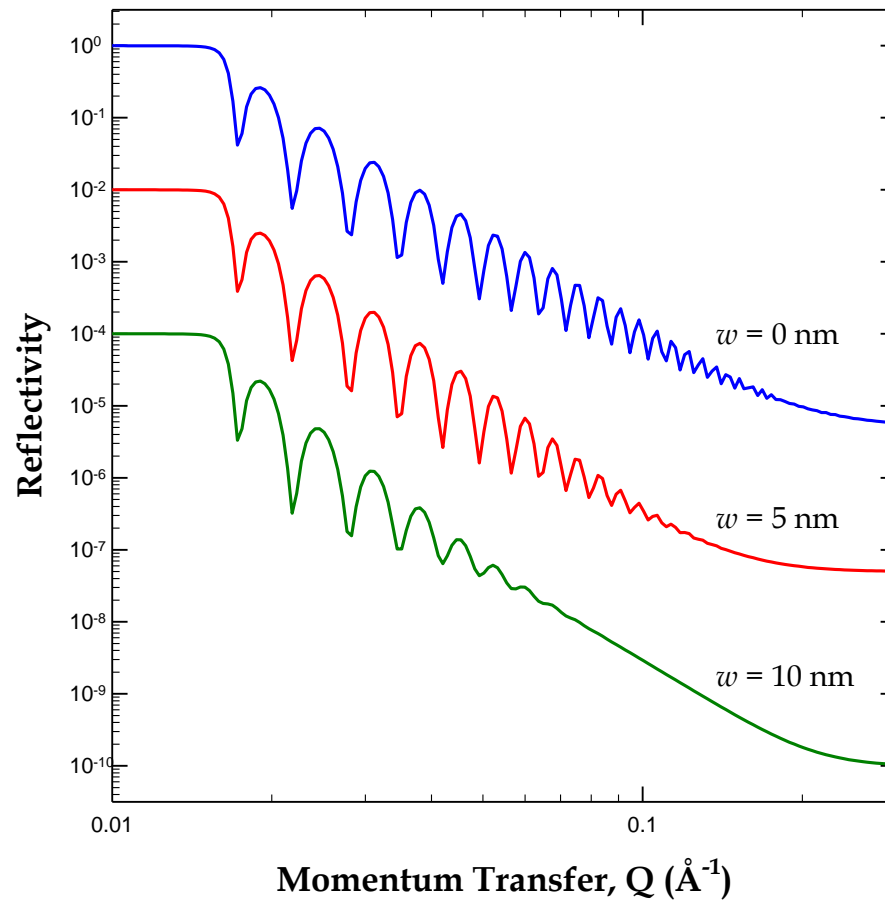
$$M = \prod_{m=0}^m c_m = \begin{bmatrix} M_{11} & M_{12} \\ M_{21} & M_{22} \end{bmatrix}$$

$$R = \left| \frac{(M_{11} + M_{12}\kappa_{m+1})\kappa_0 - (M_{21} + M_{22})\kappa_{m+1}}{(M_{11} + M_{12}\kappa_{m+1})\kappa_0 + (M_{21} + M_{22})\kappa_{m+1}} \right|^2$$

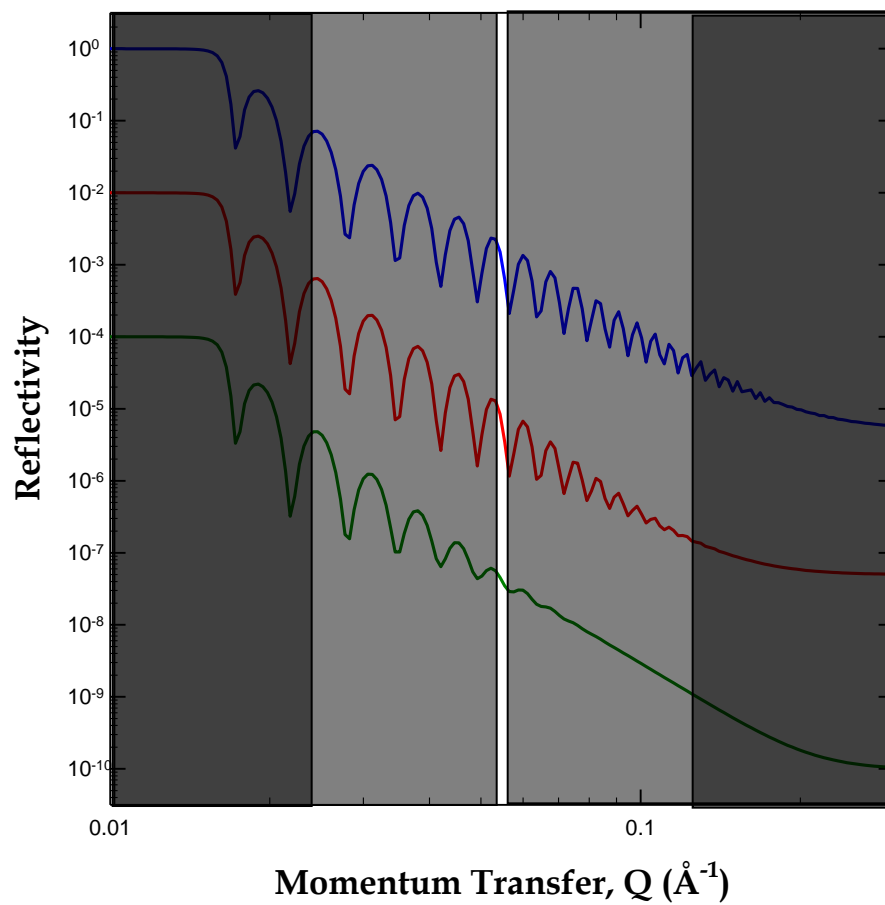
Interfacial Width - Definition



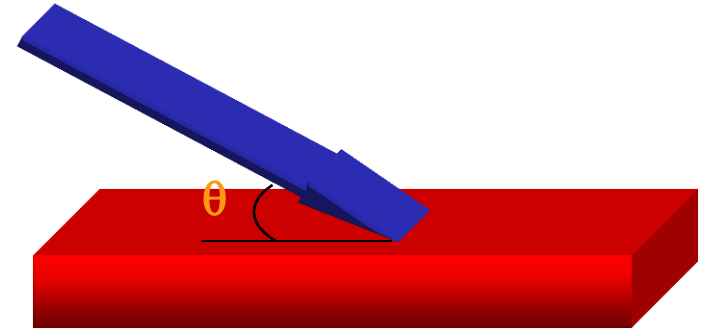
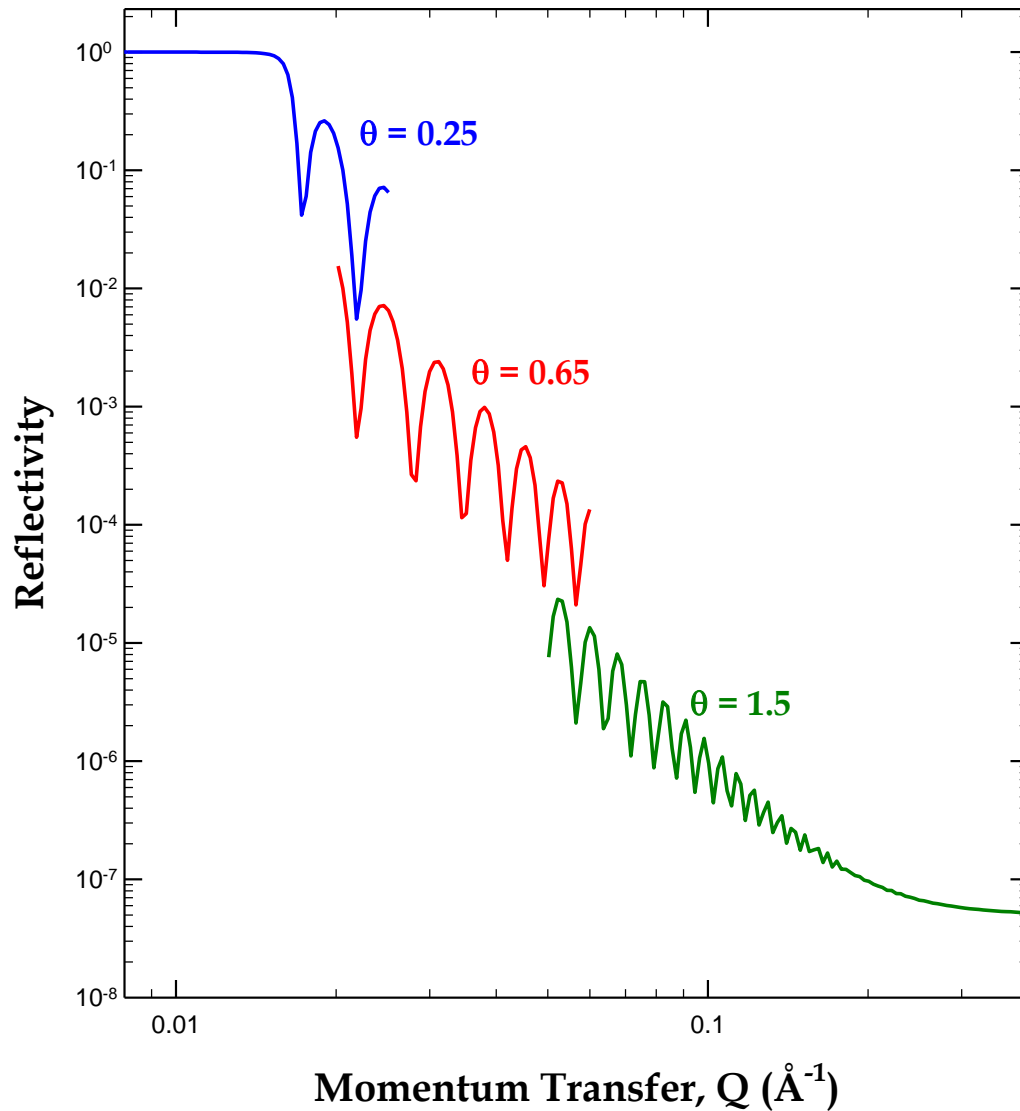
Effect of Interdiffusion on Reflectivity Profiles



Effect of Limiting Q range on Observation Window



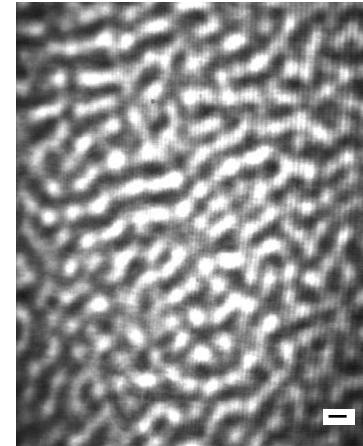
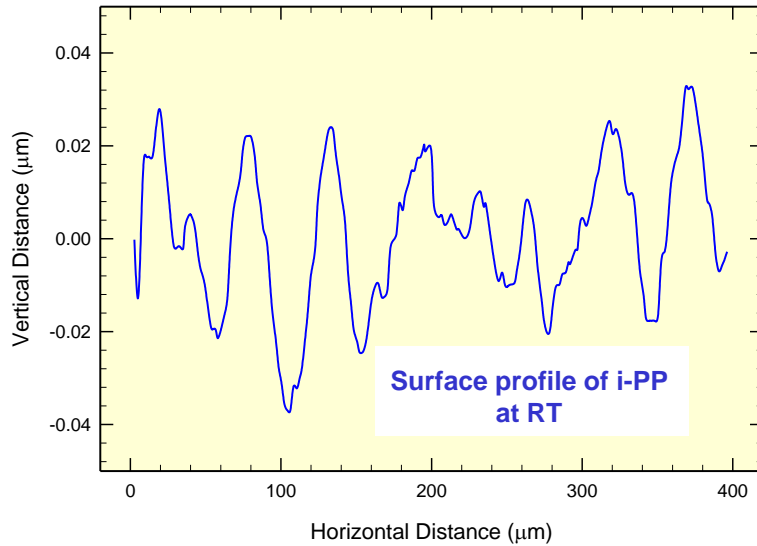
Effect of Angle on the Q Range



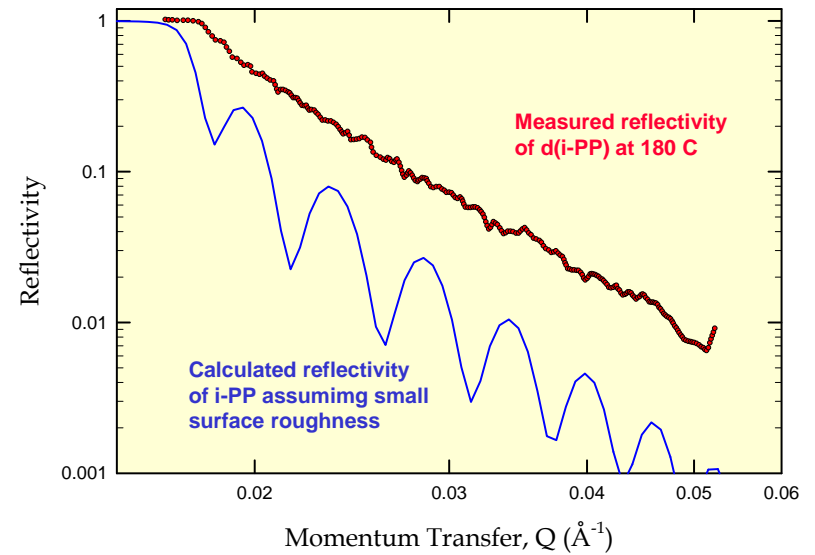
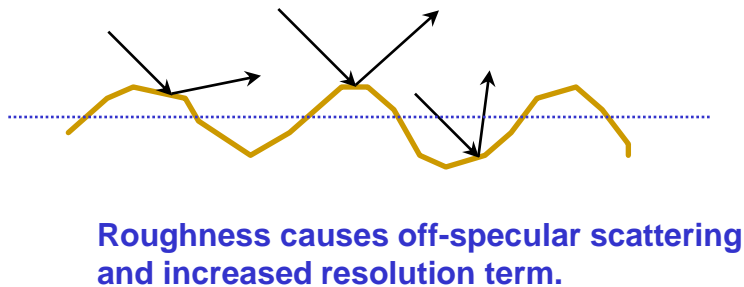
$$Q = \frac{4\pi}{\lambda} \sin \theta$$

$$0.05 < \lambda \text{ (nm)} < 0.65$$

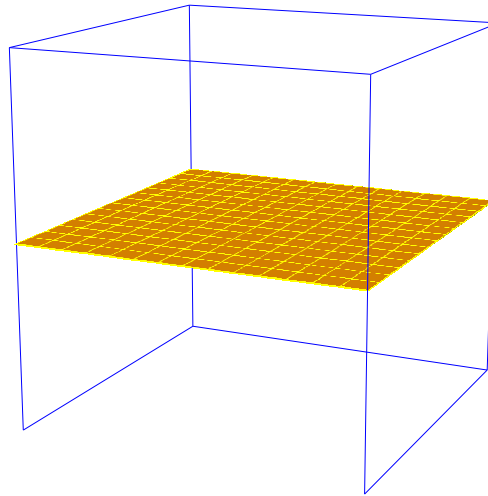
Effect of Crystallinity on reflectivity



Brewster angle micrograph of surface of i-PP (bar $20\mu\text{m}$)

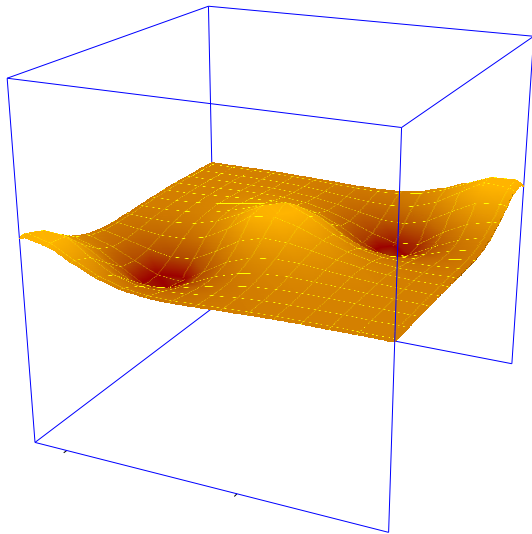


Thermally Excited Capillary Waves

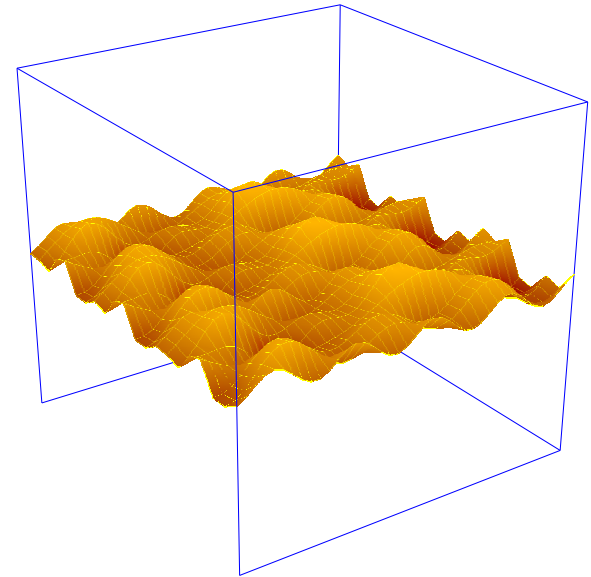
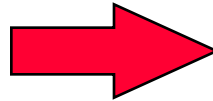


Mean field theory assumes that the interface is flat.

At equilibrium capillary waves are thermally excited.

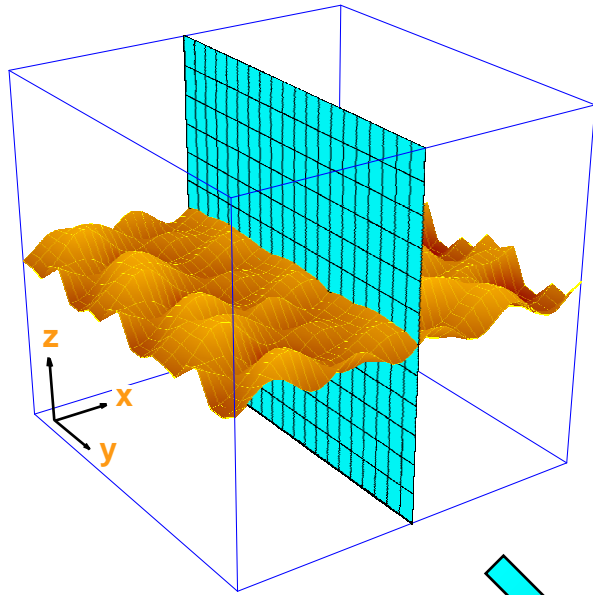


According to the equipartition theorem each mode increases the surface energy by $0.5 kT$.



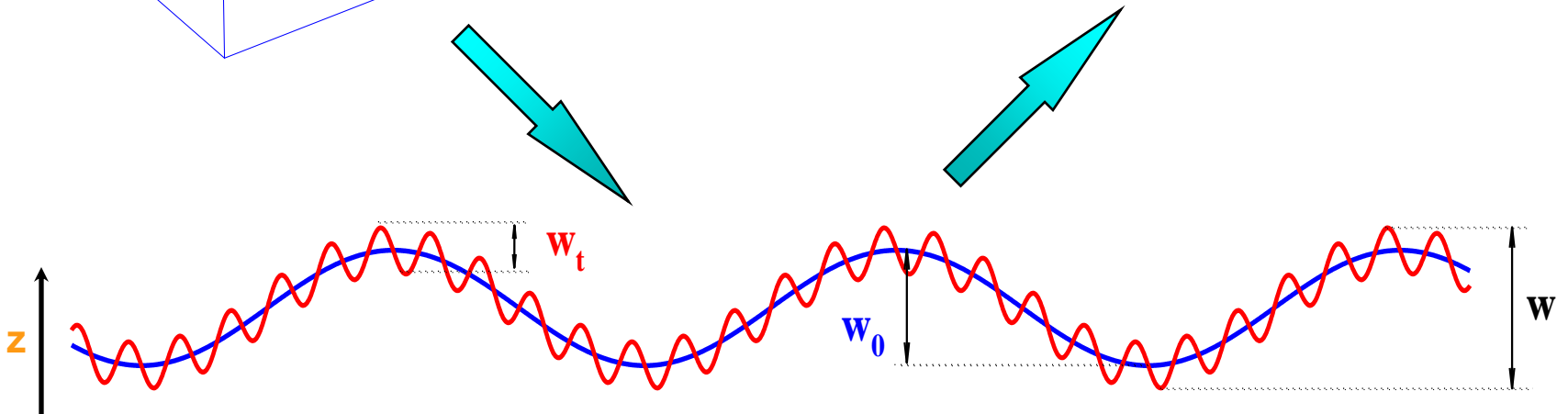
The actual surface is roughened by a superposition of all possible capillary wave modes.

NR Measured Interfacial Width



$$w = (w_t^2 + w_0^2)^{0.5}$$

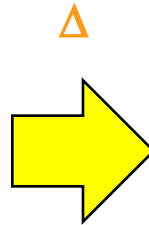
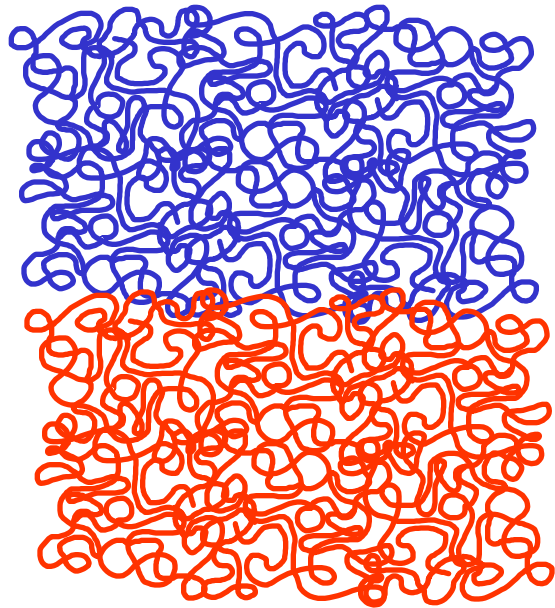
Definition of w_0 dominates derivation of w_t



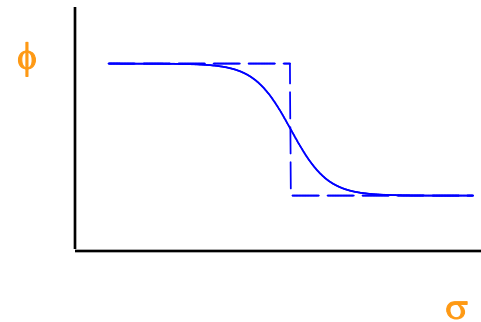
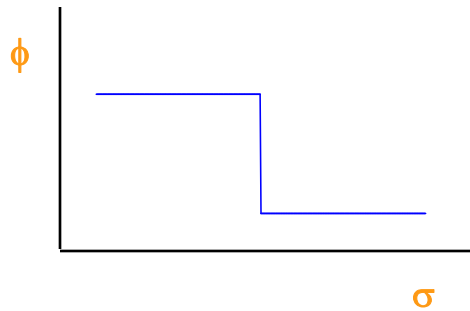
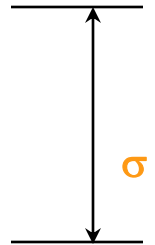
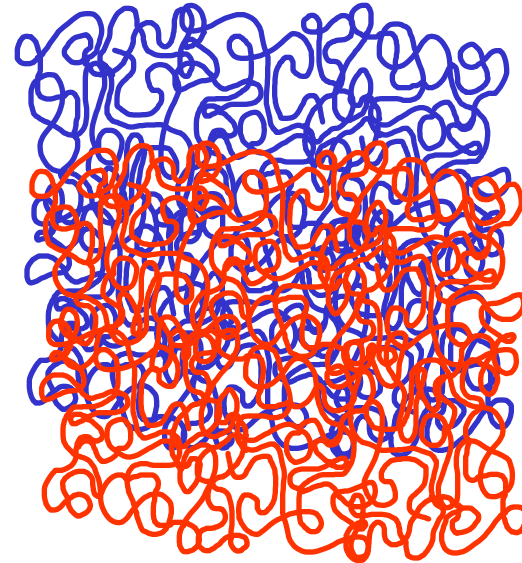
Projection onto z-y plan

Polymer Interdiffusion

As made
 $t = 0$



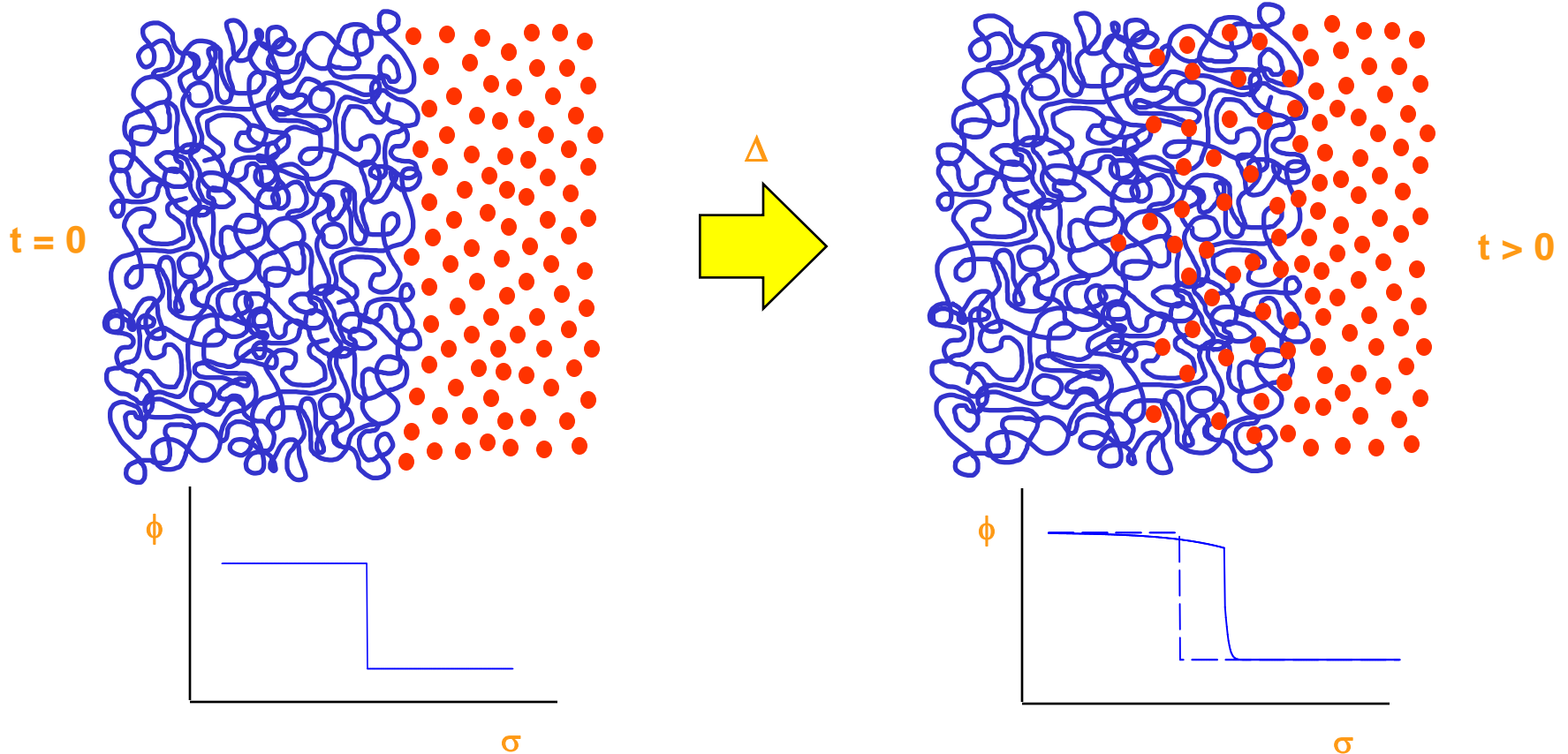
Annealed
 $t > 0$



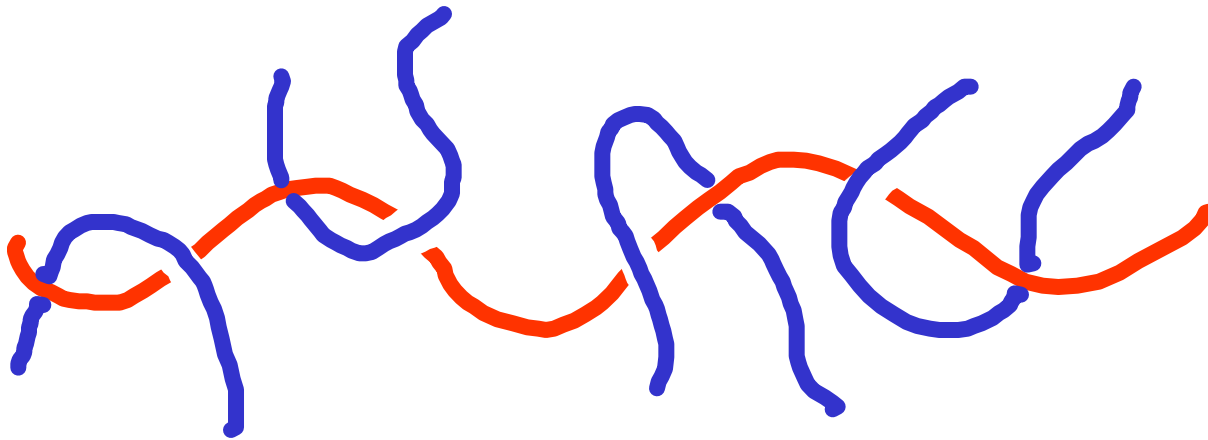
Non-Fickian Diffusion - Case II Diffusion

$$\sigma \propto t^n \quad n \neq \frac{1}{2}$$

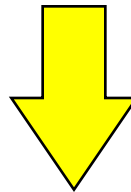
Non-Fickian Diffusion



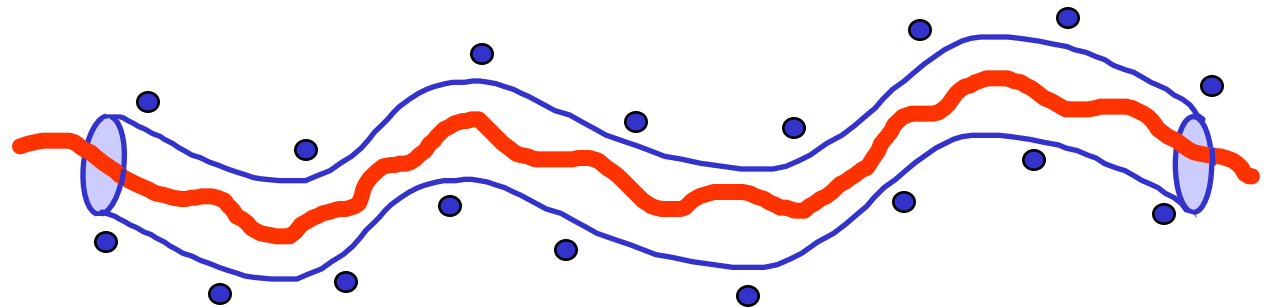
The Tube Model



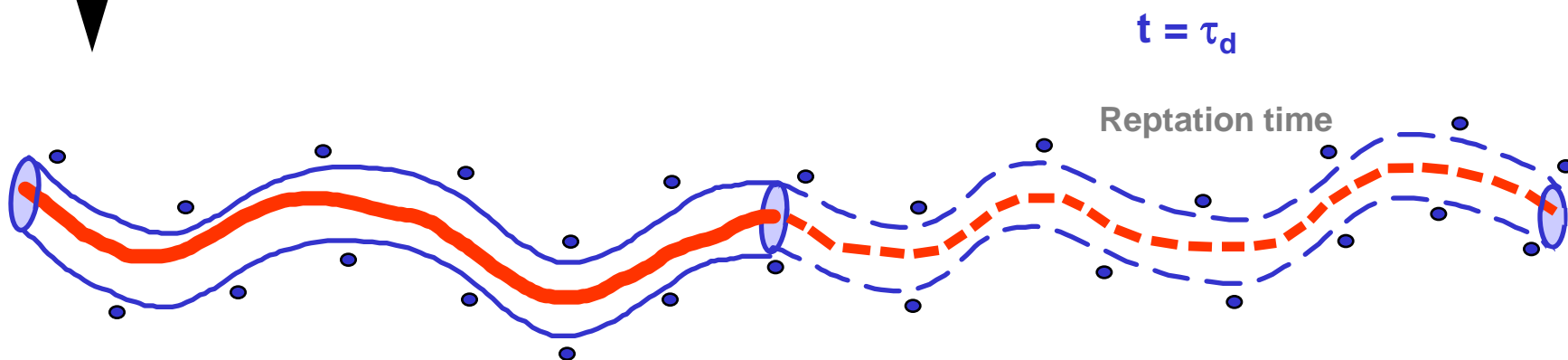
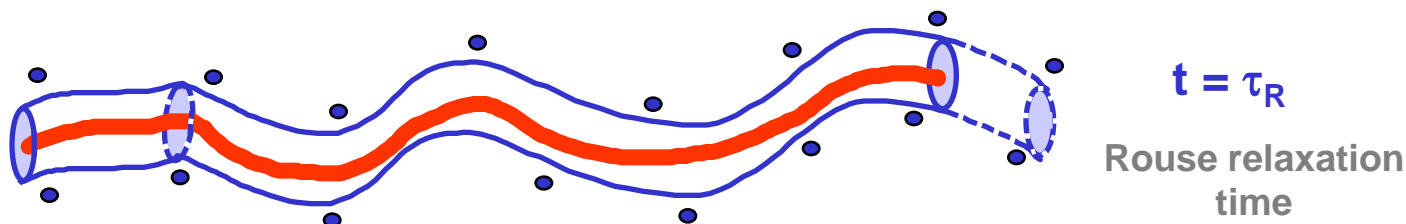
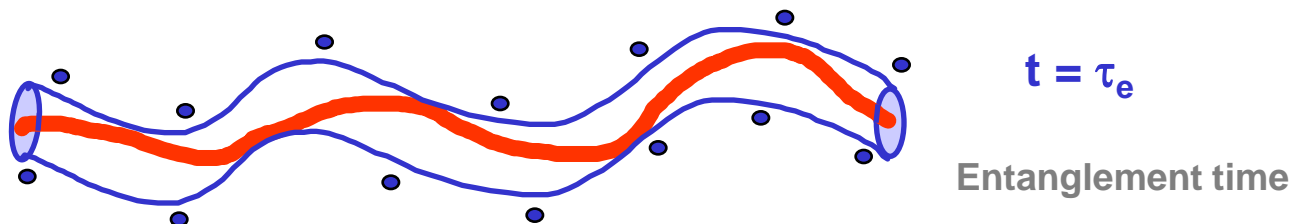
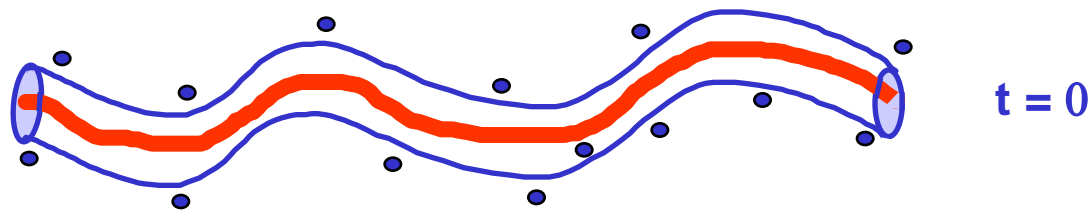
Polymer chains in the melt



Each chain can be considered to be constrained within a tube



Polymer Motion



Polymer Diffusion

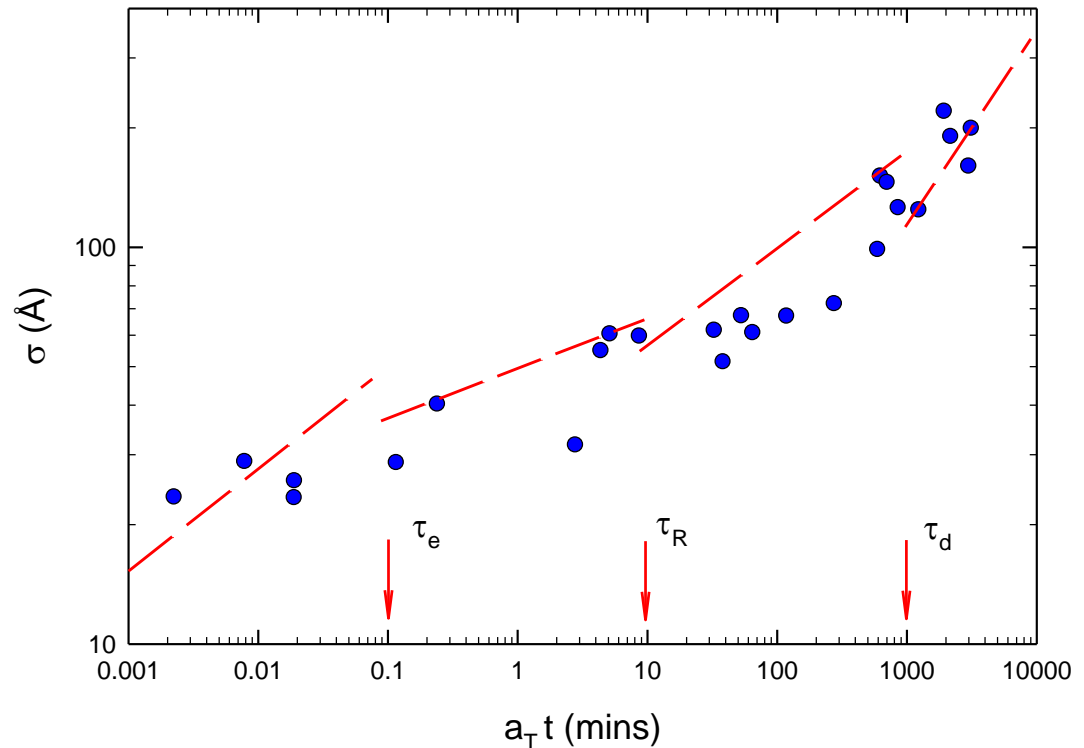
NR Results

$$t < \tau_e \quad \longrightarrow \quad \sigma \propto t^{1/4}$$

$$\tau_e < t < \tau_R \quad \longrightarrow \quad \sigma \propto t^{1/8}$$

$$\tau_R < t < \tau_d \quad \longrightarrow \quad \sigma \propto t^{1/4}$$

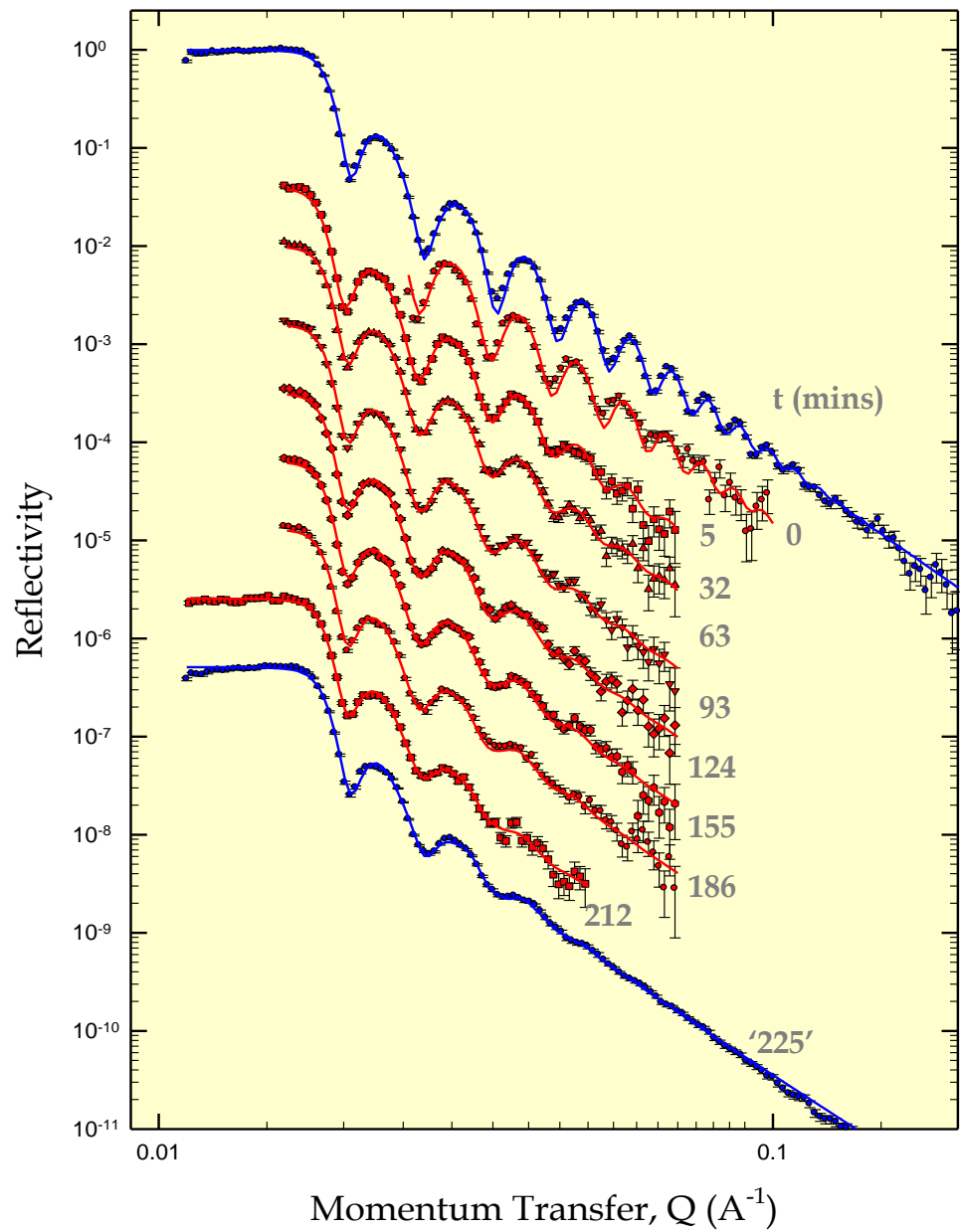
$$t > \tau_d \quad \longrightarrow \quad \sigma \propto t^{1/2}$$



A Karim et al, Phys Rev B 42 (1990) 6846

Real Time Reflectivity Measurements

Si / PS (50k) / dPS (40k) @ 115 C



Calculating a Diffusion Coefficient

$$w = \sqrt{4Dt}$$

For dPS-PS system:

$$D = (1.7 \pm 0.2) \times 10^{-17} \text{ cm}^2\text{s}^{-1}$$

$$D = \frac{k_B T d_T^2}{3N^2 \zeta b^2}$$

M Doi and SF Edwards
The Theory of Polymer Dynamics (1986)

$$D = 2.81 \times 10^{-17} \text{ cm}^2\text{s}^{-1}$$

When ζ (115C) = 0.199 dyne.s.cm⁻¹
and $d_T = 5.7 \text{ nm}$

Reptation time:

$$\tau_r = \frac{Nb^2}{3\pi^2 D}$$

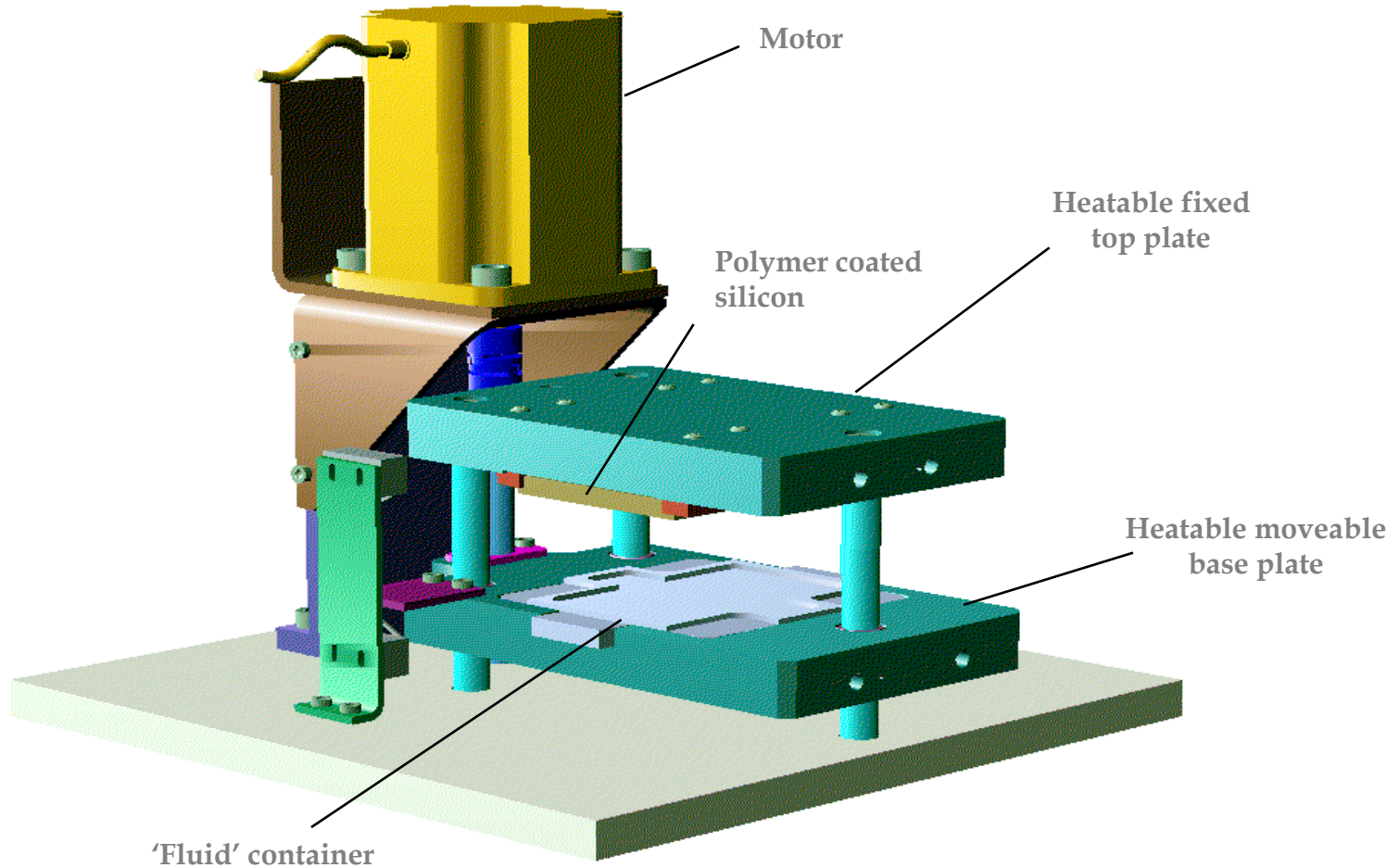
$$\begin{aligned} \tau_r &= 3223 \pm 363 \text{ s (dPS)} \\ &= 4333 \pm 489 \text{ s (hPS)} \end{aligned}$$

Rouse time:

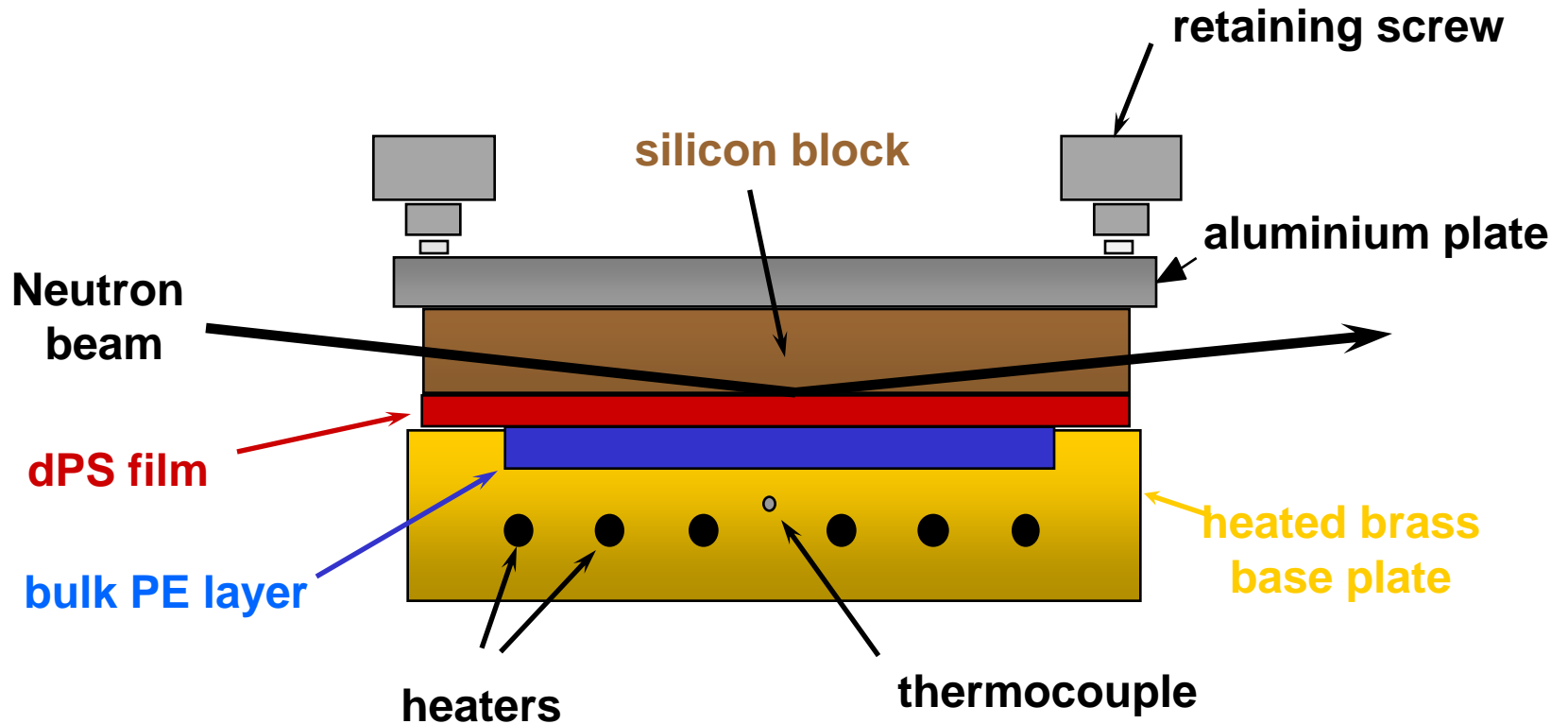
$$\tau_R = \frac{d_T^2}{9\pi^2 D}$$

$$\tau_R = 215 \pm 23 \text{ s}$$

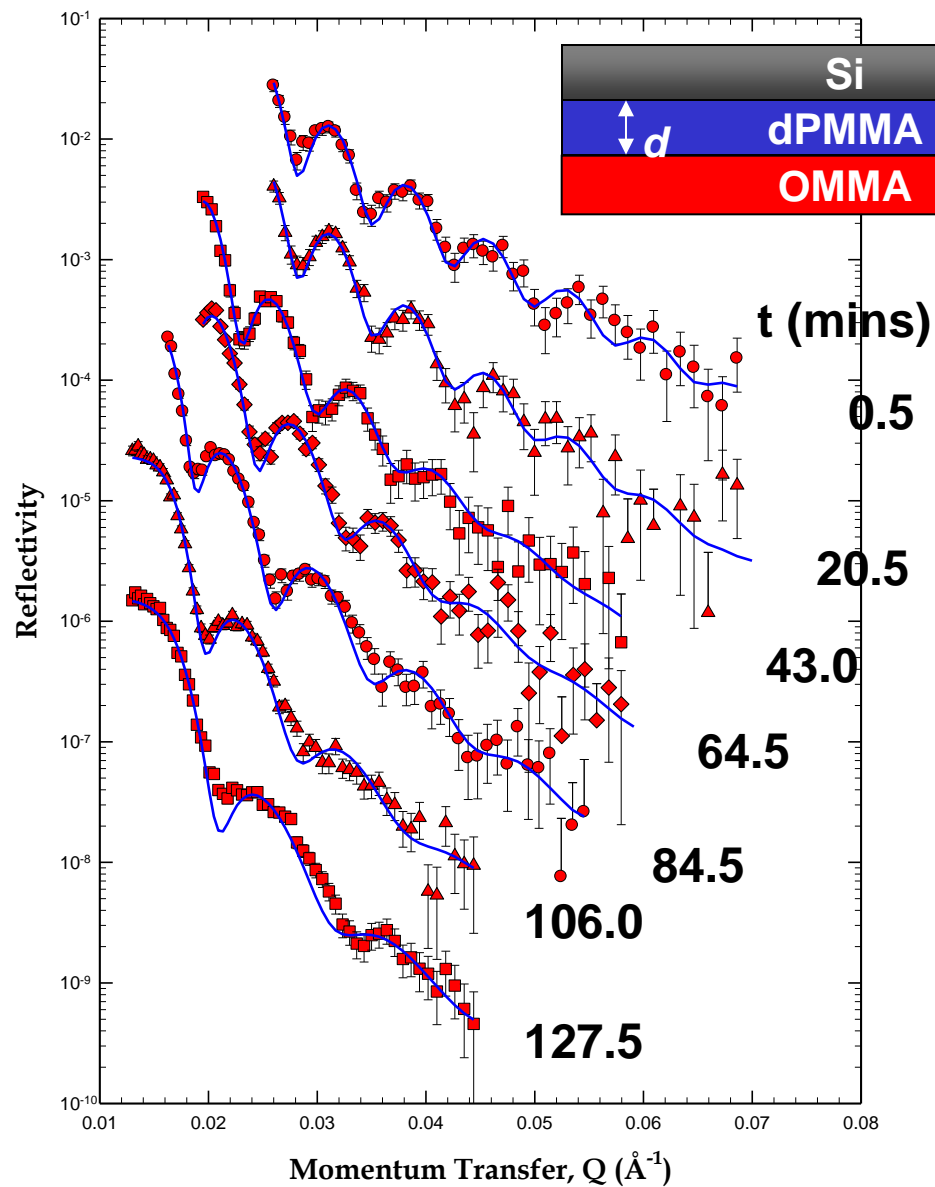
Polymer-Oligomer Interdiffusion Reflectivity Cell



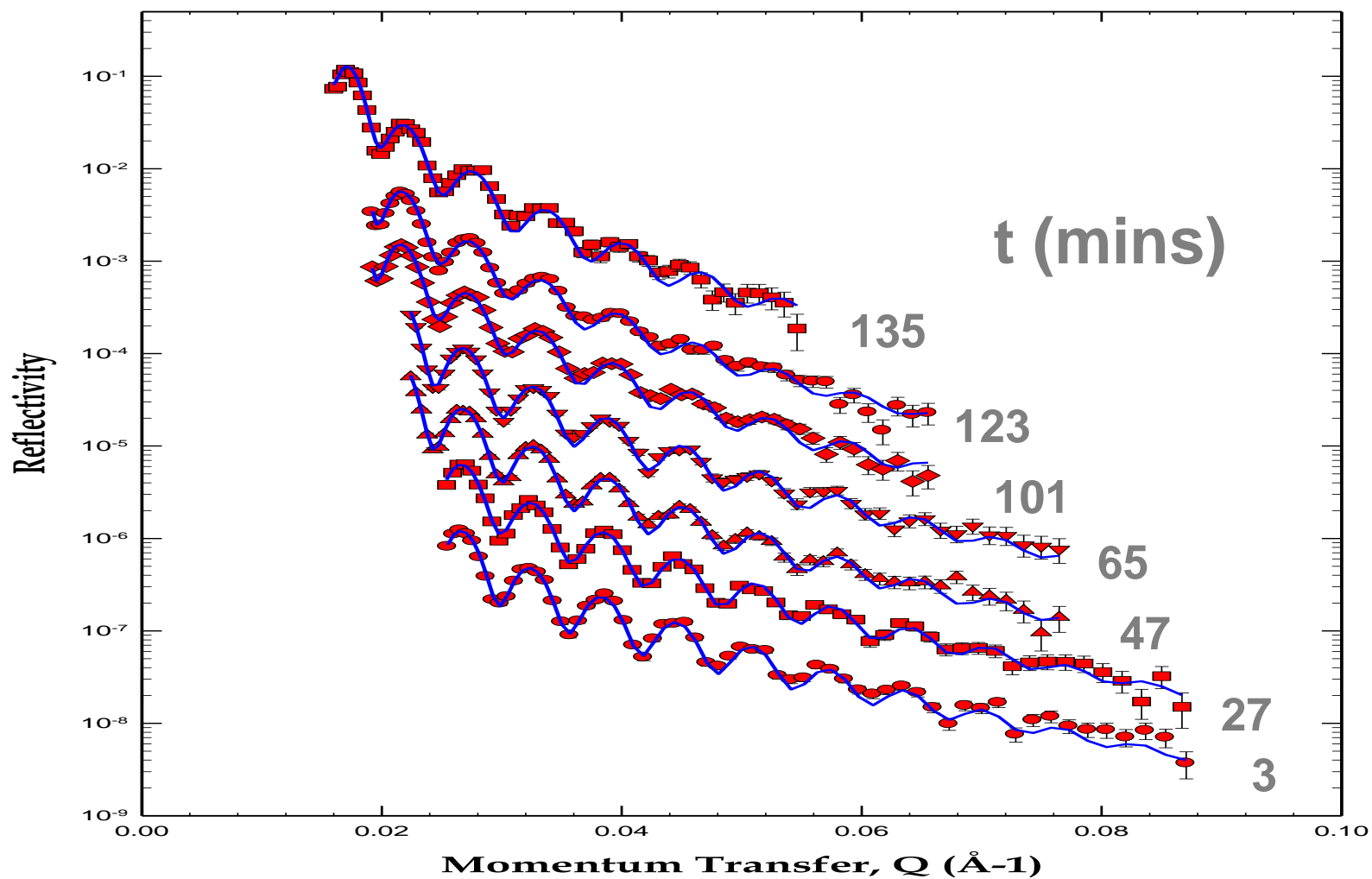
Neutron Reflectivity Melt Cell



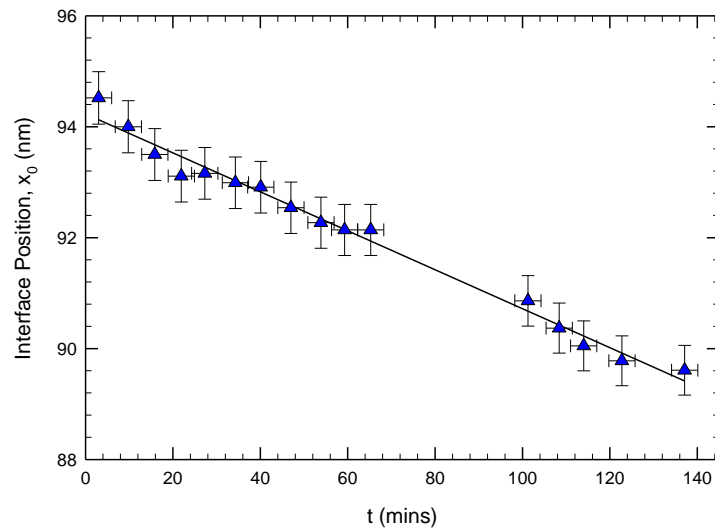
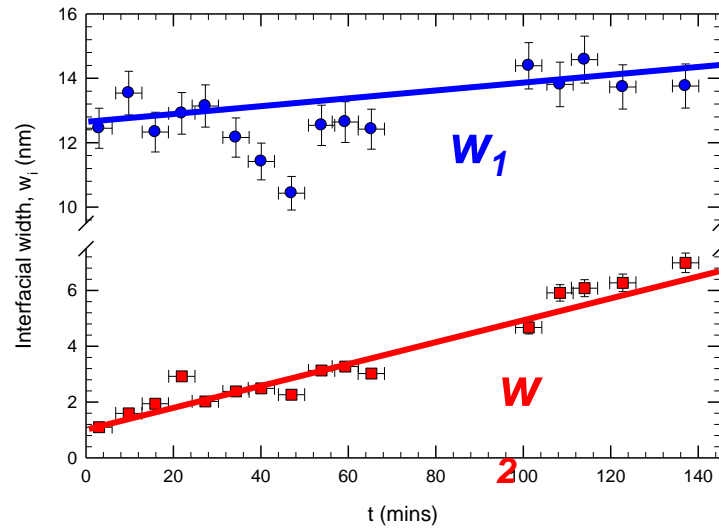
dPMMA(100k) / OMMA(510) @ 45 C



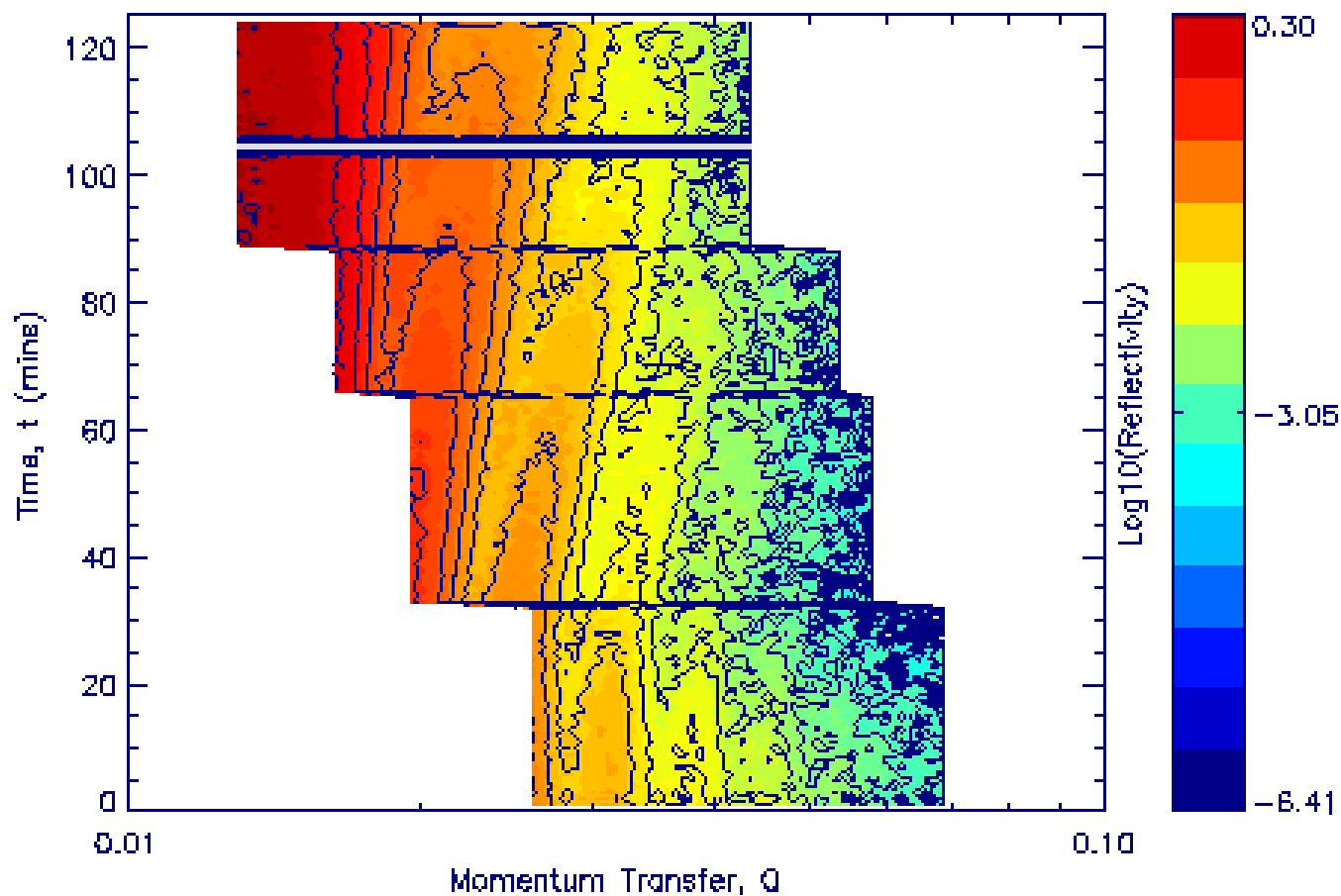
dPS (101k) / OSt (1100) Interdiffusion @ 65C



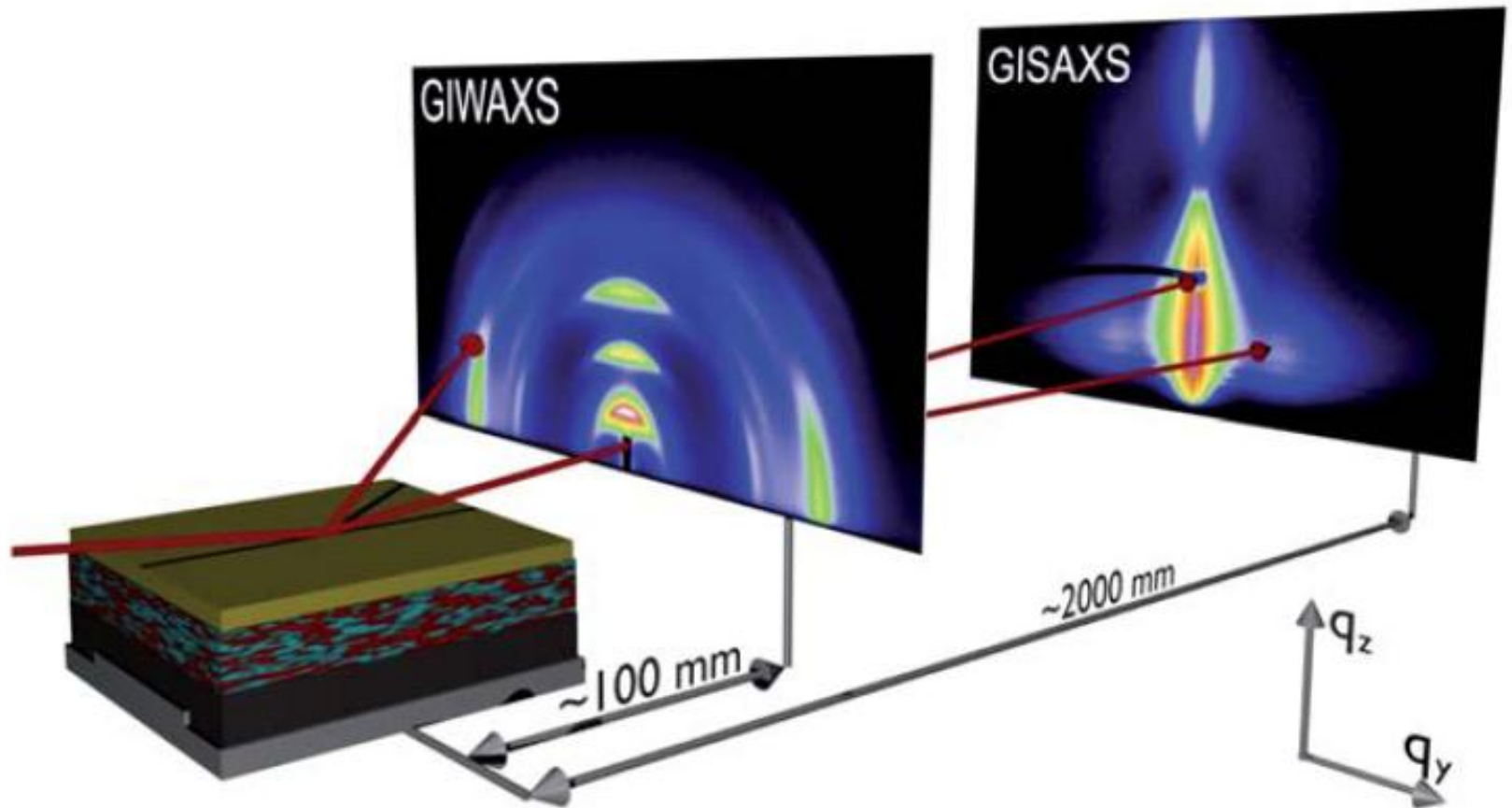
dPS (101k) / OSt (1100) Interdiffusion @ 65C



Off-specular reflection

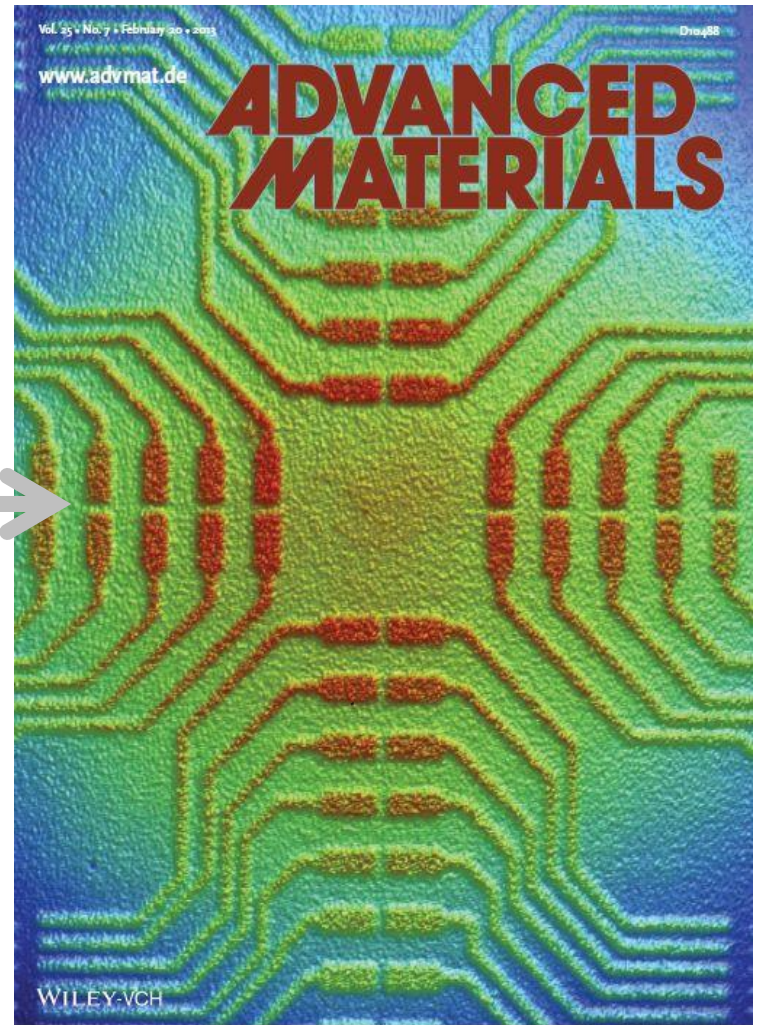
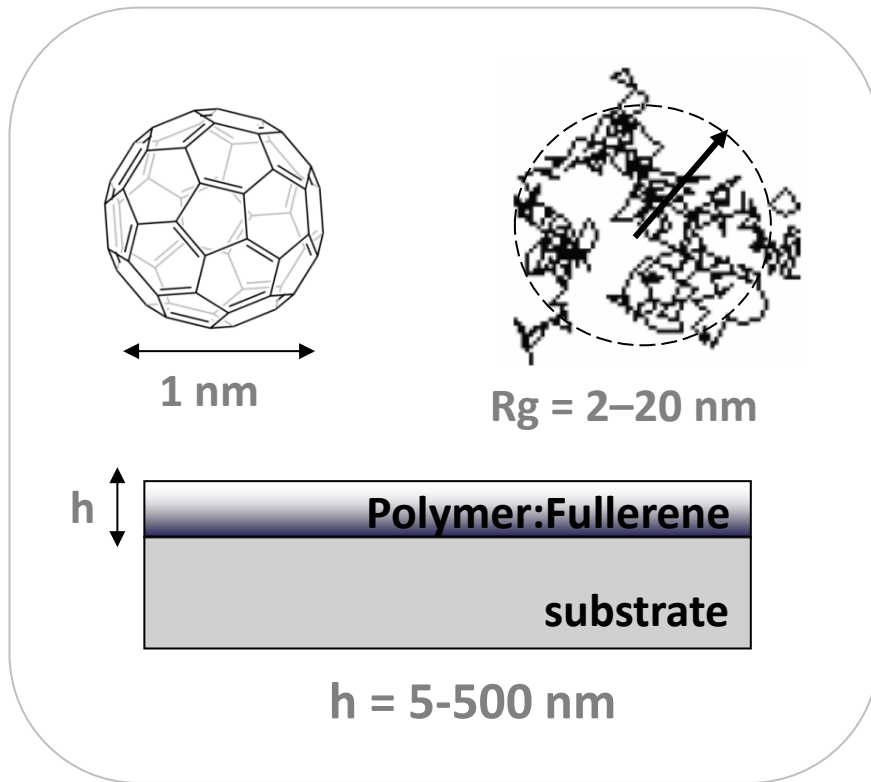


Grazing incidence: wide and small angle



P Muller (2011)

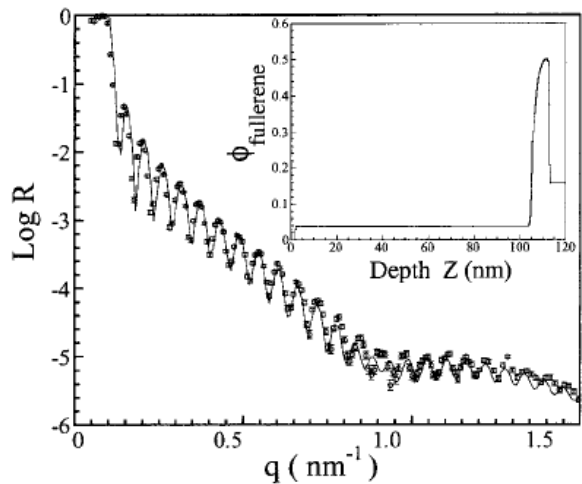
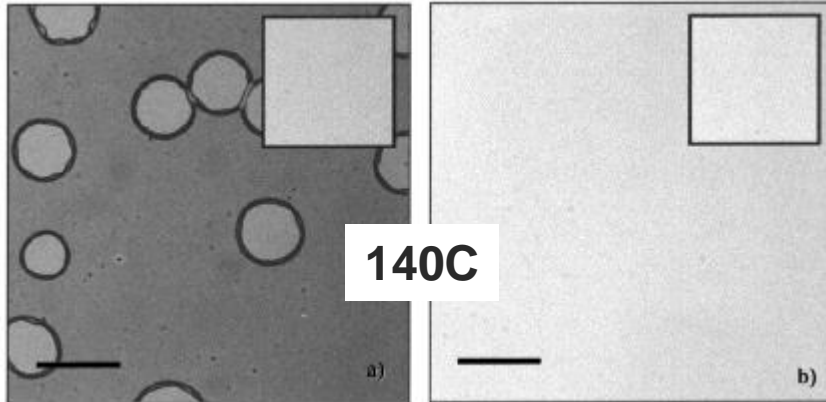
Polymer-fullerene blends



Thin films

PS (2k), 30 nm

1% C60, 30 nm

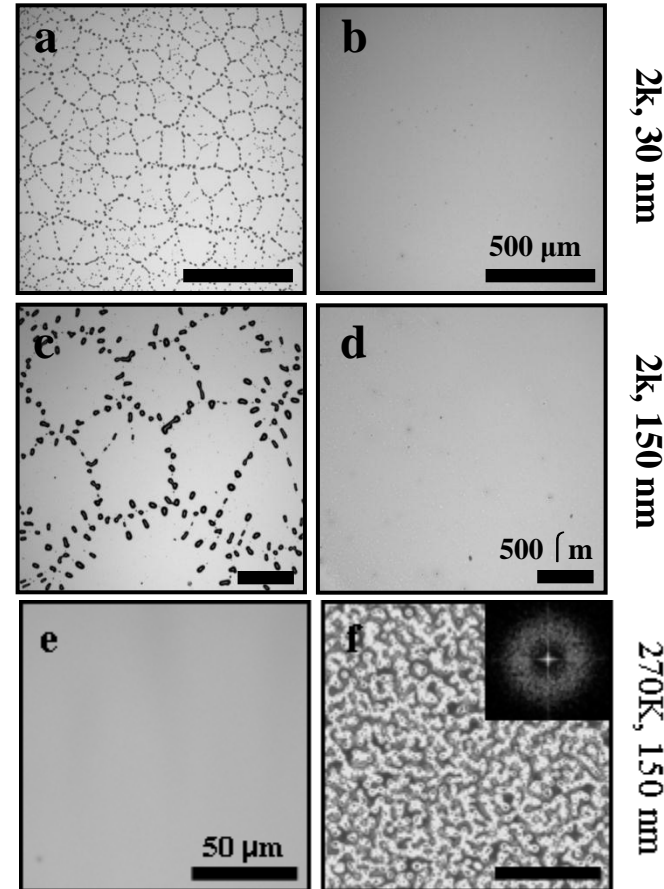


PS+5% C60 $h = 100\text{nm}$

GISANS & reflectometry

PS

PS + 5% C₆₀



140°C

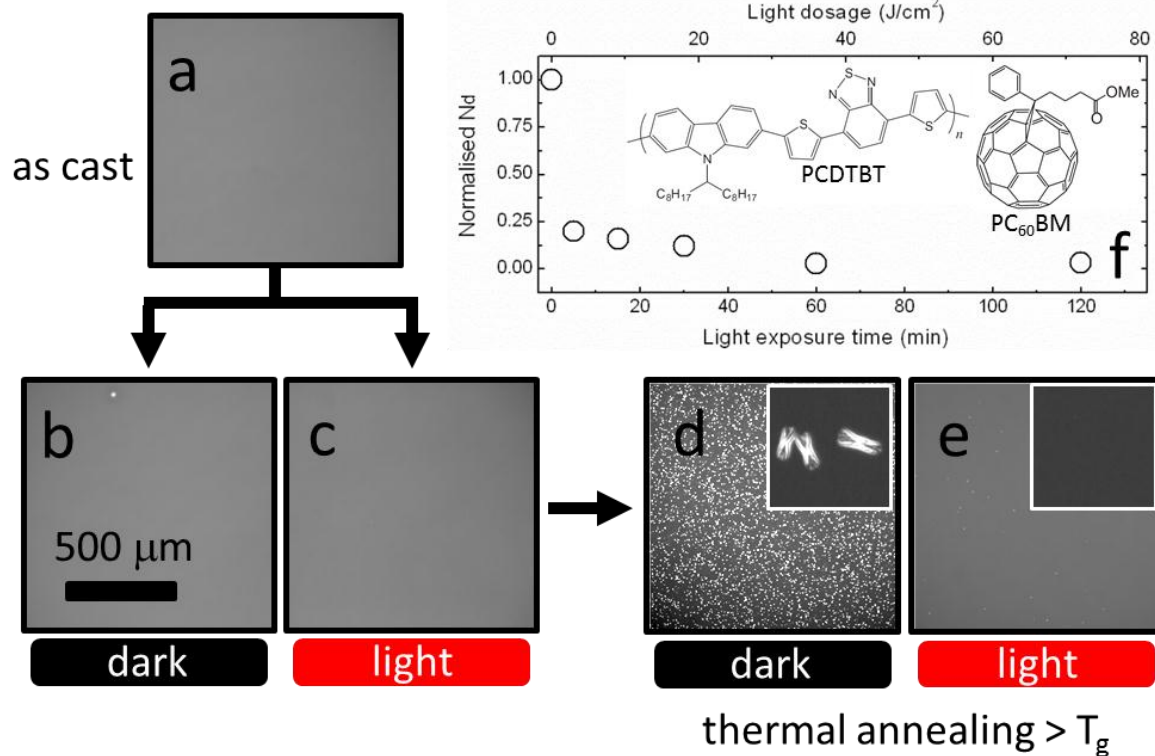
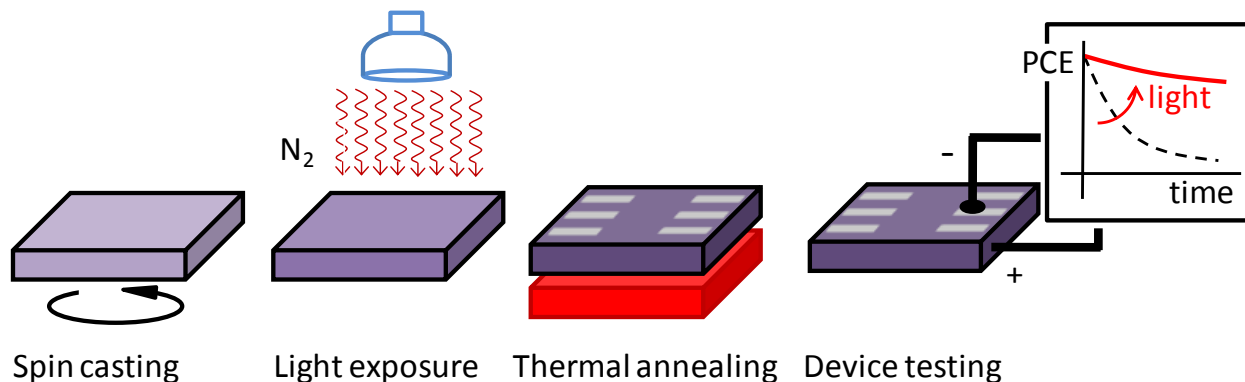
180°C

180°C

'Spinodal nucleation'

Phys. Rev. Lett. **105**, 038301 (2010)
Macromolecules **44**, 4530-4537 (2011)

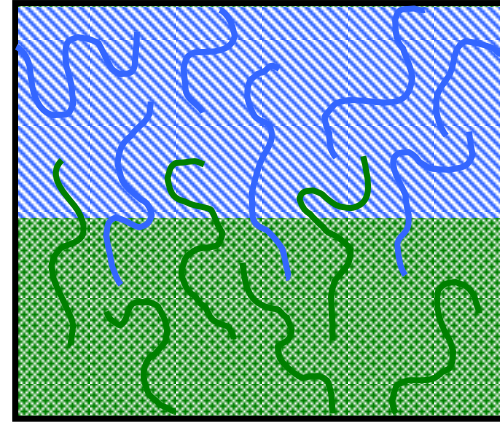
Organic Solar Cell lifetime?



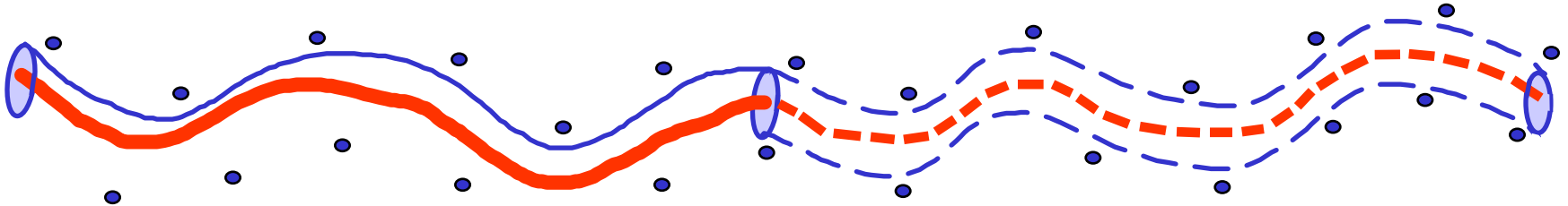
Summary

Reflectivity

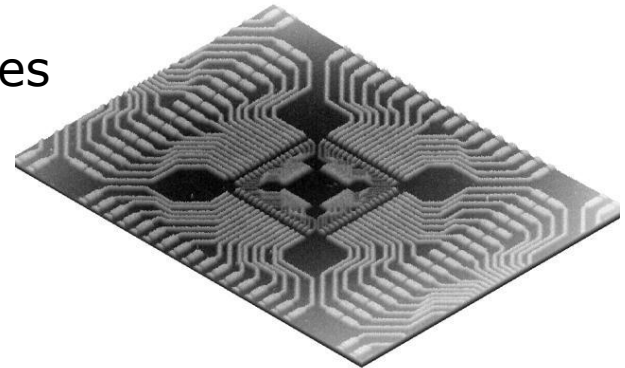
- study and design interfaces



- investigate diffusion mechanisms



- engineer 'functional' surfaces / devices



Neutrons in soft matter

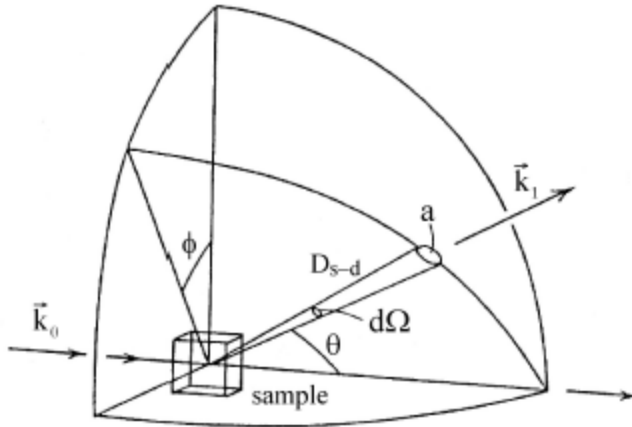
Lecture 2 (II) – Dynamics

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Scattering theory reminder

Scattering cross section



$$\frac{d^2\sigma}{d\Omega dE} = \left(\frac{d^2\sigma}{d\Omega dE} \right)_{coh} + \left(\frac{d^2\sigma}{d\Omega dE} \right)_{inc}$$

coherent incoherent

$$\left(\frac{d^2\sigma}{d\Omega dE} \right)_{coh} = \frac{1}{2\pi\hbar} \frac{k_1}{k_0} \frac{\sigma_{coh}}{4\pi} \int_{-\infty}^{+\infty} \sum_{i,j} \langle e^{-i\mathbf{q}\cdot\mathbf{R}_i(0)} e^{i\mathbf{q}\cdot\mathbf{R}_j(t)} \rangle e^{-i\omega t} dt$$

$$\left(\frac{d^2\sigma}{d\Omega dE} \right)_{inc} = \frac{1}{2\pi\hbar} \frac{k_1}{k_0} \frac{\sigma_{inc}}{4\pi} \int_{-\infty}^{+\infty} \sum_i \langle e^{-i\mathbf{q}\cdot\mathbf{R}_i(0)} e^{i\mathbf{q}\cdot\mathbf{R}_i(t)} \rangle e^{-i\omega t} dt$$

Dynamic structure factor

FT (t, ω) \Uparrow $S(\mathbf{q}, \omega) = \frac{1}{2\pi\hbar} \int_{-\infty}^{+\infty} I(\mathbf{q}, t) e^{-i\omega t} dt.$

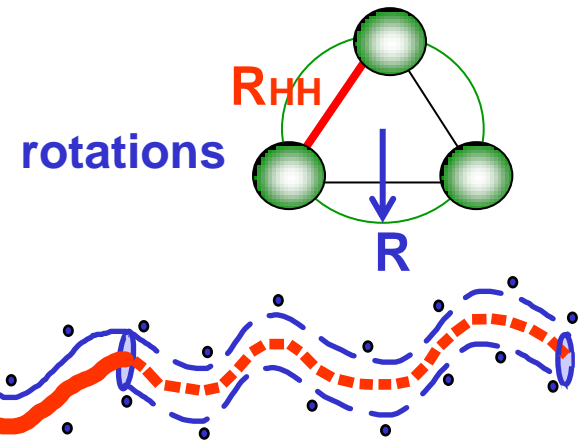
Intermediate scattering function

FT (r, q) \Uparrow $I_s(\mathbf{q}, t) = \frac{1}{N} \sum_i \langle e^{-i\mathbf{q}\cdot\mathbf{R}_i(0)} e^{i\mathbf{q}\cdot\mathbf{R}_i(t)} \rangle e^{-i\omega t}.$

Pair correlation function

$$G(\mathbf{r}, t) = \frac{1}{(2\pi)^3} \int I(\mathbf{q}, t) e^{-i\mathbf{q}\cdot\mathbf{r}} d\mathbf{q}.$$

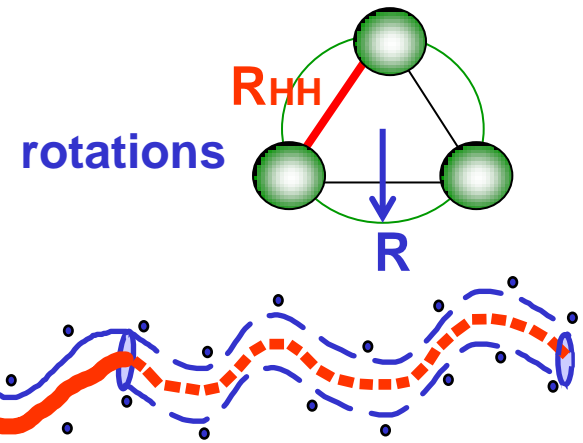
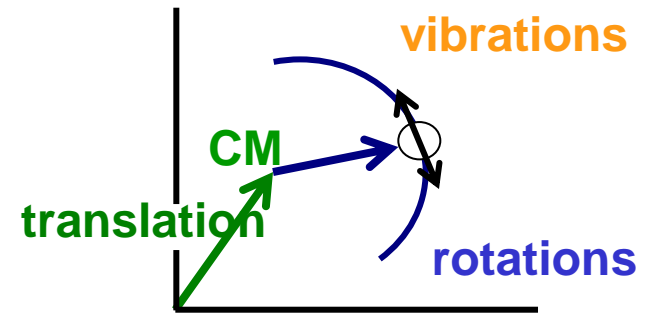
\rightleftharpoons vibrations



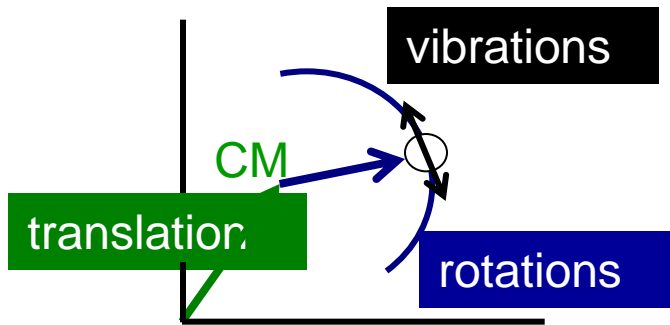
single-particle dynamics

motion decomposition

$$I_{self}(Q, t) = \frac{1}{N} \sum_i \left\langle e^{iQ \cdot [V(t) - V(0)]} \right\rangle \left\langle e^{iQ \cdot [T(t) - T(0)]} \right\rangle \left\langle e^{iQ \cdot [R(t) - R(0)]} \right\rangle$$



single-particle tools



motion decomposition

$$I_{self}(Q, t) = \frac{1}{N} \sum_i \left\langle e^{iQ \cdot [V(t) - V(0)]} \right\rangle \left\langle e^{iQ \cdot [T(t) - T(0)]} \right\rangle \left\langle e^{iQ \cdot [R(t) - R(0)]} \right\rangle$$

CM translation

frozen for polymers $T \ll T_g$.

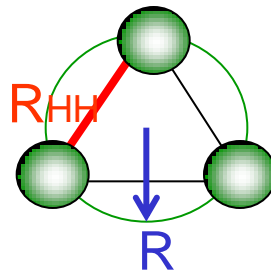
Proton delocalisation

DW factor: $e^{-\frac{1}{3}Q^2 \langle u^2 \rangle}$

relevant proton reorientations: methyl and phenyl rotations about group's axis.

Methyl protons 3-fold jumps

$$R \approx 1.032 \text{ \AA}$$



$$S_{rot}(Q, \omega) = A_0(Q)\delta(\omega) + A_1(Q) \frac{1}{\pi} \frac{3/2\tau}{(3/2\tau)^2 + \omega^2}$$

with

$$A_0(Q) = \frac{1}{3} [1 + 2j_0(Qr\sqrt{3})]$$

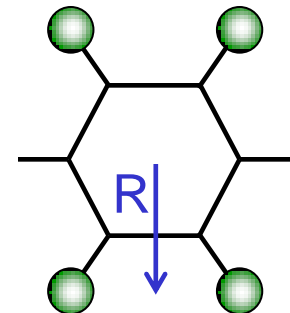
$$A_1(Q) = 1 - A_0(Q)$$

Phenyl proton 2-fold jumps

$$R \approx 2.28 \text{ \AA}$$

$$S_{rot}(Q, \omega) = A_0(Q)\delta(\omega) + A_1(Q) \frac{1}{\pi} \frac{2/\tau}{(2/\tau)^2 + \omega^2}$$

with $A_0(Q) = \frac{1}{2} [1 + j_0(2Qr)]$

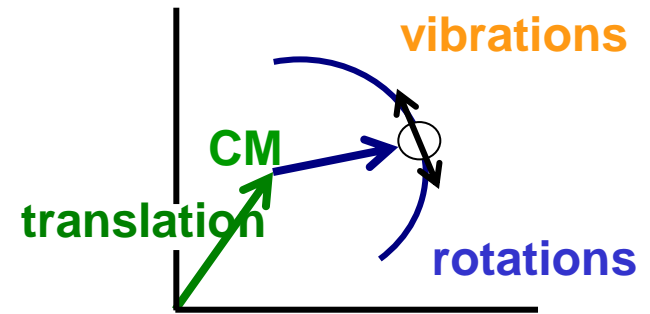


single-particle dynamics

motion decomposition in the glass

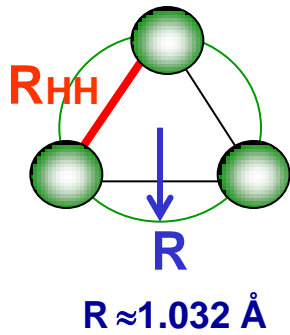
CM translation: frozen for polymers $T \ll T_g$.

Proton delocalisation: DW factor: $e^{-\frac{1}{3}Q^2 \langle u^2 \rangle}$

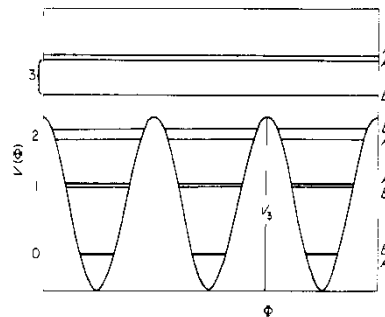


example:

Side group rotations:



3-fold CH₃ potential



Methyl protons 3-fold jumps

$$S_{\text{rot}}(Q, \omega) = A_0(Q)\delta(\omega) + A_1(Q)\frac{1}{\pi} \frac{3/2\tau}{(3/2\tau)^2 + \omega^2}$$

with

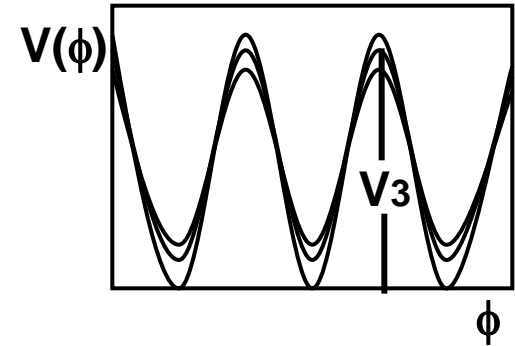
$$\begin{aligned} A_0(Q) &= \frac{1}{3} [1 + 2j_0(Qr\sqrt{3})] \\ A_1(Q) &= 1 - A_0(Q) \end{aligned}$$

distribution $\tau_{\text{correlation}}$

glassy polymers: no single relaxation time

variety local
environments

intra- molecular
inter-



(Gaussian) distribution of potential barriers:

$$g(E_i) = \frac{1}{\sigma_E \sqrt{2\pi}} e^{-\frac{(E_i - E_0)^2}{2\sigma_E^2}} \quad \text{if} \quad \Gamma = \Gamma_0 e^{-\frac{E_A}{RT}}$$

(log-Gaussian) distribution of reorientation times:

$$g(\ln \Gamma_i) = \frac{1}{\sigma \sqrt{2\pi}} e^{-\frac{\ln^2(\Gamma_i/\Gamma_0)}{2\sigma^2}} \quad \begin{array}{l} E_0: \text{average barrier height} \\ \sigma: \text{distribution width} \end{array}$$

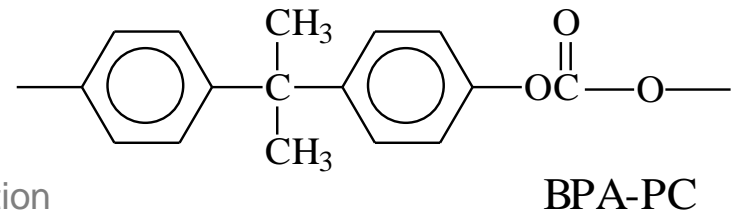
Dynamic structure factor: $S_{\text{rot}}(Q, \omega) = A_0(Q)\delta(\omega) + A_1(Q) \sum_{i=1}^N g_i L_i(\omega)$

Case study: Polycarbonates

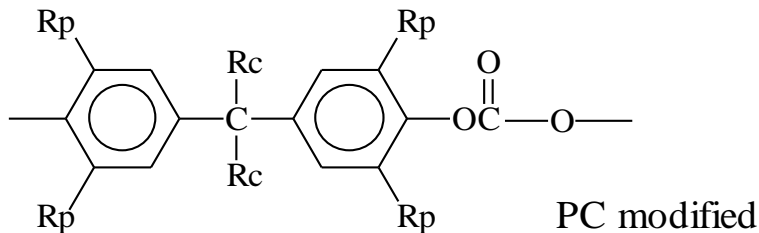
Bisphenol-A polycarbonate

thermoplastic polymer with remarkable

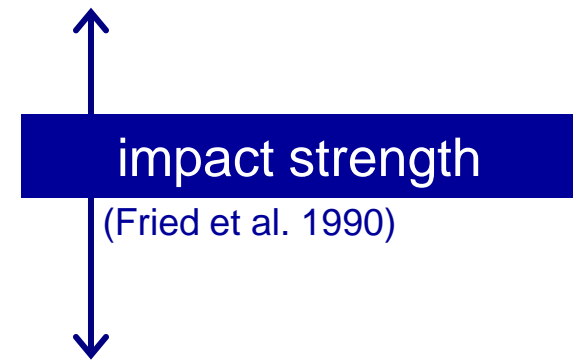
- optical clarity
- **mechanical properties**
 - high $T_{\text{glass transition}}$
 - large impact strength
 - ductility.
- commercial applications



depend strongly on architecture

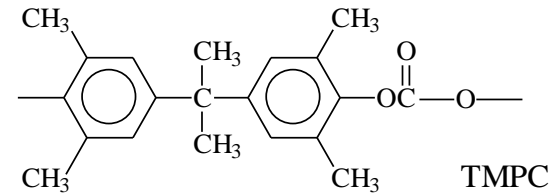
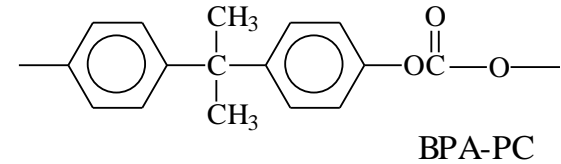
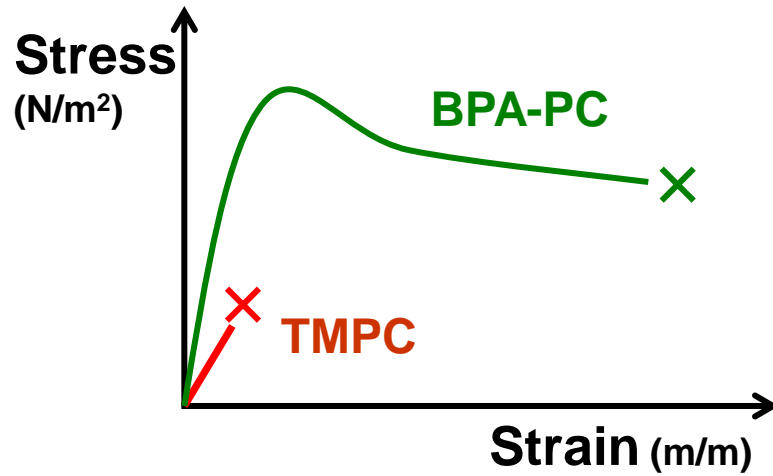


BPA-PC: ~2400 J/m



TMPC: ~70 J/m

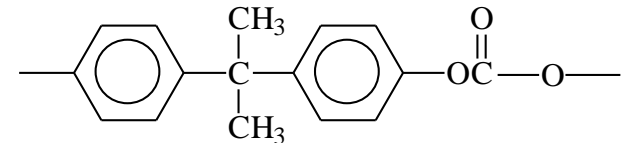
Toughness



Polycarbonates

Glassy BPAPC

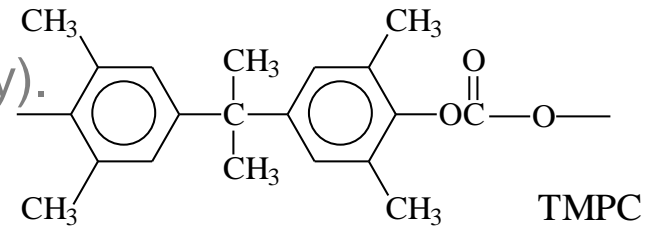
tough → co-operative phenyl motion,
involve ≥ 1 monomer
(account for dielectric/mechanical activity).



BPA-PC

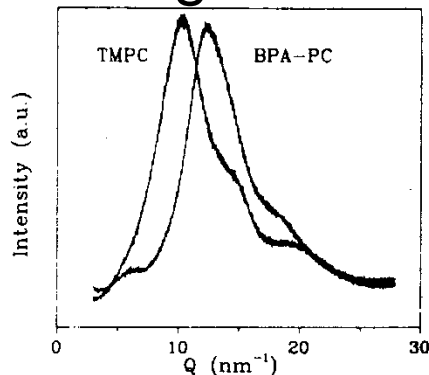
Glassy TMPC

most brittle PC → substituted CH_3 hinder backbone mobility;
poor chain packing (large free volume).



TMPC

Packing



$$\rho(\text{PC}) = 1.198 \text{ g/cm}^3$$
$$\rho(\text{TMPC}) = 1.084 \text{ g/cm}^3$$

QENS:

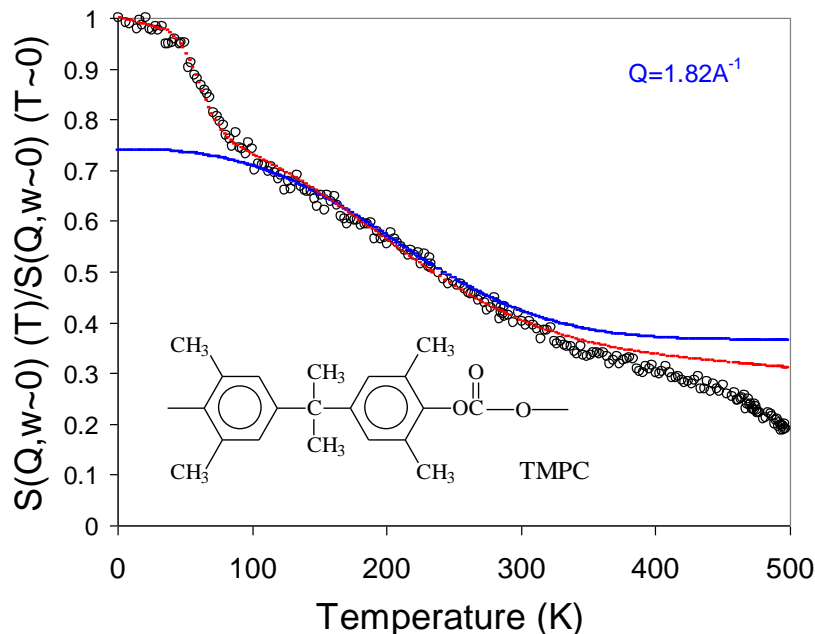
characterise dynamics of local reorientation.

quantitative window scans

Elastic scans

$$S(Q, \omega \sim 0) = \int_{-\infty}^{+\infty} S(Q, \omega') R(\omega - \omega') d\omega' \bigg|_{\omega=0}$$

for a Lorentzian resolution $S(Q, \omega \sim 0) \approx A_0(Q) + \frac{2}{\pi} [1 - A_0(Q)] \arctan\left(\frac{\Gamma_{\text{res}}}{\Gamma}\right)$



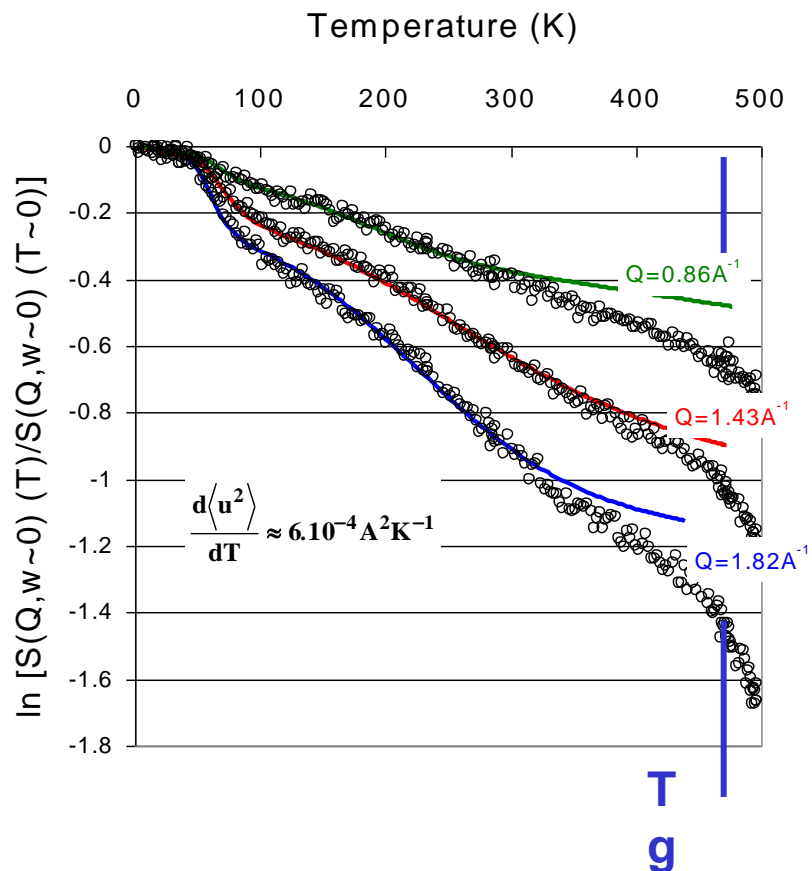
PARAMETERS

- $\langle u^2 \rangle(T) \leftarrow$ initial slope
- distribution: E_A and σ
- Γ_0

ASSUMED

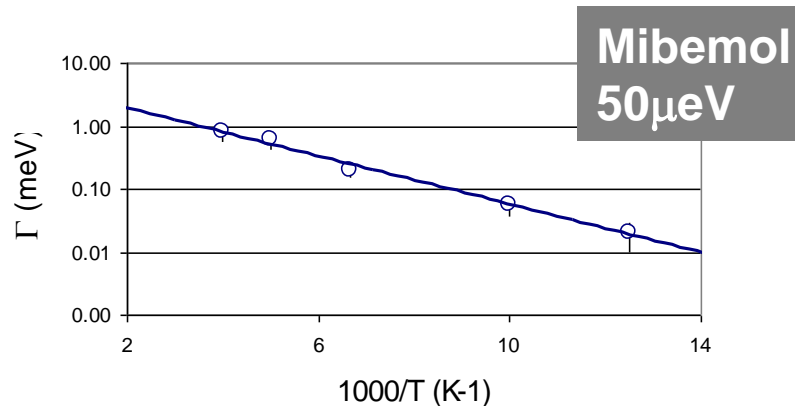
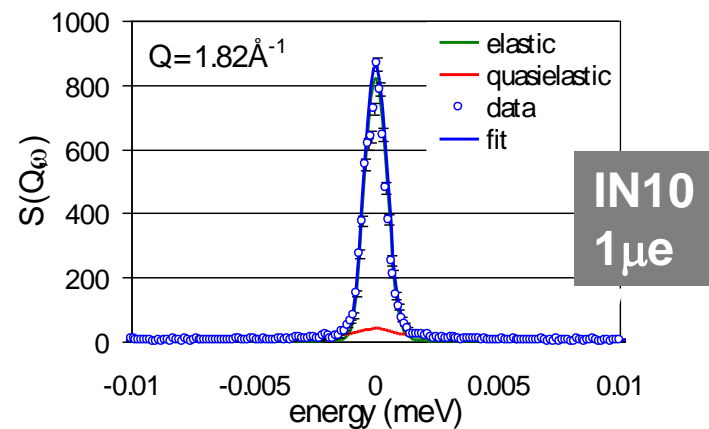
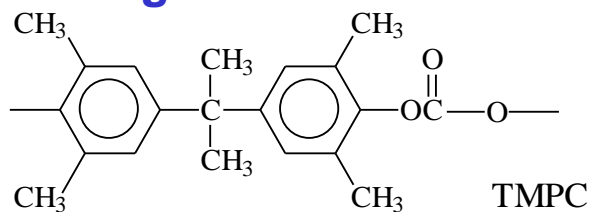
- geometry \leftarrow EISF
- activation ansatz: $\Gamma = \Gamma_0 e^{-\frac{E_A}{RT}}$

TMPC



Ea1~6
kJ/mol $\sigma 1 \sim 1$

Ea2=15
kJ/mol $\sigma 1 \sim 5$



low temperature relaxation

TMPC first relaxation step:

- very low $T \rightarrow$ low E_0
- rather sharp \rightarrow narrow

\rightarrow candidate: rotational tunneling

Mathiew equation: inelastic lines

$$S_{\text{rot}}(Q, \omega) = \frac{5 + 4j\omega(Qr)}{9} \delta(\omega) + \frac{2(1 - j\omega(Qr))}{9} [\delta(\omega - \omega_t) + \delta(\omega + \omega_t)]$$

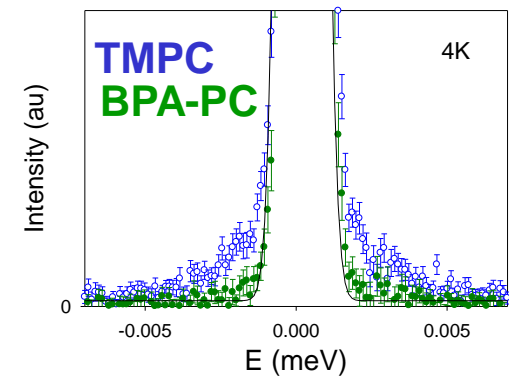
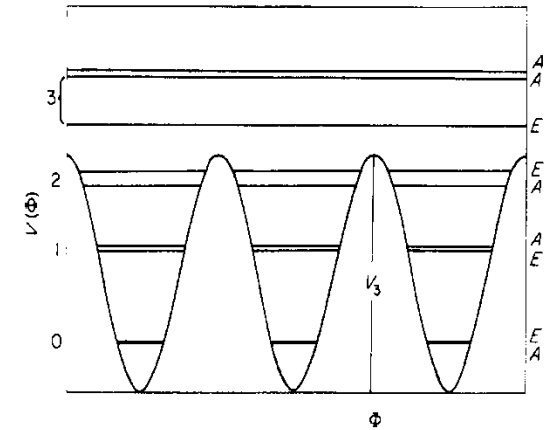
with $\hbar\omega_t \propto E_A^{3/4} e^{-\sqrt{E_A}}$

Distribution of $E_A \rightarrow$

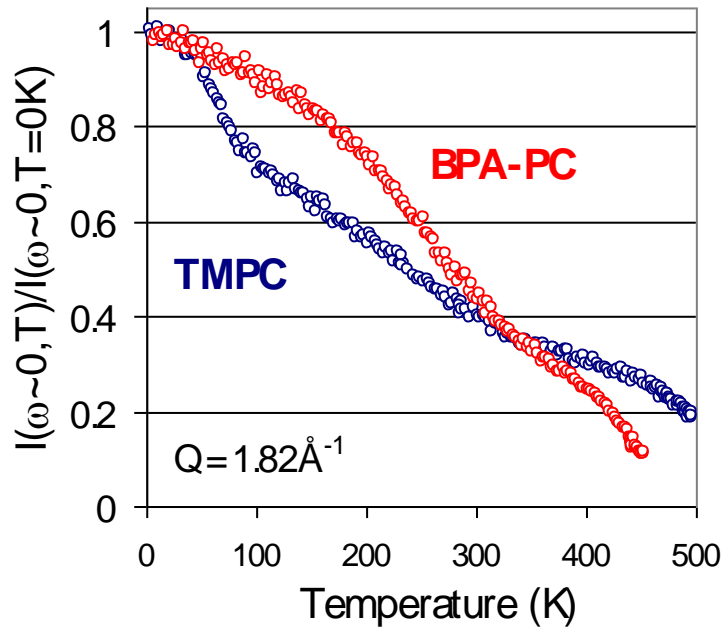
highly asymmetric distribution of ω_t

(Colmenero et al, PRL 1998)

3-fold CH_3 potential



BPA-PC

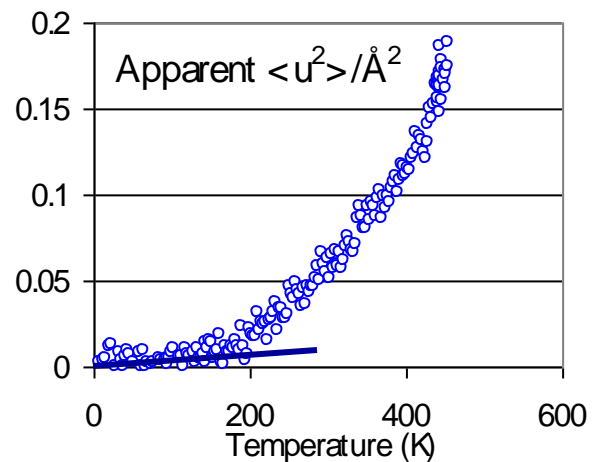
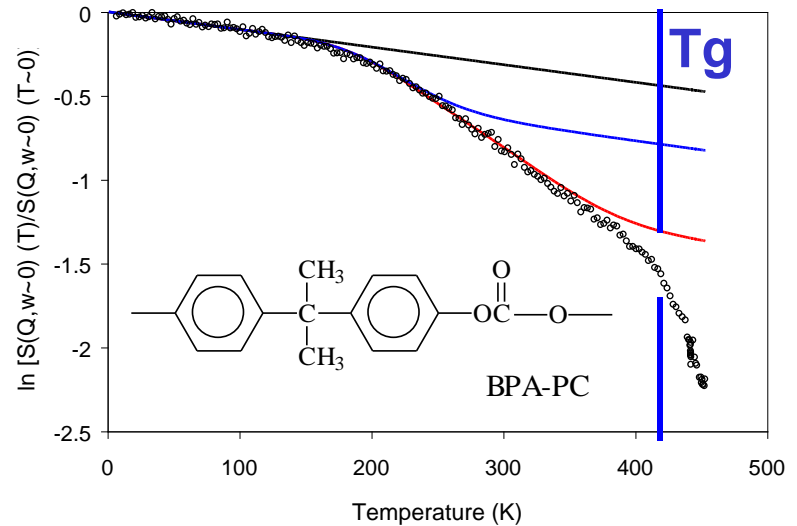


Ephenyl~37
kJ/mol $\sigma_1 \sim 6$

Ech3=15 kJ/mol
 $\sigma_1 \sim 3$

compatible with TMPC

(after Spiess et al. 1987)



Distribution?

Glassy
polymers:

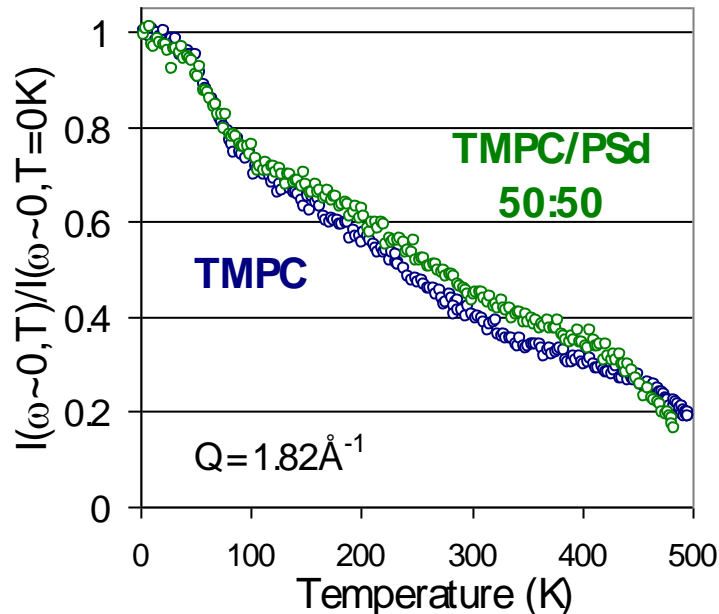
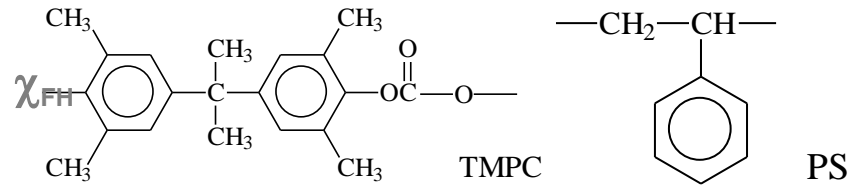
backbone chain conformation

Structural disorder

Blending

≠ inter-molecular potential.
≈ intra-

TMPC: only PC miscible with PS, large χ_{FH}



1st step: no resolvable perturbation

2nd step: broadened distribution

intramolecular environment

- average E_A
- architectural considerations

intermolecular → limited effect on σ

Conclusions: CASE STUDY

Characterisation local dynamics of PCs:

two architectures → toughest (BPA-PC) & most brittle (TMPC)

Technique combined backscattering window scans, inelastic BS & TOF

TMPC exhibits two methyl relaxations of rather different distribution of potentials

Blending affects $\sigma(E_A)$

BPA-PC Phenyl + methyl

